Superfluidity

in

ultracold Fermi gases

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- Introduction
- Overview of theory and experiment
- Polarized ultracold Fermi gases
- 1 spin \downarrow + N spins \uparrow

Review paper : S. Giorgini, L. P. Pitaevskii and S. Stringari, Rev. Mod. Phys. **80**, 1215 (2008)



- Ultracold Fermi gases, some basics
- Interatomic distance ~ 10^2 10^3 nm
- In practice alkali with even number of nucleons: ${}^{6}Li$ or ${}^{40}K$
- -Trapping (magnetic or optical)

 \rightarrow parabolic potential (harmonic oscillator) 100μ \rightarrow inhomogeneous system $N \sim 10^6 - 10^7$

- Ultracold : optical + evaporative cooling (temperature ~ 1nK 1μK)
 → degenerate Fermi gas : quantum regime
- Very low $T \rightarrow$ very low energy \rightarrow **s-wave** scattering only
- But prohibited for **fermions** by **Pauli exclusion** → **no interaction** !

(very good for atomic cloks)

⇒ Physics usually with two fermionic species : "↑ and ↓ " usually two lowest energy hyperfine states of same element Ex: ⁶Li



- Low energy s-wave scattering ⇒ single parameter scattering length a
- Short-range interaction compared to interatomic distance

 $V(\mathbf{r}) = (4\pi h^2 a/m) \,\delta(\mathbf{r})$

Very convenient ! \Rightarrow

Remarkable model system

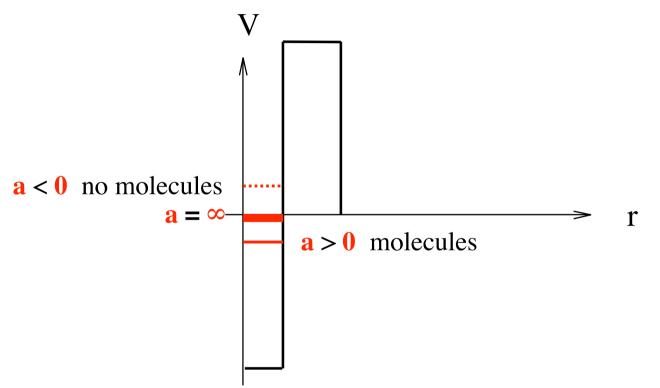
for strongly interacting (and strongly correlated) Fermi systems normal and superfluid

- BCS Superfluidity
 - Formation of Cooper pairs between ↑ and ↓ atoms
 requires attractive interaction + degenerate Fermi gas
 - Attraction \rightarrow **a** < **0**



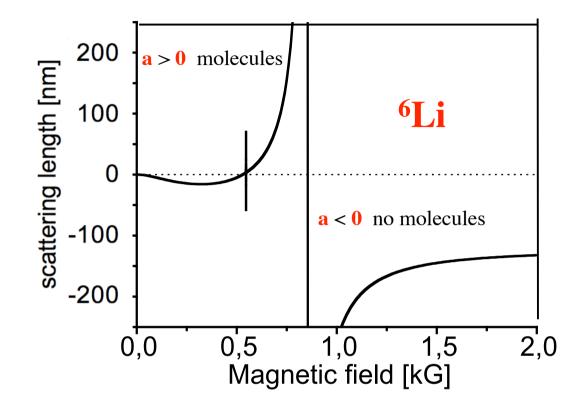
<u>Feshbach resonance</u>

- Allows to **control effective interaction** via **magnetic field** by changing **scattering length a**
- Scattering length $a = \infty$ if bound state with energy = 0 exists



Actually atoms very near each other are not in the same spin configuration as when they are very far from each other ("closed channel" and "open channel")
 → sensitivity of bound state energy to magnetic field

• Feshbach resonance for two ⁶Li particles in vacuum



• Interaction **tunable** at will by **magnetic field** !

Dream !

⇒ Allows a **physical realization** of the **BEC-BCS crossover**

"New" superfluid !



• **BEC - BCS crossover**

- BCS Ansatz for ground state wavefunction:

$$\Psi(\mathbf{r}_1,\mathbf{r}_2,\ldots) = A \left\{ \Phi(\mathbf{r}_1 - \mathbf{r}_2) \Phi(\mathbf{r}_3 - \mathbf{r}_4) \ldots \right\}$$

describes as well **dilute gas** of **molecules**, made of 2 fermions

- Known since Popov (66), Keldysh and Kozlov(68), Eagles(69).... Leggett (80), Nozières and Schmitt-Rink (85)
- Accurate in **weak-coupling**(BCS) limit and **strong-coupling**(BEC) limit In between : physically quite reasonable "**interpolation scheme**"
- \Rightarrow in between need **experiment** to tell what happens!

