

# Semiclassical Transitions With Black Hole in the Intermediate State

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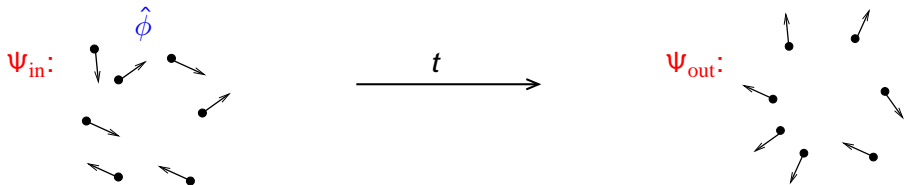


Ginzburg Conference on Physics,  
Lebedev Institute, May 29, 2012

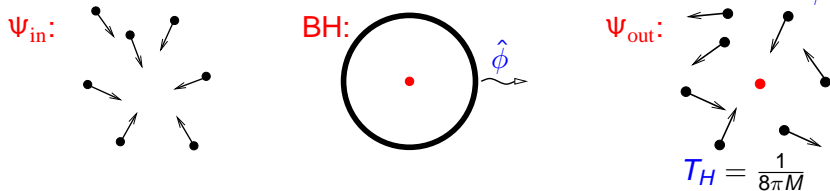
# Hawking effect and asymptotic states

Scattering:

*Hawking '75*



Collapse & Evaporation:

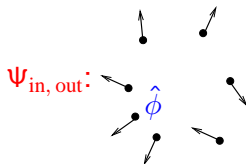


Gravitational  $\hat{S}$ -matrix:  $\Psi_{\text{out}} = \hat{S} \Psi_{\text{in}} ?$   
 flat  $\leftrightarrow$  flat

*'t Hooft '85*

# Why S-matrix?

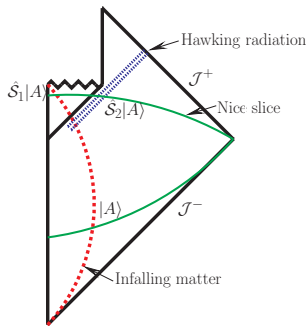
- $\Psi_{\text{in, out}}$  are well understood.



- Information paradox

$$\hat{S}^\dagger \hat{S} \neq 0 ?$$

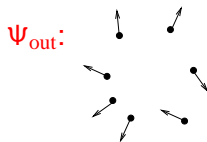
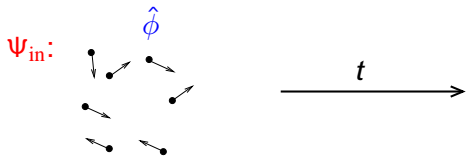
It is worth calculating  $\hat{S}$  !



Hand-waving argument!

# Calculating $\hat{S}$ -matrix elements

Scattering:

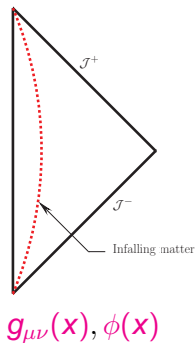


$$\langle \Psi_{out} | \hat{S} | \Psi_{in} \rangle = \int dg_{\mu\nu} d\phi \Psi_{out}^* e^{i(S_g[g] + S_\phi[\phi])} \Psi_{in}$$

Semiclassics:  $S_g + S_\phi \gg 1 \Rightarrow$   $g_{\mu\nu}(x), \phi(x)$   
saddle point

$$\frac{\delta}{\delta g_{\mu\nu}}, \frac{\delta}{\delta \phi} [S_g + S_\phi - i \ln \Psi_{in, out}] = 0$$

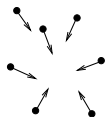
$$\Rightarrow \langle \Psi_{out} | \hat{S} | \Psi_{in} \rangle = \mathcal{A} \cdot e^{i(S_g[g] + S_\phi[\phi]) + \ln \Psi_{out} + \ln \Psi_{in}}$$



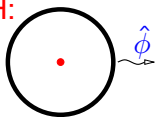
# Calculating $\hat{S}$ -matrix elements

