# A Comparative Study of Laplacians in Riemannian and Antisymplectic Geometry 

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## A Comparative Study of Laplacians in Riemannian and Antisymplectic Geometry

(1) Anti-Poisson Geometry
(2) Riemannian Geometry
(3) Conclusions

## Darboux Coordinates

## Poisson

## Darboux Coordinates

## Poisson

## Coordinates <br> Momenta <br> $p_{i}$

## Darboux Coordinates

Poisson

| Coordinates | $q^{i}$ | Boson | Fermion |
| :---: | :---: | :---: | :---: |
| Momenta | $p_{i}$ | $\downarrow$ | Boson |
| Fermion |  |  |  |

## Darboux Coordinates

Poisson

| Coordinates | $q^{i}$ | Boson | Fermion |
| :---: | :---: | :---: | :---: |
|  |  | $\downarrow$ | $\uparrow$ |
| Momenta | $p_{i}$ | Boson | Fermion |

## Anti-Poisson

## Darboux Coordinates

Poisson

| Coordinates | $q^{i}$ | Boson | Fermion |
| :---: | :---: | :---: | :---: |
|  |  | $\downarrow$ | $\uparrow$ |
| Momenta | $p_{i}$ | Boson | Fermion |

## Anti-Poisson

## Fields $\phi^{\alpha}$ <br> Antifields $\phi_{\alpha}^{*}$

## Darboux Coordinates

Poisson

| Coordinates | $q^{i}$ | Boson | Fermion |
| :---: | :---: | :---: | :---: |
|  |  | $\downarrow$ | $\uparrow$ |
| Momenta | $p_{i}$ | Boson | Fermion |

Anti-Poisson

| Fields | $\phi^{\alpha}$ | Boson | Fermion |
| :---: | :---: | :---: | :---: |
|  |  | $\uparrow$ | $\downarrow$ |
| Antifields | $\phi_{\alpha}^{*}$ | Fermion | Boson |

## Darboux Coordinates

## Poisson

| Coordinates | $q^{i}$ | Boson | Fermion |
| :---: | :---: | :---: | :---: |
|  |  | $\uparrow$ | $\uparrow$ |
| Momenta | $p_{i}$ | Boson | Fermion |

## Anti-Poisson

| Fields | $\phi^{\alpha}$ | Boson | Fermion |
| :---: | :---: | :---: | :---: |
|  |  | $\downarrow$ | $\downarrow$ |
| Antifields | $\phi_{\alpha}^{*}$ | Fermion | Boson |

## Darboux Coordinates

## Poisson

| Coordinates | $q^{i}$ | Boson | Fermion |
| :---: | :---: | :---: | :---: |
|  |  | $\uparrow$ | $\uparrow$ |
| Momenta | $p_{i}$ | Boson | Fermion |

## Poisson Bracket $\{,\}_{\text {PB }}$

$$
\begin{aligned}
\left\{q^{i}, q^{j}\right\}_{P B} & =0 \\
\left\{q^{i}, p_{j}\right\}_{P B} & =\delta_{j}^{i} \\
\left\{p_{i}, p_{j}\right\}_{P B} & =0
\end{aligned}
$$

## Anti-Poisson

| Fields | $\phi^{\alpha}$ | Boson | Fermion |
| :---: | :---: | :---: | :---: |
|  |  | $\uparrow$ | $\downarrow$ |
| Antifields | $\phi_{\alpha}^{*}$ | Fermion | Boson |

## Darboux Coordinates

## Poisson

| Coordinates | $q^{i}$ | Boson | Fermion |
| :---: | :---: | :---: | :---: |
|  |  | $\uparrow$ | $\uparrow$ |
| Momenta | $p_{i}$ | Boson | Fermion |

## Anti-Poisson

| Fields | $\phi^{\alpha}$ | Boson | Fermion |
| :---: | :---: | :---: | :---: |
|  |  | $\uparrow$ | $\uparrow$ |

## Poisson Bracket $\{,\}_{P B}$

$$
\begin{aligned}
& \left\{q^{i}, q^{j}\right\}_{P B}=0 \\
& \left\{q^{i}, p_{j}\right\}_{P B}=\delta_{j}^{i} \\
& \left\{p_{i}, p_{j}\right\}_{P B}=0
\end{aligned}
$$

## Antibracket (, ) $A B$

Antifields $\phi_{\alpha}^{*}$ Fermion Boson

## Darboux Coordinates

## Poisson

| Coordinates | $q^{i}$ | Boson | Fermion |
| :---: | :---: | :---: | :---: |
| Momenta | $p_{i}$ | $\downarrow$ | $\uparrow$ |
| Boson | Fermion |  |  |

## Anti-Poisson

| Fields | $\phi^{\alpha}$ | Boson | Fermion |
| :---: | :---: | :---: | :---: |
|  |  | $\uparrow$ | $\uparrow$ |
| Antifields | $\phi^{*}$ | Fermion | Boson |

## Poisson Bracket $\{,\}_{P B}$

$$
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& \left\{q^{i}, q^{j}\right\}_{P B}=0 \\
& \left\{q^{i}, p_{j}\right\}_{P B}=\delta_{j}^{i} \\
& \left\{p_{i}, p_{j}\right\}_{P B}=0
\end{aligned}
$$

## Antibracket (, ) ${ }_{A B}$

$$
\begin{aligned}
\left(\phi^{\alpha}, \phi^{\beta}\right)_{A B} & =0 \\
\left(\phi^{\alpha}, \phi_{\beta}^{*}\right)_{A B} & =\delta_{\beta}^{\alpha} \\
\left(\phi_{\alpha}^{*}, \phi_{\beta}^{*}\right)_{A B} & =0
\end{aligned}
$$

## Darboux Coordinates

## Poisson

| Coordinates | $q^{i}$ | Boson | Fermion |
| :---: | :---: | :---: | :---: |
| Momenta | $p_{i}$ | Boson | $\uparrow$ |
| Fermion |  |  |  |

## Poisson Bracket $\{,\}_{P B}$

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& \left\{q^{i}, q^{j}\right\}_{P B}=0 \\
& \left\{q^{i}, p_{j}\right\}_{P B}=\delta_{j}^{i} \\
& \left\{p_{i}, p_{j}\right\}_{P B}=0
\end{aligned}
$$

## "Comma is a Boson"

## Antibracket (, ) $A B$

$$
\begin{aligned}
\left(\phi^{\alpha}, \phi^{\beta}\right)_{A B} & =0 \\
\left(\phi^{\alpha}, \phi_{\beta}^{*}\right)_{A B} & =\delta_{\beta}^{\alpha} \\
\left(\phi_{\alpha}^{*}, \phi_{\beta}^{*}\right)_{A B} & =0
\end{aligned}
$$