No exact theory



<u>Theoretical interest</u>

- General interest : **strongly interacting superfluid** systems (**unitarity**)
- Interesting for high T_c superconductivity : very tight pairs, pseudogap ~ preformed pairs ? ⇒ nearly BEC ?
- Ultracold gases results not in favor of this model:
 BEC physics pushed beyond unitarity by Fermi sea
- Need to control normal state for good understanding of superfluid

 $\epsilon_k = k^2/2m$ $\xi_k = \epsilon_k - \mu$

 $\frac{m}{4\pi a} = \sum_{k} \left(\frac{1}{2\epsilon_k} - \frac{1}{2E_k} \right)$

 $n \equiv \frac{k_F^3}{6\pi^2} = \sum_k \frac{1}{2} \left(1 - \frac{\xi_k}{E_k} \right)$

 $E_k = \sqrt{\xi_k^2 + \Delta^2}$

 $E_F \equiv \frac{k_F^2}{2m}$

- Theoretical treatments
 - BCS (mean field) or equivalent
 Use scattering length a known experimentally
 instead of BCS potential V (cf. Galitski and Belaiev)
 single parameter 1/k_Ea
 - fairly reasonable !

- Quantum Monte-Carlo quite reliable, but numbers
- T-matrix

Quite often used (High T_c) Quite satisfactory for ultracold gases

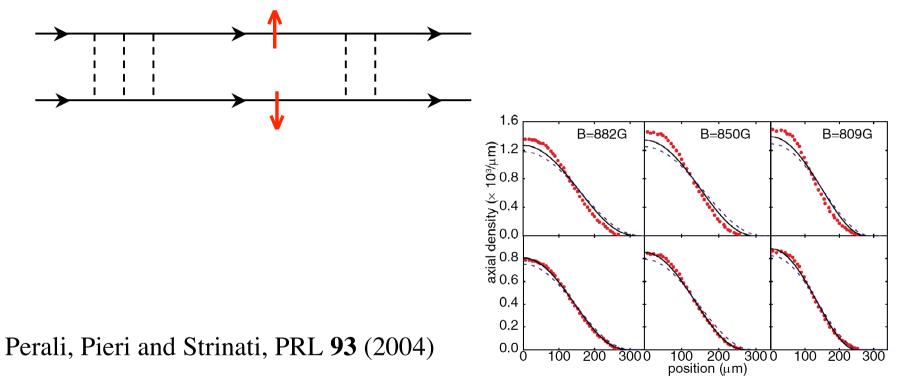


FIG. 1 (color online). Comparison between experimental and theoretical axial density profiles. Experimental data from Ref. [17] (dots) are shown for three different values of the magnetic field *B* tuning the FF resonance. Theoretical results at T = 0 obtained by our theory (solid lines) and by BCS mean field (dashed lines) are shown for the corresponding couplings $(k_F a_F)^{-1}$ given in the text. The upper (lower) panel refers to the estimated number of atoms $N = 4 \times 10^5$ ($N = 2.3 \times 10^5$).

• Unitarity

- Quite interesting case: very strongly interacting
- parameter : $1/k_F a = 0$
- $\Rightarrow single parameter k_F (or E_F) left \Rightarrow dimensional analysis$ Example :

$$\mu = \xi E_{\rm F}$$

BCS $\xi = 0.59$ QMC $\xi = 0.42-0.44$ T-matrix $\xi = 0.455$ Exp $\xi = 0.27-0.51(\pm 10)$



Away from unitarity: chemical potential

Pieri, Pisani and Strinati, PRB 72 (2005)

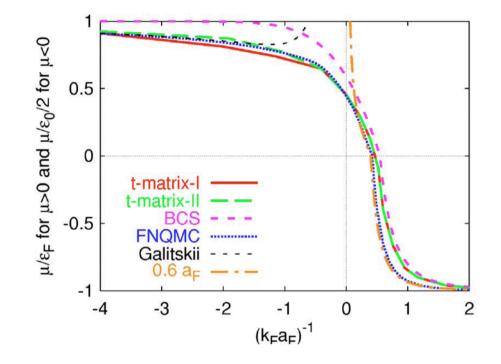


FIG. 1. (Color online) Chemical potential at zero temperature vs the coupling parameter $(k_F a_F)^{-1}$. The results of the present theory (*t*-matrix-I) and of its version without the inclusion of the selfenergy shift Σ_0 (*t*-matrix-II) are compared with the BCS mean field (BCS), the fixed-node QMC data from Ref. 11 (FNQMC), the Galitskii's expression for the dilute Fermi gas (Galitskii), and the asymptotic expression for strong coupling using the result a_B =0.6 a_F .