## Collapse & Evaporation:

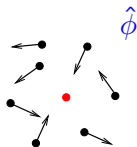
$\Psi_{\text{in}}$ :



BH:



$\Psi_{\text{out}}$ :

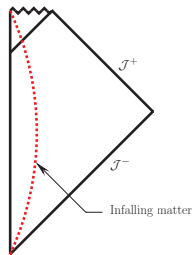


## The same integral:

$$\langle \Psi_{\text{out}} | \hat{S} | \Psi_{\text{in}} \rangle = \int dg_{\mu\nu} d\phi \Psi_{\text{out}}^* e^{i(S_g[g] + S_\phi[\phi])} \Psi_{\text{in}}$$

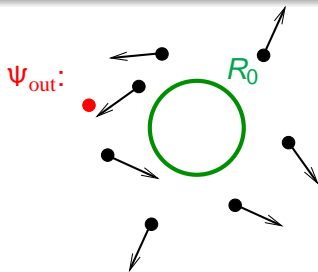
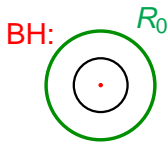
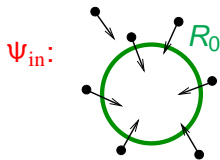
**BUT:** No saddle point  $g_{\mu\nu}(x), \phi(x)$   
describing complete evolution !

$$\langle \Psi_{\text{out}} | \hat{S} | \Psi_{\text{in}} \rangle = ?$$



$g_{\mu\nu}(x), \phi(x)$

# Adding constraint

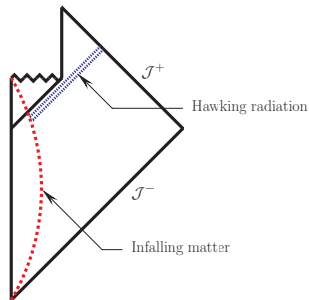


$$\Sigma[g, \phi] = \int dt M_{R_0}(t)$$

Add constraint  $\Sigma[g, \phi] = \Sigma_0$   
to GR equations



Correct asymptotics of solution!



$$g_{\mu\nu}(x), \phi(x)$$

# Faddeev-Popov method

$$1 = \int_0^\infty d\Sigma_0 \delta(\Sigma_0 - \Sigma[g, \phi]) = \int_0^\infty d\Sigma_0 \int_{-i\infty}^{i\infty} d\epsilon e^{\epsilon(\Sigma_0 - \Sigma[g, \phi])}$$
$$\langle \Psi_{\text{in}} | \hat{\mathcal{S}} | \Psi_{\text{out}} \rangle = \int dg d\phi \cdot e^{i(S_g + S_\phi)} \cdot \Psi_{\text{in}} \Psi_{\text{out}}^*$$

## Strategy:

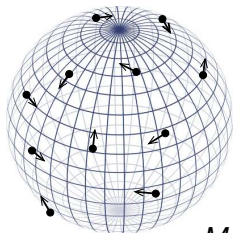
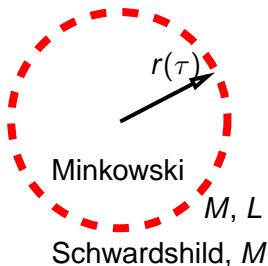
Bezrukov, DL '03

DL, Panin, Sibiryakov '07

- 1 Evaluate path integral at **fixed**  $\Sigma_0$ .  
 $\Psi_{\text{in}}$  and  $\Psi_{\text{out}}$  live in **flat space!**

$$S_\epsilon = S + i\epsilon\Sigma$$

- 2 Evaluate integral over  $\epsilon$ :  $\Sigma_0 = \Sigma[g, \phi]$
- 3 Integrate over  $\Sigma_0$ :  $\epsilon \rightarrow +0$ .
- 4  $\mathcal{P} = \sum_{\text{out}} |\langle \Psi_{\text{in}} | \hat{\mathcal{S}} | \Psi_{\text{out}} \rangle|^2 \approx \lim_{\epsilon \rightarrow 0} \exp[-2\text{Im}(S_\epsilon - i \ln \Psi_i)]$