## Darboux Coordinates

## Poisson

| Coordinates | $q^{i}$ | Boson | Fermion |
| :---: | :---: | :---: | :---: |
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| Momenta | $p_{i}$ | Boson | Fermion |

## Poisson Bracket $\{,\}_{P B}$

$$
\begin{aligned}
& \left\{q^{i}, q^{j}\right\}_{P B}=0 \\
& \left\{q^{i}, p_{j}\right\}_{P B}=\delta_{j}^{i} \\
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\end{aligned}
$$

## "Comma is a Boson"

## Antibracket (, ) ${ }_{A B}$

$$
\begin{aligned}
\left(\phi^{\alpha}, \phi^{\beta}\right)_{A B} & =0 \\
\left(\phi^{\alpha}, \phi_{\beta}^{*}\right)_{A B} & =\delta_{\beta}^{\alpha} \\
\left(\phi_{\alpha}^{*}, \phi_{\beta}^{*}\right)_{A B} & =0
\end{aligned}
$$

## "Comma is a Fermion"

## General Coordinates

## General Coordinates

Poisson Bracket

$$
\begin{aligned}
\left\{z^{A}, z^{B}\right\}_{P B} & =\omega^{A B} \\
\{F, G\}_{P B} & =\left(F \frac{\overleftarrow{\partial^{r}}}{\partial z^{A}}\right) \omega^{A B}\left(\frac{\overrightarrow{\partial^{\ell}}}{\partial z^{B}} G\right)
\end{aligned}
$$

## General Coordinates

Poisson Bracket

$$
\begin{aligned}
\left\{z^{A}, z^{B}\right\}_{P B} & =\omega^{A B} \\
\{F, G\}_{P B} & =\left(F \frac{\overleftarrow{\partial^{r}}}{\partial z^{A}}\right) \omega^{A B}\left(\frac{\overrightarrow{\partial^{\ell}}}{\partial z^{B}} G\right)
\end{aligned}
$$

## Antibracket

$$
\begin{aligned}
\left(z^{A}, z^{B}\right)_{A B} & =E^{A B} \\
(F, G)_{A B} & =\left(F \frac{\overleftarrow{\partial^{r}}}{\partial z^{A}}\right) E^{A B}\left(\frac{\overrightarrow{\partial^{\ell}}}{\partial z^{B}} G\right)
\end{aligned}
$$

## General Coordinates

Poisson Bracket

$$
\begin{aligned}
\left\{z^{A}, z^{B}\right\}_{P B} & =\omega^{A B} \\
\{F, G\}_{P B} & =\left(F \frac{\overleftarrow{\partial^{r}}}{\partial z^{A}}\right) \omega^{A B}\left(\frac{\overrightarrow{\partial^{\ell}}}{\partial z^{B}} G\right)
\end{aligned}
$$

Grassmann-parity

$$
\varepsilon\left(\omega^{A B}\right)=\varepsilon_{A}+\varepsilon_{B}
$$

## Even

## Antibracket

$$
\begin{aligned}
\left(z^{A}, z^{B}\right)_{A B} & =E^{A B} \\
(F, G)_{A B} & =\left(F \frac{\overleftarrow{\partial^{r}}}{\partial z^{A}}\right) E^{A B}\left(\frac{\overrightarrow{\partial^{\ell}}}{\partial z^{B}} G\right)
\end{aligned}
$$

## General Coordinates

Poisson Bracket

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\left\{z^{A}, z^{B}\right\}_{P B} & =\omega^{A B} \\
\{F, G\}_{P B} & =\left(F \frac{\overleftarrow{\partial^{r}}}{\partial z^{A}}\right) \omega^{A B}\left(\frac{\overrightarrow{\partial^{\ell}}}{\partial z^{B}} G\right)
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Grassmann-parity

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\varepsilon\left(\omega^{A B}\right)=\varepsilon_{A}+\varepsilon_{B}
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\begin{aligned}
\left(z^{A}, z^{B}\right)_{A B} & =E^{A B} \\
(F, G)_{A B} & =\left(F \frac{\overleftarrow{\partial^{r}}}{\partial z^{A}}\right) E^{A B}\left(\frac{\overrightarrow{\partial^{\ell}}}{\partial z^{B}} G\right)
\end{aligned}
$$

Grassmann-parity

$$
\varepsilon\left(E^{A B}\right)=\varepsilon_{A}+\varepsilon_{B}+1
$$

Odd

## General Coordinates

## General Coordinates

## Poisson Case

$\omega^{B A}=-(-1)^{\varepsilon_{A} \varepsilon_{B}} \omega^{A B}$

- Antisymmetric


## General Coordinates

Poisson Case

$$
\omega^{B A}=-(-1)^{\varepsilon_{A} \varepsilon_{B}} \omega^{A B}
$$

- Antisymmetric


## Inverse 2-form <br> $\omega=\frac{1}{2} d z^{A} \omega_{A B} \wedge d z^{B}$

$\omega_{B A}=(-1)^{\left(\varepsilon_{A}+1\right)\left(\varepsilon_{B}+1\right)} \omega_{A B}$

- Skewsymmetric


## General Coordinates

Poisson Case

$$
\omega^{B A}=-(-1)^{\varepsilon_{A} \varepsilon_{B}} \omega^{A B}
$$

- Antisymmetric


## Anti-Poisson Case

$E^{B A}=-(-1)^{\left(\varepsilon_{A}+1\right)\left(\varepsilon_{B}+1\right)} E^{A B}$

$$
\omega_{B A}=(-1)^{\left(\varepsilon_{A}+1\right)\left(\varepsilon_{B}+1\right)} \omega_{A B}
$$

- Skewsymmetric

Inverse 2-form $\omega=\frac{1}{2} d z^{A} \omega_{A B} \wedge d z^{B}$

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## General Coordinates

## Poisson Case

$$
\omega^{B A}=-(-1)^{\varepsilon_{A} \varepsilon_{B}} \omega^{A B}
$$

- Antisymmetric

$$
\text { Inverse 2-form } \omega=\frac{1}{2} d z^{A} \omega_{A B} \wedge d z^{B}
$$

$$
\omega_{B A}=(-1)^{\left(\varepsilon_{A}+1\right)\left(\varepsilon_{B}+1\right)} \omega_{A B}
$$

- Skewsymmetric


## Anti-Poisson Case

$E^{B A}=-(-1)^{\left(\varepsilon_{A}+1\right)\left(\varepsilon_{B}+1\right)} E^{A B}$

- Antiskewsymmetric
- Morally Symmetric like the Riemannian Case


## General Coordinates

## Poisson Case

$$
\omega^{B A}=-(-1)^{\varepsilon_{A} \varepsilon_{B}} \omega^{A B}
$$

- Antisymmetric

$$
\text { Inverse 2-form } \omega=\frac{1}{2} d z^{A} \omega_{A B} \wedge d z^{B}
$$

$$
\omega_{B A}=(-1)^{\left(\varepsilon_{A}+1\right)\left(\varepsilon_{B}+1\right)} \omega_{A B}
$$