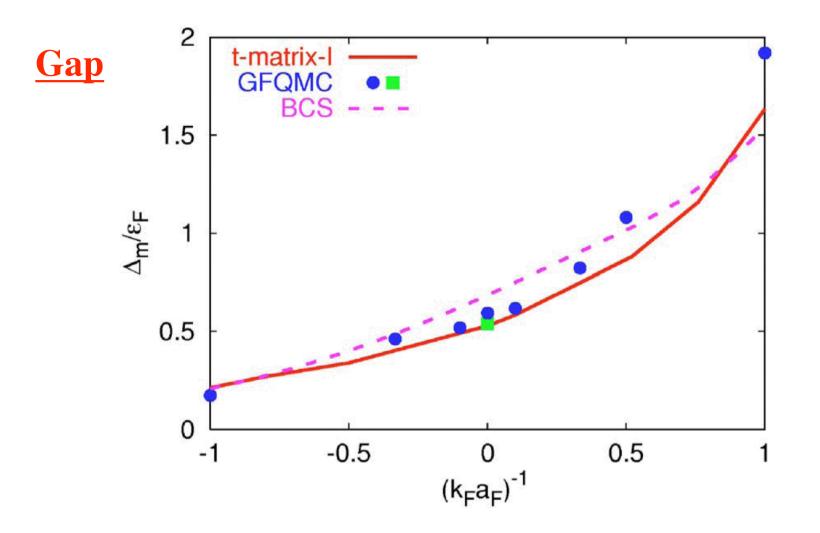
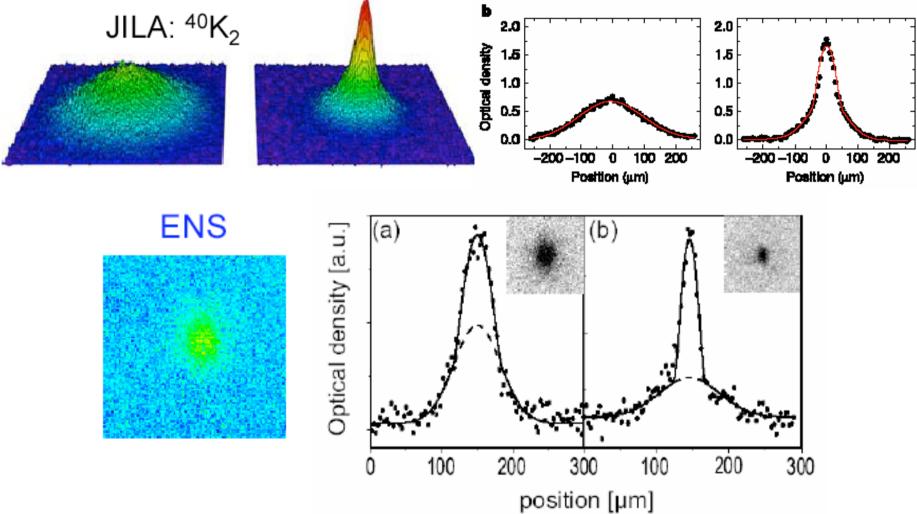


FIG. 2. (Color online) Excitation gap $\Delta_{\rm m}$ at zero temperature vs the coupling parameter $(k_F a_F)^{-1}$. The results of the present theory (*t*-matrix-I) are compared with the Green's function QMC data of Refs. 8 and 10 (GFQMC) as well as with the BCS mean field (BCS).

• **Bose-Einstein condensates of molecules** (2003-2004)



- Also **MIT** and **Innsbrück**
- First Bose-Einstein condensates of molecules made of fermions !