$M, L, m = 0$

$$S_m = - \int d\tau \sqrt{m^2 + \frac{L^2}{r^2}}$$

$$S_g = \frac{1}{16\pi} \int d^4x \mathcal{R} \sqrt{-g} + \text{boundary terms}$$

$$\left(\frac{dr}{d\tau}\right)^2 - \left(\frac{rM}{L} + \frac{L}{2r^2}\right)^2 + 1 = 0$$

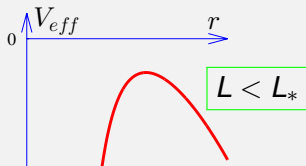
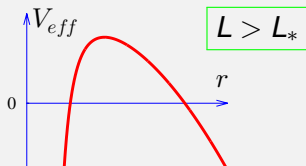
$V_{\text{eff}}$



# Shell evolution

$$\epsilon = 0$$

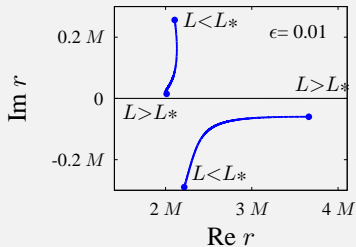
$$V_{eff} = - \left( \frac{rM}{L} + \frac{L}{2r^2} \right)^2 + 1 = 0$$



$$\epsilon \neq 0$$

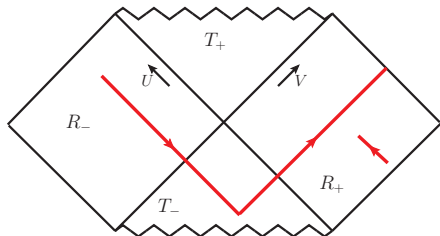
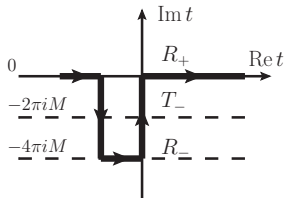
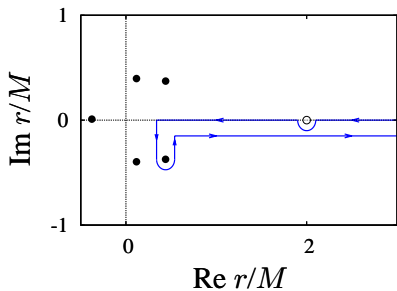
$$S_\epsilon = S + i\epsilon\Sigma$$

$$\Sigma[g, r] = \int dt \theta(R_0 - r(\tau))$$



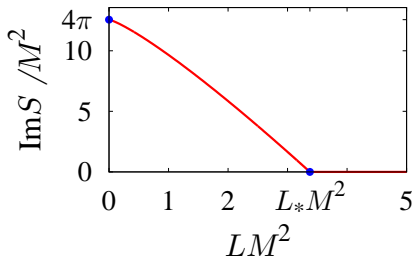
Always reflection!

# Spacetime structure



$U, V$  – Kruskal coordinates

Result:  $\epsilon \rightarrow 0$



Probability at  $L = 0$ :  $\mathcal{P} = \lim_{\epsilon \rightarrow +0} e^{-2\text{Im} S_\epsilon + \text{b.t.}} = e^{-4\pi M^2}$

$$\mathcal{P} = \exp(-A/4) = \exp(-S_B)$$

*Parikh, Wilczek '00*

## New semiclassical method:

- Derived from first principles.
- Gives  $\langle \Psi_{\text{out}} | \hat{S} | \Psi_{\text{in}} \rangle$ .
- Includes backreaction.
- Easy to implement: solve classical Eqs. with complex terms.
- Can be applied to complex models:  $\phi(r, t)$ ,  $g_{\mu\nu}(r, t)$ .

*Cf. Berezin, Boyarsky, Neronov '99*

## Dust shell:

- Our result is **consistent** with Hawking radiation & Parikh, Wilzcek & Bekeinstein entropy.
- Black hole is a true energy eigenstate:  $\mathcal{P} \neq 1$

*Cf. Vilkovisky '06*

**Other models?**