- Skewsymmetric


## Anti-Poisson Case

$$
E^{B A}=-(-1)^{\left(\varepsilon_{A}+1\right)\left(\varepsilon_{B}+1\right)} E^{A B}
$$

- Antiskewsymmetric
- Morally Symmetric like the Riemannian Case


## Inverse 2-form

$E=\frac{1}{2} d z^{A} E_{A B} \wedge d z^{B}$

$$
E_{B A}=-(-1)^{\varepsilon_{A} \varepsilon_{B}} E_{A B}
$$

- Antisymmetric


## General Coordinates

## Poisson Case

$$
\omega^{B A}=-(-1)^{\varepsilon_{A} \varepsilon_{B}} \omega^{A B}
$$

- Antisymmetric

$$
\text { Inverse 2-form } \omega=\frac{1}{2} d z^{A} \omega_{A B} \wedge d z^{B}
$$

$$
\omega_{B A}=(-1)^{\left(\varepsilon_{A}+1\right)\left(\varepsilon_{B}+1\right)} \omega_{A B}
$$

- Skewsymmetric

Inverse 2-form $E=\frac{1}{2} d z^{A} E_{A B} \wedge d z^{B}$

$$
E_{B A}=-(-1)^{\varepsilon_{A} \varepsilon_{B}} E_{A B}
$$

- Antisymmetric
- Morally

Skewsymmetric like the Symplectic Case

## Jacobi Identity and Closeness Relation

## Jacobi Identity for Poisson Bracket

$$
\sum_{\text {cycl. } f, g, h}(-1)^{\varepsilon_{f} \varepsilon_{h}}\{f,\{g, h\}\}=0
$$

## Jacobi Identity and Closeness Relation

$$
\begin{aligned}
& \text { Jacobi Identity for Poisson Bracket } \\
& \qquad \sum_{\text {cycl. } f, g, h}(-1)^{\varepsilon_{f} \varepsilon_{h}}\{f,\{g, h\}\}=0
\end{aligned}
$$

## Jacobi Identity for Antibracket

$$
\sum(-1)^{\left(\varepsilon_{f}+1\right)\left(\varepsilon_{h}+1\right)}(f,(g, h))=0
$$ cycl. $f, g, h$

## Jacobi Identity and Closeness Relation

## Jacobi Identity for Poisson Bracket

$$
\sum(-1)^{\varepsilon_{f} \varepsilon_{h}}\{f,\{g, h\}\}=0
$$ cycl. $f, g, h$

Jacobi Identity for Antibracket

$$
\sum(-1)^{\left(\varepsilon_{f}+1\right)\left(\varepsilon_{h}+1\right)}(f,(g, h))=0
$$

$$
\operatorname{cycl.} f, g, h
$$

## Symplectic Case <br> Closed 2-form:

$$
d \omega=0
$$

## Jacobi Identity and Closeness Relation

## Jacobi Identity for Poisson Bracket

$$
\sum(-1)^{\varepsilon_{f} \varepsilon_{h}}\{f,\{g, h\}\}=0
$$

$$
\operatorname{cycl.} f, g, h
$$

## Symplectic Case

Closed 2-form:

$$
d \omega=0
$$

## Anti-Symplectic Case

Closed 2-form:

$$
d E=0
$$

## Jacobi Identity and Closeness Relation

## Jacobi Identity for Poisson Bracket

$$
\sum(-1)^{\varepsilon_{f} \varepsilon_{h}}\{f,\{g, h\}\}=0
$$

$$
\operatorname{cycl.f}, g, h
$$

## Const. rank $\Rightarrow$ Darboux Thm.

Jacobi Identity for Antibracket

## Symplectic Case

Closed 2-form:

$$
d \omega=0
$$

Anti-Symplectic Case Closed 2-form:

$$
d E=0
$$

cycl. $f, g, h$

## Const. rank $\Rightarrow$ Darboux Thm.

## Jacobi Identity and Closeness Relation

Jacobi Identity for Poisson Bracket

$$
\sum_{\text {cycl. } f, g, h}(-1)^{\varepsilon_{f} \varepsilon_{h}}\{f,\{g, h\}\}=0
$$

Const. rank $\Rightarrow$ Darboux Thm.
Jacobi Identity for Antibracket

## Symplectic Case

Closed 2-form:

$$
d \omega=0
$$

Anti-Symplectic Case Closed 2-form:

$$
d E=0
$$

$$
(-1)^{\left(\varepsilon_{f}+1\right)\left(\varepsilon_{h}+1\right)}(f,(g, h))=0
$$

Const. rank $\Rightarrow$ Darboux Thm.

## Riemannian Case

No Jacobi Identity, no Closeness Relation, and no Darboux Thm.

## Laplacian

## Even Laplacian in Riemannian Case

$$
\Delta_{\rho}=\frac{(-1)^{\varepsilon_{A}}}{\rho} \frac{\overrightarrow{\partial^{\ell}}}{\partial z^{A}} \rho g^{A B} \frac{\overrightarrow{\partial^{\ell}}}{\partial z^{B}}
$$

## Laplacian

## Even Laplacian in Riemannian Case

$$
\Delta_{\rho}=\frac{(-1)^{\varepsilon_{A}}}{\rho} \frac{\overrightarrow{\partial^{\ell}}}{\partial z^{A}} \rho g^{A B} \frac{\overrightarrow{\partial^{\ell}}}{\partial z^{B}}
$$

## Odd Laplacian in Anti-Poisson Case

$$
2 \Delta_{\rho}=\frac{(-1)^{\varepsilon_{A}}}{\rho} \frac{\overrightarrow{\partial^{\ell}}}{\partial z^{A}} \rho E^{A B} \frac{\overrightarrow{\partial^{\ell}}}{\partial z^{B}}
$$

## Laplacian

## Even Laplacian in Riemannian Case

$$
\Delta_{\rho}=\frac{(-1)^{\varepsilon_{A}}}{\rho} \frac{\overrightarrow{\partial^{\ell}}}{\partial z^{A}} \rho g^{A B} \frac{\overrightarrow{\partial^{\ell}}}{\partial z^{B}}
$$

- Canonical density

$$
\rho_{g}:=\sqrt{g}:=\sqrt{\operatorname{sdet}\left(g_{A B}\right)}
$$

Odd Laplacian in Anti-Poisson Case

$$
2 \Delta_{\rho}=\frac{(-1)^{\varepsilon_{A}}}{\rho} \frac{\overrightarrow{\partial^{\ell}}}{\partial z^{A}} \rho^{A B} \frac{\overrightarrow{\partial^{\ell}}}{\partial z^{B}}
$$

## Laplacian

## Even Laplacian in Riemannian Case

$$
\Delta_{\rho}=\frac{(-1)^{\varepsilon_{A}}}{\rho} \frac{\overrightarrow{\partial^{\ell}}}{\partial z^{A}} \rho g^{A B} \frac{\overrightarrow{\partial^{\ell}}}{\partial z^{B}}
$$

- Canonical density $\rho_{g}:=\sqrt{g}:=\sqrt{\operatorname{sdet}\left(g_{A B}\right)}$

Odd Laplacian in Anti-Poisson Case

$$
2 \Delta_{\rho}=\frac{(-1)^{\varepsilon_{A}}}{\rho} \frac{\overrightarrow{\partial^{\ell}}}{\partial z^{A}} \rho E^{A B} \frac{\overrightarrow{\partial^{\ell}}}{\partial z^{B}}
$$

- No canonical density!