- Vortices as evidence for superfluidity (2005)
- Before : Anisotropic expansion ~ BEC
 Collective mode damping (much more convincing)

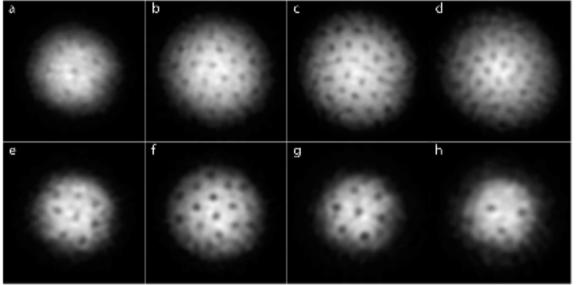
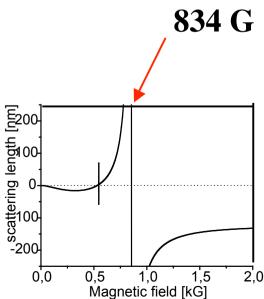


Fig. 2: Vortices in a strongly interacting gas of fermionic atoms on the BEC- and the BCS-side of the Feshbach resonance. At the given field, the cloud of lithium atoms was stirred for 300 ms (a) to 500 ms (b-h) followed by an equilibration time of 500 ms. After 2 ms of ballistic expansion, the magnetic field was ramped to 735 G for imaging (see text for details). The magnetic fields were (a) 740 G, (b) 766 G, (c) 792 G, (d) 812 G, (e) 833 G, (f) 843 G, (g) 853 G and (h) 863 G. The field of view of each image is 880 m × 880 m.

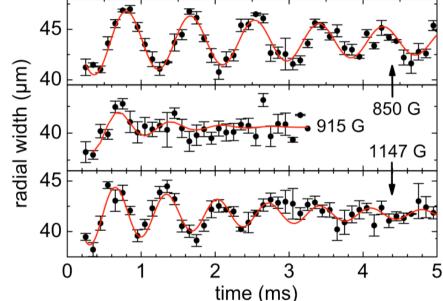


MIT

• Collective oscillations in harmonic trap (superfluid state)

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Cigar geometry \omega_z \ll \omega_{x,y}
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- In situ experiments (no need for interpretation)
- High experimental precision possible

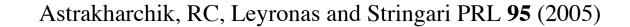


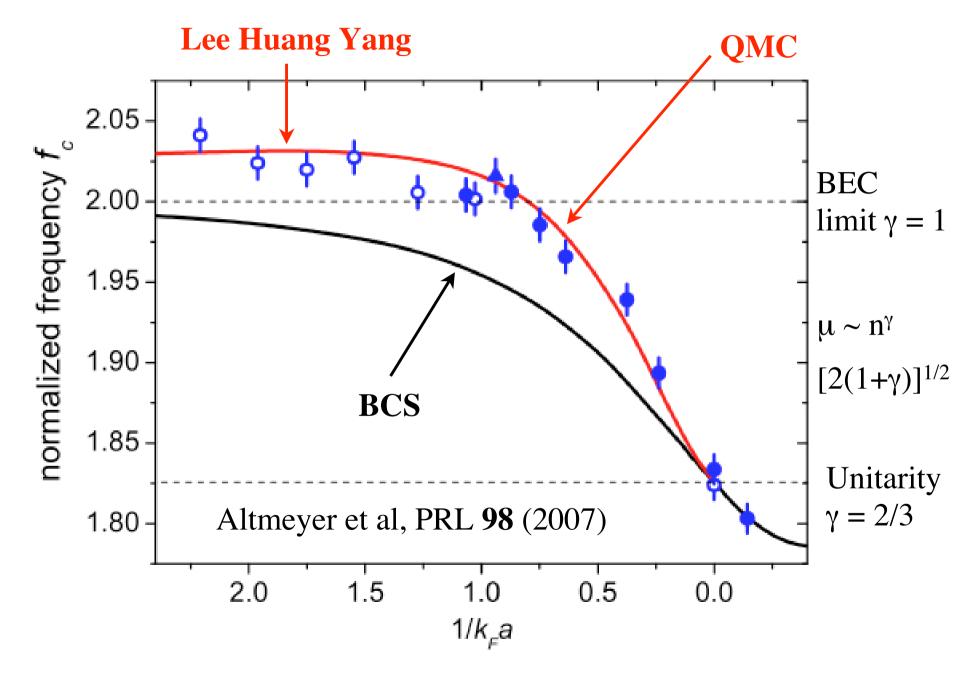
Most experiments require expansion

- Direct access to equation of state $\mu(n)$



- Equations of state
- Monte-Carlo : should be reasonably accurate
- 0.42 0.68 1.03 1.56 2.57 5.80).2 0.6 - BCS equation of state 0.59 1/k₋a 0.5 0.44 0.4 μ / Ε_F 0.3 Mean-Field-BCS - Hydrodynamics $a_{\rm M} = 0.6 a$ $a_{\rm M} = 2. a$ 0.2 Monte-Carlo (PSS) 0.1 - $T \rightarrow 0$ Superfluid satisfies 0 0.2 0.4 1.2 1.4 0 0.6 0.8 1.6 hydrodynamics, but $\omega \ll E_{\rm h}$ atan(1/k_a) 2.2 (**pair binding** energy, i.e. collisionless limit **BEC** limit Ţ 2.0 $\omega \ll \Delta$ on **BCS side**) $\Omega_r \, / \omega_r$ 1.8 1.6 - Radial geometry 1.4 0.8 0.6 - Strong attenuation at 910 G ັ¹ 0.6 ສ/ັ່ນ ມີ 0.4 \rightarrow pair-breaking peak $\Omega_r = 2 \Delta (T,B)$ 0.2 \rightarrow superfluid ! 0.0 1000 1200 600 800 magnetic field (G)

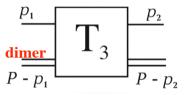


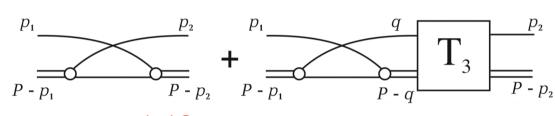


- Equation of state for composite bosons
 - Purely fermionic exact theory same spirit as **Keldysh and Kozlov,** Sov. Phys. JETP **23** (1968)
- **Diagrammatic formulation of 4-body problem**

Brodsky, Klaptsov, Kagan, RC and Leyronas, JETP Letters 82 (2005)

- 3-body



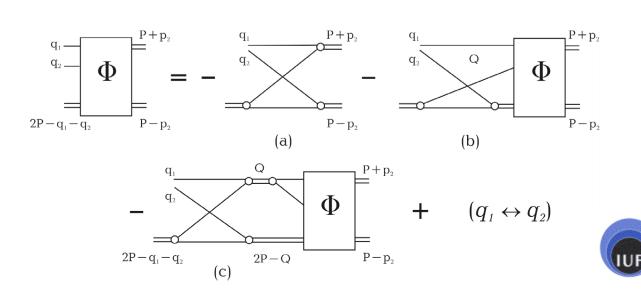


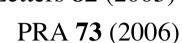
leads to

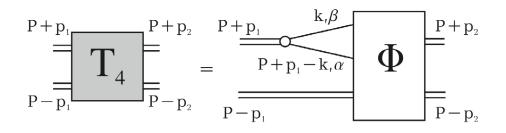
 $a_3 = 1.18 a$

Skorniakov and Ter-Martirosian, Sov. Phys. JETP 4 (1957)

- **4-body**



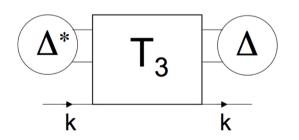




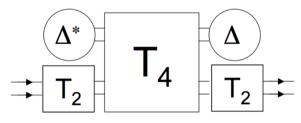
leads to $a_{\rm M} = 0.60 a$

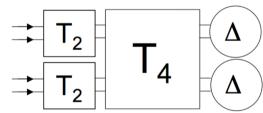
Petrov, Salomon and Shlyapnikov, PRL 93 (2004) (solving Schrödinger equation)

- Many-body (superfluid T=0)
- Systematic expansion in powers of anomalous self-energy $\Delta(k)$
- Collective mode contributions