## Laplacian

## Even Laplacian in Riemannian Case

$$
\Delta_{\rho}=\frac{(-1)^{\varepsilon_{A}}}{\rho} \frac{\overrightarrow{\partial^{\ell}}}{\partial z^{A}} \rho g^{A B} \frac{\overrightarrow{\partial^{\ell}}}{\partial z^{B}}
$$

- Canonical density $\rho_{g}:=\sqrt{g}:=\sqrt{\operatorname{sdet}\left(g_{A B}\right)}$
- $\Delta_{\rho}^{2}$ is a 4th-order operator.


## Odd Laplacian in Anti-Poisson Case

$$
2 \Delta_{\rho}=\frac{(-1)^{\varepsilon_{A}}}{\rho} \frac{\overrightarrow{\partial^{\ell}}}{\partial z^{A}} \rho E^{A B} \frac{\overrightarrow{\partial^{\ell}}}{\partial z^{B}}
$$

- No canonical density!


## Laplacian

## Even Laplacian in Riemannian Case

$$
\Delta_{\rho}=\frac{(-1)^{\varepsilon_{A}}}{\rho} \frac{\overrightarrow{\partial^{\ell}}}{\partial z^{A}} \rho g^{A B} \frac{\overrightarrow{\partial^{\ell}}}{\partial z^{B}}
$$

- Canonical density $\rho_{g}:=\sqrt{g}:=\sqrt{\operatorname{sdet}\left(g_{A B}\right)}$
- $\Delta_{\rho}^{2}$ is a 4th-order operator.

Odd Laplacian in Anti-Poisson Case

$$
2 \Delta_{\rho}=\frac{(-1)^{\varepsilon_{A}}}{\rho} \frac{\overrightarrow{\partial^{\ell}}}{\partial z^{A}} \rho E^{A B} \frac{\overrightarrow{\partial^{\ell}}}{\partial z^{B}}
$$

- No canonical density!
- $\Delta_{\rho}^{2}$ is a 1st-order operator.


## Laplacian

## Even Laplacian in Riemannian Case

$$
\Delta_{\rho}=\frac{(-1)^{\varepsilon_{A}}}{\rho} \frac{\overrightarrow{\partial^{\ell}}}{\partial z^{A}} \rho g^{A B} \frac{\overrightarrow{\partial^{\ell}}}{\partial z^{B}}
$$

- Canonical density $\rho_{g}:=\sqrt{g}:=\sqrt{\operatorname{sdet}\left(g_{A B}\right)}$
- $\Delta_{\rho}^{2}$ is a 4th-order operator.

Odd Laplacian in Anti-Poisson Case

$$
2 \Delta_{\rho}=\frac{(-1)^{\varepsilon_{A}}}{\rho} \frac{\overrightarrow{\partial^{\ell}}}{\partial z^{A}} \rho E^{A B} \frac{\overrightarrow{\partial^{\ell}}}{\partial z^{B}}
$$

- No canonical density!
- $\Delta_{\rho}^{2}$ is a 1 st-order operator.
- When is $\Delta_{\rho}^{2}=0$ nilpotent?


## Laplacian

## Even Laplacian in Riemannian Case

$$
\Delta_{\rho}=\frac{(-1)^{\varepsilon_{A}}}{\rho} \frac{\overrightarrow{\partial^{\ell}}}{\partial z^{A}} \rho g^{A B} \frac{\overrightarrow{\partial^{\ell}}}{\partial z^{B}}
$$

- Canonical density $\rho_{g}:=\sqrt{g}:=\sqrt{\operatorname{sdet}\left(g_{A B}\right)}$
- $\Delta_{\rho}^{2}$ is a 4th-order operator.

Odd Laplacian in Anti-Poisson Case

$$
2 \Delta_{\rho}=\frac{(-1)^{\varepsilon_{A}}}{\rho} \frac{\overrightarrow{\partial^{\ell}}}{\partial z^{A}} \rho E^{A B} \frac{\overrightarrow{\partial^{\ell}}}{\partial z^{B}}
$$

- No canonical density!
- $\Delta_{\rho}^{2}$ is a 1 st-order operator.
- When is $\Delta_{\rho}^{2}=0$ nilpotent?


## The $2 \times 2=4$ Classical Geometries and their Symmetries

|  | Even Geometry | Odd Geometry |
| :---: | :---: | :---: |
|  | $g=d z^{A} g_{A B} \vee d z^{B}$ | $g=d z^{A} g_{A B} \vee d z^{B}$ |
| Riemannian | $\varepsilon\left(g_{A B}\right)=\varepsilon_{A}+\varepsilon_{B}$ | $\varepsilon\left(g_{A B}\right)=\varepsilon_{A}+\varepsilon_{B}+1$ |
| Covariant | $g_{B A}=-(-1)^{\left(\varepsilon_{A}+1\right)\left(\varepsilon_{B}+1\right)} g_{A B}$ | $g_{B A}=(-1)^{\varepsilon_{A} \varepsilon_{B}} g_{A B}$ |
| Metric | Antiskewsymmetric | Symmetric |
|  | No Closeness Relation | No Closeness Relation |
| Inverse | $\varepsilon\left(g^{A B}\right)=\varepsilon_{A}+\varepsilon_{B}$ | $\varepsilon\left(g^{A B}\right)=\varepsilon_{A}+\varepsilon_{B}+1$ |
| Riemannian | $g^{B A}=(-1)^{\varepsilon_{A} \varepsilon_{B}} g^{A B}$ | $g^{B A}=(-1)^{\left(\varepsilon_{A}+1\right)\left(\varepsilon_{B}+1\right)} g^{A B}$ |
| Contravariant | Symmetric | Skewsymmetric |
| Metric | Even Laplacian | No Laplacian |
|  | $\omega=\frac{1}{2} d z^{A} \omega_{A B} \wedge d z^{B}$ | $E=\frac{1}{2} d z^{A} E_{A B} \wedge d z^{B}$ |
| Symplectic | $\varepsilon\left(\omega_{A B}\right)=\varepsilon_{A}+\varepsilon_{B}$ | $\varepsilon\left(E_{A B}\right)=\varepsilon_{A}+\varepsilon_{B}+1$ |
| Covariant | $\omega_{B A}=(-1)^{\left(\varepsilon_{A}+1\right)\left(\varepsilon_{B}+1\right)} \omega_{A B}$ | $E_{B A}=-(-1)^{\varepsilon_{A} \varepsilon_{B} E_{A B}}$ |
| Two-Form | Skewsymmetric | Antisymmetric |
|  | Closeness Relation | Closeness Relation |
| Inverse | $\varepsilon\left(\omega^{A B}\right)=\varepsilon_{A}+\varepsilon_{B}$ | $\varepsilon\left(E^{A B}\right)=\varepsilon_{A}+\varepsilon_{B}+1$ |
| Symplectic | $\omega^{B A}=-(-1)^{\varepsilon} \varepsilon_{A} \varepsilon_{B} \omega_{A B}$ | $E^{B A}=-(-1)^{\left(\varepsilon_{A}+1\right)\left(\varepsilon_{B}+1\right)} E^{A B}$ |
| Contravariant | Antisymmetric | Symmetric |
| Tensor | No Laplacian | Odd Laplacian |