Lowest order normal self-energy





Irreducible vertices for collective mode propagator

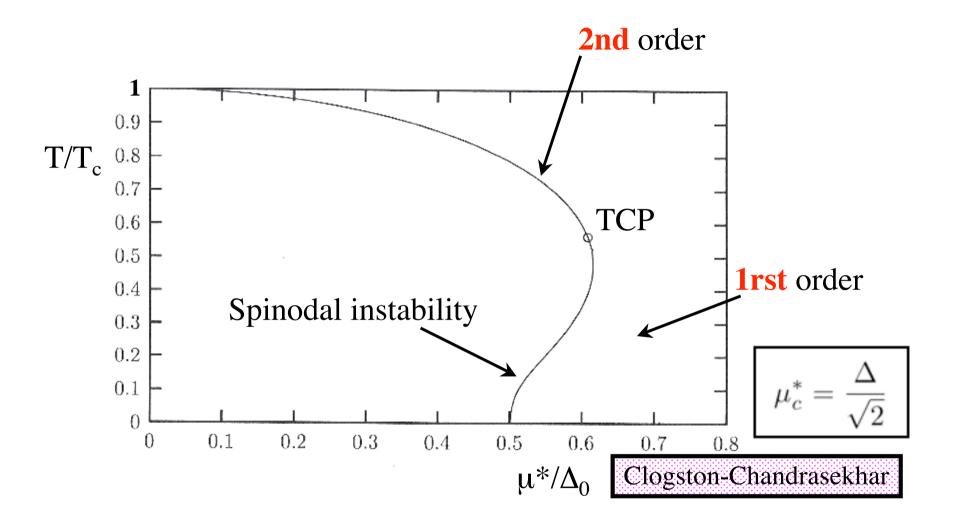
$$\cdots \qquad \Rightarrow \qquad \left[\mu = -\frac{E_b}{2} + \frac{\pi a_M}{m} n \left[1 + \frac{32}{3\sqrt{\pi}} (n a_M^3)^{1/2} \right] \right] \quad LHY$$



- <u>New physical situations</u>
 - Unequal masses : $m_{\uparrow} \neq m_{\downarrow}$ experiments just beginning
 - Bose Fermi mixtures
- <u>" Polarized " gases</u>
 - Strong imbalance $n_{\uparrow} \neq n_{\downarrow}$ easily achieved (stable) $\Rightarrow \mu_{\uparrow} \neq \mu_{\downarrow}$
 - Weak imbalance in superconductors in magnetic fieldby coupling to electronic spinE = -M.B(with orbital currents suppressed in planar geometry)
 - Very interesting for high critical field ($High T_c !$)
 - Very interesting for **quark matter** :

superfluid core of neutron stars, heavy ions collisions Casalbuoni and Nardulli Rev.Mod.Phys. 76 (2004)

- Imbalance breaks pairs since pairing \Rightarrow $n_{\uparrow} = n_{\downarrow}$ \Rightarrow critical "field" μ_{c}^{*} $\mu^{*} = (\mu_{\uparrow} - \mu_{\downarrow})/2$ • Weak coupling BCS





• Fulde - Ferrell and Larkin - Ovchinnikov (1964)

- FFLO or LOFF phases
- Pairs with nonzero total momentum $q \neq 0$

better in high effective field \Rightarrow

extension of superfluid stability domain

- spontaneous symmetry breaking ! (~vorticity)
- No very clear observation in standard superconductors
- Larkin Ovchinnikov :
- T = 0 Ginzburg-Landau investigation : 2nd order transition
- Best LO solution :

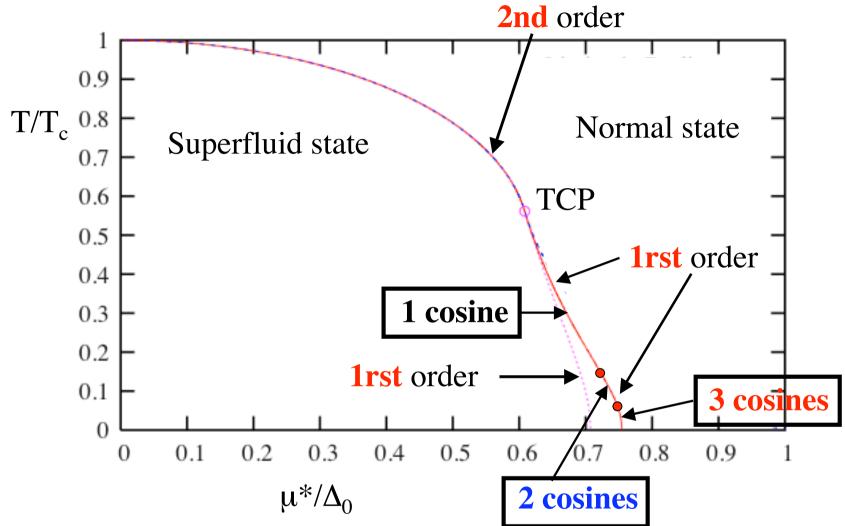
$$\Delta(\mathbf{r}) \sim \cos(\mathbf{q}.\mathbf{r})$$

• Reinvestigated for 1rst order transition

better ! but transition line very near LO

$$\Delta(\mathbf{r}) = \sum_{i=x,y,z} \cos(\mathbf{q}_i \cdot \mathbf{r}) \begin{bmatrix} T=0 \\ I = 0 \end{bmatrix}$$

RC and C. Mora, Europhys. Lett. 68 (2004)



• No LOFF state seen (yet?) in ultracold gases:

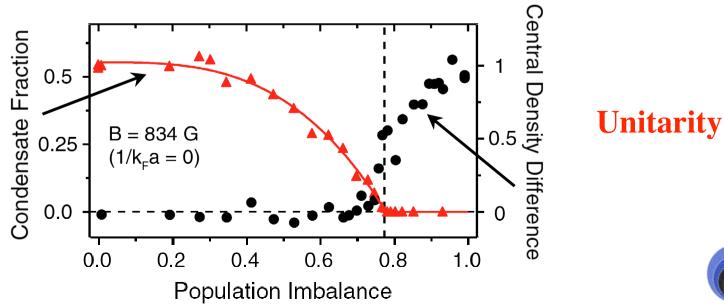
only for **weakly interacting** systems ??

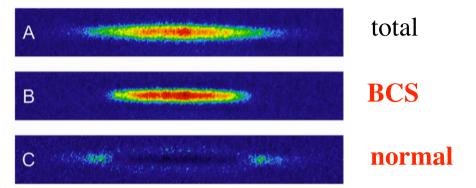
- Difficult: inhomogeneous systems
- Mostly ignored by theories
- 2D or 1D better



• Experiments on polarized ultracold Fermi gases

- Exp. systems inhomogeneous !
- Partridge et al. (Rice) Science **311** (2006)
 Zwierlein et al. (MIT) Science **311** (2006)
 Shin et al. (MIT) PRL **97** (2006)
- Phase separation seen between
- BCS phase $\mathbf{n}_{\uparrow} = \mathbf{n}_{\downarrow}$ and (strongly polarized) normal phase $\mathbf{n}_{\uparrow} \neq \mathbf{n}_{\downarrow}$
- Disagreement between Rice and MIT : Rice smaller and more elongated ?







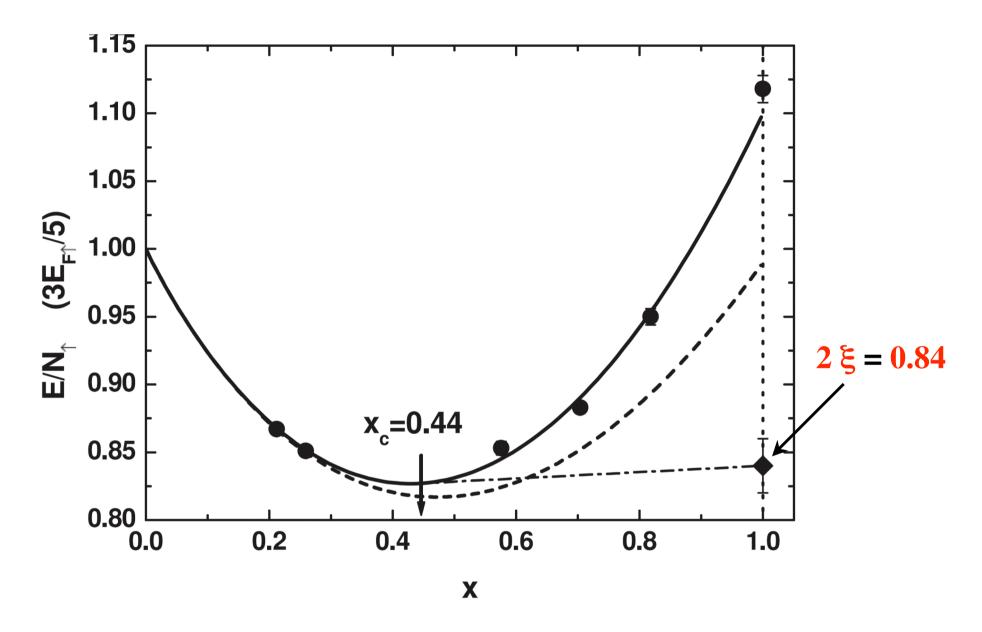
- Many theoretical papers !
- C. Lobo, A. Recati, S. Giorgini and S. Stringari, PRL 97 (2006)
- Generalizes Clogston-Chandrasekhar for strongly interacting system
- Strongly polarized normal state at unitarity (for simplicity)
 - single \downarrow spin with (non interacting) Fermi sea of \uparrow spins $e_{\downarrow}(k) = \mu_{\downarrow} + \frac{k^2}{2m^*}$
- small Fermi sea of \downarrow spin with (non interacting) Fermi sea of \uparrow spins