## The $2 \times 2=4$ Classical Geometries and their Symmetries

|  | Even Geometry | Odd Geometry |
| :---: | :---: | :---: |
|  | $g=d z^{A} g_{A B} \vee d z^{B}$ | $g=d z^{A} g_{A B} \vee d z^{B}$ |
| Riemannian | $\varepsilon\left(g_{A B}\right)=\varepsilon_{A}+\varepsilon_{B}$ | $\varepsilon\left(g_{A B}\right)=\varepsilon_{A}+\varepsilon_{B}+1$ |
| Covariant | $g_{B A}=-(-1)^{\left(\varepsilon_{A}+1\right)\left(\varepsilon_{B}+1\right)} g_{A B}$ | $g_{B A}=(-1)^{\varepsilon} \varepsilon_{B} \varepsilon_{B} g_{A B}$ |
| Metric | Antiskewsymmetric | Symmetric |
|  | No Closeness Relation | No Closeness Relation |
| Inverse | $\varepsilon\left(g^{A B}\right)=\varepsilon_{A}+\varepsilon_{B}$ | $\varepsilon\left(g^{A B}\right)=\varepsilon_{A}+\varepsilon_{B}+1$ |
| Riemannian | $g^{B A}=(-1)^{\varepsilon_{A} \varepsilon_{B}} g^{A B}$ | $g^{B A}=(-1)^{\left(\varepsilon_{A}+1\right)\left(\varepsilon_{B}+1\right)} g^{A B}$ |
| Contravariant | Symmetric | Skewsymmetric |
| Metric | Even Laplacian | No Laplacian |
| Symplectic | $\omega=\frac{1}{2} d z^{A} \omega_{A B} \wedge d z^{B}$ | $E=\frac{1}{2} d z^{A} E_{A B} \wedge d z^{B}$ |
| Covariant | $\varepsilon\left(\omega_{A B}\right)=\varepsilon_{A}+\varepsilon_{B}$ | $\varepsilon\left(E_{A B}\right)=\varepsilon_{A}+\varepsilon_{B}+1$ |
| Two-Form | $\omega_{B A}=(-1)^{\left(\varepsilon_{A}+1\right)\left(\varepsilon_{B}+1\right)} \omega_{A B}$ | $E_{B A}=-(-1)_{A} \varepsilon_{B} E_{A B}$ |
|  | Skewsymmetric | Antisymmetric |
| Closeness Relation | Closeness Relation |  |
| Inverse | $\varepsilon\left(\omega^{A B}\right)=\varepsilon_{A}+\varepsilon_{B}$ | $\varepsilon\left(E^{A B}\right)=\varepsilon_{A}+\varepsilon_{B}+1$ |
| Symplectic | $\omega^{B A}=-(-1)^{\varepsilon_{A} \varepsilon_{B} \omega^{A B}}$ | $E^{B A}=-(-1)^{\left(\varepsilon_{A}+1\right)\left(\varepsilon_{B}+1\right)} E^{A B}$ |
| Contravariant | Antisymmetric | Symmetric |
| Tensor | No Laplacian | Odd Laplacian |

## The $2 \times 2=4$ Classical Geometries and their Symmetries

|  | Even Geometry | Odd Geometry |
| :---: | :---: | :---: |
| Riemannian Covariant Metric | $\begin{gathered} \hline \hline g=d z^{A} g_{A B} \vee d z^{B} \\ \varepsilon\left(g_{A B}\right)=\varepsilon_{A}+\varepsilon_{B} \\ g_{B A}=-(-1)^{\left(\varepsilon_{A}+1\right)\left(\varepsilon_{B}+1\right)} g_{A B} \\ \text { Antiskewsymmetric } \\ \text { No Closeness Relation } \end{gathered}$ | $\begin{gathered} g=d z^{A} g_{A B} \vee d z^{B} \\ \varepsilon\left(g_{A B}\right)=\varepsilon_{A}+\varepsilon_{B}+1 \\ g_{B A}=(-1)^{\varepsilon} \varepsilon_{B} g_{A B} \\ \text { Symmetric } \end{gathered}$ <br> No Closeness Relation |
| Inverse Riemannian Contravariant Metric | $\begin{gathered} \varepsilon\left(g^{A B}\right)=\varepsilon_{A}+\varepsilon_{B} \\ g^{B A}=(-1)^{\varepsilon_{A} \varepsilon_{B}} g^{A B} \end{gathered}$ <br> Symmetric <br> Even Laplacian | $\begin{gathered} \varepsilon\left(g^{A B}\right)=\varepsilon_{A}+\varepsilon_{B}+1 \\ g^{B A}=(-1)^{\left(\varepsilon_{A}+1\right)\left(\varepsilon_{B}+1\right)} g^{A B} \\ \text { Skewsymmetric } \\ \text { No Laplacian } \end{gathered}$ |
| Symplectic <br> Covariant <br> Two-Form | $\begin{gathered} \hline \hline \omega=\frac{1}{2} d z^{A} \omega_{A B} \wedge d z^{B} \\ \varepsilon\left(\omega_{A B}\right)=\varepsilon_{A}+\varepsilon_{B} \\ \omega_{B A}=(-1){ }^{\left(\varepsilon_{A}+1\right)\left(\varepsilon_{B}+1\right)} \omega_{A B} \\ \text { Skewsymmetric } \\ \text { Closeness Relation } \end{gathered}$ | $\begin{gathered} \hline E=\frac{1}{2} d z^{A} E_{A B} \wedge d z^{B} \\ \varepsilon\left(E_{A B}\right)=\varepsilon_{A}+\varepsilon_{B}+1 \\ E_{B A}=-(-1)^{\varepsilon} \varepsilon_{B} E_{A B} \\ \text { Antisymmetric } \end{gathered}$ <br> Closeness Relation |
| Inverse Symplectic Contravariant Tensor | $\begin{gathered} \varepsilon\left(\omega^{A B}\right)=\varepsilon_{A}+\varepsilon_{B} \\ \omega^{B A}=-(-1)^{\varepsilon_{A} \varepsilon_{B}}{ }^{A B} \end{gathered}$ <br> Antisymmetric No Laplacian | $\begin{gathered} \varepsilon\left(E^{A B}\right)=\varepsilon_{A}+\varepsilon_{B}+1 \\ E^{B A}=-(-1)^{\left(\varepsilon_{A}+1\right)\left(\varepsilon_{B}+1\right)} E^{A B} \\ \text { Symmetric } \\ \text { Odd Laplacian } \end{gathered}$ |

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| Inverse Riemannian <br> Contravariant Metric | $\begin{gathered} \varepsilon\left(g^{A B}\right)=\varepsilon_{A}+\varepsilon_{B} \\ g^{B A}=(-1)^{\varepsilon_{A} \varepsilon_{B}} g^{A B} \\ \text { Symmetric } \\ \text { Even Laplacian } \end{gathered}$ | $\begin{gathered} \varepsilon\left(g^{A B}\right)=\varepsilon_{A}+\varepsilon_{B}+1 \\ g^{B A}=(-1)^{\left(\varepsilon_{A}+1\right)\left(\varepsilon_{B}+1\right)} g^{A B} \\ \text { Skewsymmetric } \\ \text { No Laplacian } \end{gathered}$ |
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| Inverse <br> Symplectic Contravariant Tensor | $\begin{gathered} \varepsilon\left(\omega^{A B}\right)=\varepsilon_{A}+\varepsilon_{B} \\ \omega^{B A}=-(-1)^{\varepsilon} \varepsilon^{\varepsilon_{B}} \omega^{A B} \\ \text { Antisymmetric } \\ \text { No Laplacian } \end{gathered}$ | $\begin{gathered} \varepsilon\left(E^{A B}\right)=\varepsilon_{A}+\varepsilon_{B}+1 \\ E^{B A}=-(-1)^{\left(\varepsilon_{A}+1\right)\left(\varepsilon_{B}+1\right)} E^{A B} \\ \text { Symmetric } \\ \text { Odd Laplacian } \end{gathered}$ |

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| Covariant | $g_{B A}=-(-1)^{\left(\varepsilon_{A}+1\right)\left(\varepsilon_{B}+1\right)} g_{A B}$ | $g_{B A}=(-1)^{\varepsilon_{A} \varepsilon_{B}} g_{A B}$ |
| Metric | Antiskewsymmetric | Symmetric |
|  | No Closeness Relation | No Closeness Relation |
| Inverse | $\varepsilon\left(g^{A B}\right)=\varepsilon_{A}+\varepsilon_{B}$ | $\varepsilon\left(g^{A B}\right)=\varepsilon_{A}+\varepsilon_{B}+1$ |
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| Contravariant | Symmetric | Skewsymmetric |
| Metric | Even Laplacian | No Laplacian |
|  | $\omega=\frac{1}{2} d z^{A} \omega_{A B} \wedge d z^{B}$ | $E=\frac{1}{2} d z^{A} E_{A B} \wedge d z^{B}$ |
| Symplectic | $\varepsilon\left(\omega_{A B}\right)=\varepsilon_{A}+\varepsilon_{B}$ | $\varepsilon\left(E_{A B}\right)=\varepsilon_{A}+\varepsilon_{B}+1$ |
| Covariant | $\omega_{B A}=(-1)^{\left(\varepsilon_{A}+1\right)\left(\varepsilon_{B}+1\right)} \omega_{A B}$ | $E_{B A}=-(-1)^{\varepsilon_{A} \varepsilon_{B} E_{A B}}$ |
| Two-Form | Skewsymmetric | Antisymmetric |
|  | Closeness Relation | Closeness Relation |
| Inverse | $\varepsilon\left(\omega^{A B}\right)=\varepsilon_{A}+\varepsilon_{B}$ | $\varepsilon\left(E^{A B}\right)=\varepsilon_{A}+\varepsilon_{B}+1$ |
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| Contravariant | Antisymmetric | Symmetric |
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## The Odd Scalar $\nu_{\rho}$

Odd Scalar in Antisymplectic Geometry (KB 2006)

$$
\nu_{\rho}:=\nu_{\rho}^{(0)}+\frac{\nu^{(1)}}{8}-\frac{\nu^{(2)}}{24}
$$

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## Terms built from $E$ and $\rho$

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\nu^{(2)} & :=-(-1)^{\varepsilon_{B}}\left(\frac{\partial^{\ell}}{\partial z^{A}} E_{B C}\right)\left(z^{C},\left(z^{B}, z^{A}\right)\right)
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& =(-1)^{\varepsilon_{A} \varepsilon_{D}}\left(\frac{\partial^{\ell}}{\partial z^{D}} E^{A B}\right) E_{B C}\left(E^{C D} \frac{\overleftarrow{\partial^{r}}}{\partial z^{A}}\right)
\end{aligned}
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## Classification of 2nd-order Differential Invariants

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- $\nu$ scales as $\nu \rightarrow \lambda \nu$ under constant Weyl scaling $E^{A B} \rightarrow \lambda E^{A B}$, where $\lambda$ is a $z$-independent parameter,


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- $\nu$ scales as $\nu \rightarrow \lambda \nu$ under constant Weyl scaling $E^{A B} \rightarrow \lambda E^{A B}$, where $\lambda$ is a z-independent parameter,
- and each term in $\nu$ contains precisely two $z$-derivatives?


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- and each term in $\nu$ contains precisely two $z$-derivatives?


## Unique Answer (up to scaling)

$$
\nu=\alpha \nu_{\rho}
$$

## The $\triangle$ Operator

## Question <br> What is (the local form of) the most general differential operator $\Delta$

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- that takes scalar functions to scalar functions,
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- that is of 2 nd-order,
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## Unique Answer (modulo an odd constant) (Batalin,KB 2007)

$$
\Delta=\Delta_{\rho}+\nu_{\rho}
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## The $\triangle$ Operator

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What is (the local form of) the most general differential operator $\Delta$

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$$
\Delta=\Delta_{\rho}+\nu_{\rho}
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The $\Delta$ Operator $=$ Odd Laplacian + Odd Scalar.

## Quantum Master Equation

## Exponential form

$$
\Delta e^{\frac{i}{\hbar} W}=0
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## Additive form

$\frac{1}{2}(W, W)=i \hbar \Delta_{\rho} W+\hbar^{2} \nu_{\rho}$

## Quantum Master Equation

Exponential form

$$
\Delta e^{\frac{i}{\hbar} W}=0
$$

## Additive form

$$
\begin{aligned}
& \qquad \frac{1}{2}(W, W)=i \hbar \Delta_{\rho} W+\hbar^{2} \nu_{\rho} \\
& \text { Odd scalar } \nu_{\rho} \text { enters at } \\
& \text { 2-loop. }
\end{aligned}
$$

## Quantum Master Equation

Exponential form

$$
\Delta e^{\frac{i}{\hbar} W}=0
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## Quantum Master

 Action$$
W=S+\sum_{n=1}^{\infty} \hbar^{n} M_{n}
$$

## Additive form

$$
\frac{1}{2}(W, W)=i \hbar \Delta_{\rho} W+\hbar^{2} \nu_{\rho}
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Odd scalar $\nu_{\rho}$ enters at 2-loop.

## Quantum Master Equation

Exponential form

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Quantum Master Action

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## Infinite Tower of Master Equations

$$
(S, S)=0
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Odd scalar $\nu_{\rho}$ enters at 2-loop.

Infinite Tower of Master Equations

$$
(S, S)=0 \leftarrow \text { Classical Master Equation }
$$

## Quantum Master Equation

Exponential form

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Quantum Master Action

$$
W=S+\sum_{n=1}^{\infty} \hbar^{n} M_{n}
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## Additive form

$$
\frac{1}{2}(W, W)=i \hbar \Delta_{\rho} W+\hbar^{2} \nu_{\rho}
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Odd scalar $\nu_{\rho}$ enters at 2-loop.

Infinite Tower of Master Equations

$$
\begin{aligned}
(S, S) & =0 \quad \leftarrow \quad \text { Classical Master Equation } \\
\left(M_{1}, S\right) & =i\left(\Delta_{\rho} S\right)
\end{aligned}
$$

## Quantum Master Equation

Exponential form

$$
\Delta e^{\frac{i}{\hbar} W}=0
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Quantum Master Action

$$
W=S+\sum_{n=1}^{\infty} \hbar^{n} M_{n}
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(S, S) & =0 \\
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\left(M_{2}, S\right) & =i\left(\Delta_{\rho} M_{1}\right)-\frac{1}{2}\left(M_{1}, M_{1}\right)+\nu_{\rho}
\end{aligned}
$$

## Quantum Master Equation

Exponential form

$$
\Delta e^{\frac{i}{\hbar} W}=0
$$

Quantum Master Action

$$
W=S+\sum_{n=1}^{\infty} \hbar^{n} M_{n}
$$

## Additive form

$$
\frac{1}{2}(W, W)=i \hbar \Delta_{\rho} W+\hbar^{2} \nu_{\rho}
$$

Odd scalar $\nu_{\rho}$ enters at 2-loop.

Infinite Tower of Master Equations

$$
\begin{aligned}
(S, S) & =0 \\
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\forall n \geq 3: \quad\left(M_{n}, S\right) & =i\left(\Delta_{\rho} M_{n-1}\right)-\frac{1}{2} \sum_{r=1}^{n-1}\left(M_{r}, M_{n-r}\right)
\end{aligned}
$$

## Khudaverdian's $\triangle_{E}$ Operator

## The $\triangle$ Operator

$$
\Delta=\Delta_{\rho}+\nu_{\rho}
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$=$ Odd Laplacian + Odd Scalar
$=$ built from $E$ and $\rho$.

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\sqrt{\rho} \Delta \frac{1}{\sqrt{\rho}}
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is independent of
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\Delta_{E}:=\Delta_{1}=(-1)^{\varepsilon_{\alpha}} \frac{\overrightarrow{\partial^{\ell}}}{\partial \phi^{\alpha}} \frac{\overrightarrow{\partial^{\ell}}}{\partial \phi_{\alpha}^{*}}
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- In General Coordinates:

$$
\begin{array}{r}
\Delta_{E}:=\Delta_{1}+\frac{\nu^{(1)}}{8}-\frac{\nu^{(2)}}{24} \\
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## A Comparative Study of Laplacians in Riemannian and Antisymplectic Geometry

(1) Anti-Poisson Geometry
(2) Riemannian Geometry
(3) Conclusions

## The Even Scalar $\nu_{\rho}$

Even Scalar in Riemannian Geometry with density $\rho$

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\nu_{\rho}:=\nu_{\rho}^{(0)}+\frac{\nu^{(1)}}{4}-\frac{\nu^{(2)}}{8}-\frac{\nu^{(3)}}{16}
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$$
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$$

$$
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$$

$$
\nu^{(3)}:=(-1)^{\varepsilon_{A}}\left(g_{A B}, g^{B A}\right) \leftarrow \text { bracket wrt. metric } \mathrm{g} \text {. }
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## Interpretation of $\nu_{\rho}$ in terms of Scalar Curvature $R$

## Riemannian Case

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\nu_{\rho_{g}}=-\frac{R}{4}
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## Even Scalar Curvature

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## for the Levi-Civita

Connection $\nabla$, i.e., $\nabla$ is:

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Riemannian Case /w General $\rho$

$$
\nu_{\rho}=\sqrt{\frac{\rho_{g}}{\rho}}\left(\Delta_{\rho_{g}} \sqrt{\frac{\rho}{\rho_{g}}}\right)-\frac{R}{4}
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## Complete Solution

$$
\nu=\alpha \nu_{\rho}+\beta \nu_{\rho_{g}}+\gamma\left(\ln \frac{\rho}{\rho_{g}}, \ln \frac{\rho}{\rho_{g}}\right)
$$

## The Even $\triangle$ Operator

| Even $\Delta$ |  | Even |  | Even |
| :---: | :---: | :---: | :---: | :---: |
| Operator |  | Laplacian |  | Scalar |
| $\Delta$ | $:=$ | $\Delta_{\rho}$ | + | $\nu_{\rho}$ |

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| Operator |  | Laplacian |  | Scalar |  |
| $\Delta$ | := | $\Delta_{\rho}$ | $+$ | $\nu_{\rho}$ |  |
|  |  | $\downarrow$ |  | $\downarrow$ | for $\quad \rho \rightarrow \rho_{g}$ |
|  |  | $\Delta_{\rho_{g}}$ | - | $\frac{R}{4}$ |  |
|  |  | Laplace- |  | a quarter |  |
|  |  | Beltrami |  | of the Scalar |  |
|  |  | Operator |  | Curvature |  |

## The Even $\triangle$ Operator



For comparison: Conformally Covariant Laplacian

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\Delta_{\rho_{g}}-\frac{(N-2) R}{(N-1) 4} \quad \rightarrow \quad \Delta_{\rho_{g}}-\frac{R}{4} \quad \text { for } \quad N \rightarrow \infty
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## Particle in Curved Space

## Classical Hamiltonian Action

$$
\begin{aligned}
S_{\mathrm{cl}} & =\int_{1} d t\left(p_{A} \dot{z}^{A}-H_{\mathrm{cl}}\right) \\
H_{\mathrm{cl}} & =\frac{1}{2} p_{A} p_{B} g^{B A} \\
\left\{z^{A}, p_{B}\right\}_{P B} & =\delta_{B}^{A}
\end{aligned}
$$

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\begin{aligned}
S_{\mathrm{cl}} & =\int d t\left(p_{A} \dot{z}^{A}-H_{\mathrm{cl}}\right) \\
H_{\mathrm{cl}} & =\frac{1}{2} p_{A} p_{B} g^{B A} \\
\left\{z^{A}, p_{B}\right\}_{P B} & =\delta_{B}^{A}
\end{aligned}
$$

## Naive Quantum Hamiltonian

$$
\hat{H}_{\rho}=\frac{1}{2 \sqrt{\rho(\hat{z})}} \hat{p}_{A} \rho(\hat{z}) g^{A B}(\hat{z}) \hat{p}_{B} \frac{(-1)^{\varepsilon_{B}}}{\sqrt{\rho(\hat{z})}}
$$

## Particle in Curved Space

## Classical Hamiltonian Action

$$
\begin{aligned}
S_{\mathrm{cl}} & =\int d t\left(p_{A} \dot{z}^{A}-H_{\mathrm{cl}}\right) \\
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## Full Quantum Hamiltonian

$$
\hat{H}=\hat{H}_{\rho}-\frac{\hbar^{2}}{2} \nu_{\rho}(\hat{z}) \sim T\left(H_{\mathrm{cl}}\right)
$$

## Particle in Curved Space

Classical Hamiltonian Action

$$
\begin{aligned}
S_{\mathrm{cl}} & =\int_{H_{\mathrm{cl}}} d t\left(p_{A} \dot{z}^{A}-H_{\mathrm{cl}}\right) \\
& =\frac{1}{2} p_{A} p_{B} g^{B A} \\
\left\{z^{A}, p_{B}\right\}_{P B} & =\delta_{B}^{A}
\end{aligned}
$$

## Schrödinger Representation

$\frac{\hbar}{i} \frac{\overrightarrow{\partial^{\ell}}}{\partial z^{A}}=\sqrt{\rho(\hat{z})} \hat{p}_{A} \frac{(-1)^{\varepsilon_{A}}}{\sqrt{\rho(\hat{z})}}$

Naive Quantum Hamiltonian

$$
\hat{H}_{\rho}=\frac{1}{2 \sqrt{\rho(\hat{z})}} \hat{p}_{A} \rho(\hat{z}) g^{A B}(\hat{z}) \hat{p}_{B} \frac{(-1)^{\varepsilon_{B}}}{\sqrt{\rho(\hat{z})}}
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\hat{H}=\hat{H}_{\rho}-\frac{\hbar^{2}}{2} \nu_{\rho}(\hat{z}) \sim T\left(H_{\mathrm{cl}}\right)
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Classical Hamiltonian Action

$$
\begin{aligned}
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\left\{z^{A}, p_{B}\right\}_{P B} & =\delta_{B}^{A}
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$$

## Laplacian

Full Quantum Hamiltonian

$$
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## Particle in Curved Space

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Full Quantum Hamiltonian

$$
\hat{H}=\hat{H}_{\rho}-\frac{\hbar^{2}}{2} \nu_{\rho}(\hat{z}) \sim T\left(H_{\mathrm{cl}}\right)
$$

$$
\Delta=\Delta_{\rho}+\nu_{\rho}
$$

## Operator Formalism $\leftrightarrow$ Path Integral Formalism

## (starting with DeWitt 1957)

## The operator formalism

with the full Hamiltonian operator $\hat{H}$
$\downarrow$ corresponds to $\downarrow$
a Hamiltonian path integral formulation where the path integral action is the pure classical action $S_{\mathrm{cl}}$ with no quantum corrections.

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$$
\left\langle z_{f}\right| \exp \left[-\frac{i}{\hbar} \hat{H} \Delta t\right]\left|z_{i}\right\rangle \sim \int_{z\left(t_{i}\right)=z_{i}}^{z\left(t_{f}\right)=z_{f}}[d z][d p] \exp \left[\frac{i}{\hbar} S_{\mathrm{cl}}[z, p]\right]
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Full Quantum
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## Classical Action

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S_{\mathrm{cl}}[z, p]=\int d t\left(p_{A} \dot{z}^{A}-H_{\mathrm{cl}}\right)
$$

## A Comparative Study of Laplacians in Riemannian and Antisymplectic Geometry

(1) Anti-Poisson Geometry
(2) Riemannian Geometry
(3) Conclusions

## Conclusions

## Odd $\triangle$ Operator in Antisymplectic Geometry

$$
2 \Delta=2 \Delta_{\rho}+2 \nu_{\rho}=2 \Delta_{\rho}-\frac{R}{4}
$$

## Conclusions

Odd $\triangle$ Operator in
Antisymplectic Geometry

$$
\begin{aligned}
& 2 \Delta=2 \Delta_{\rho}+2 \nu_{\rho}=2 \Delta_{\rho}-\frac{R}{4} \\
& \text { Characterized by } \\
& \text { nilpotency }
\end{aligned}
$$

## Conclusions

Odd $\triangle$ Operator in
Antisymplectic Geometry

$$
2 \Delta=2 \Delta_{\rho}+2 \nu_{\rho}=2 \Delta_{\rho}-\frac{R}{4}
$$

- Characterized by nilpotency
- and characterized by a $\rho$ independence argument.


## Conclusions

## Even $\triangle$ Operator in Riemannian Geometry

$\Delta=\Delta_{\rho}+\nu_{\rho} \rightarrow \Delta_{\rho_{g}}-\frac{R}{4}$

Odd $\triangle$ Operator in
Antisymplectic Geometry

$$
2 \Delta=2 \Delta_{\rho}+2 \nu_{\rho}=2 \Delta_{\rho}-\frac{R}{4}
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## Conclusions

Even $\triangle$ Operator in
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## Particle in Curved Space

## $\Delta$ is the full quantum

 Hamiltonian$$
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in the Schrödinger representation.

## Conclusions

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Riemannian Geometry

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2 \Delta=2 \Delta_{\rho}+2 \nu_{\rho}=2 \Delta_{\rho}-\frac{R}{4}
$$

- Characterized by nilpotency
- and characterized by a $\rho$ independence argument.

Curvature term in
Quantum Master Equation

$$
(W, W)=2 i \hbar \Delta_{\rho} W-\hbar^{2} \frac{R}{4}
$$

Important 2-loop effect.

## References

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