$$E_{\downarrow} = \mu_{\downarrow} n_{\downarrow} + \frac{3}{5} \frac{k_{F\downarrow}^2}{2m^*} n_{\downarrow} \qquad x = \frac{n_{\downarrow}}{n_{\uparrow}}$$
$$\frac{E(x)}{n_{\uparrow}} = \frac{3}{5} E_F \left(1 - \frac{5}{3} \frac{|\mu_{\downarrow}|}{E_F} x + \frac{m}{m^*} x^{5/3} \right)$$

- QMC: $\frac{|\mu_{\downarrow}|}{E_F} = 0.58$ $\frac{m}{m^*} = 1/1.04$ Superfluid: x = 1 $\frac{E}{n_*} = 2\xi \frac{3}{5}E_F$

$$\overline{n_{\uparrow}}$$
 =





First order transition for $x_c = 0.44$



• **Trapped gas** with **Local Density Approximation** (LDA)

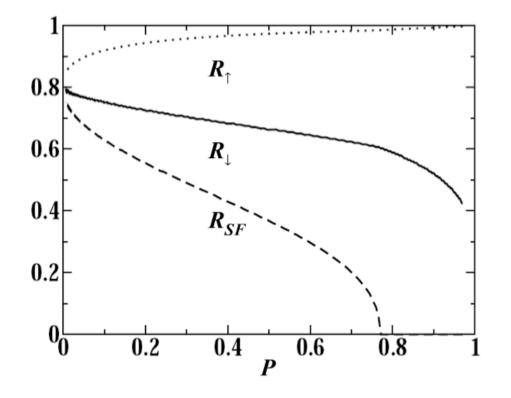
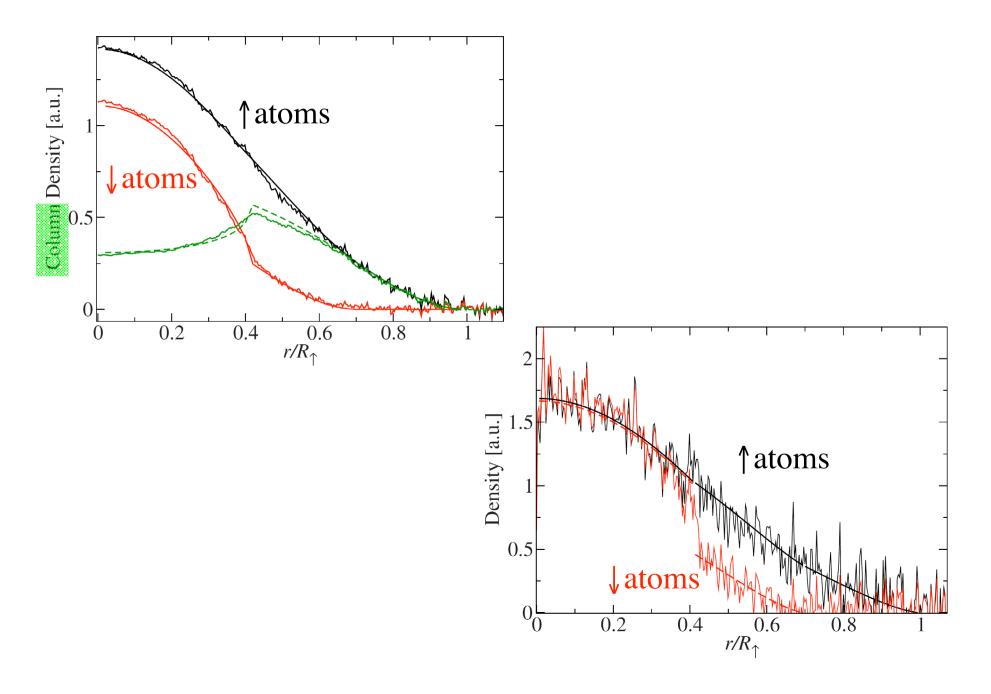


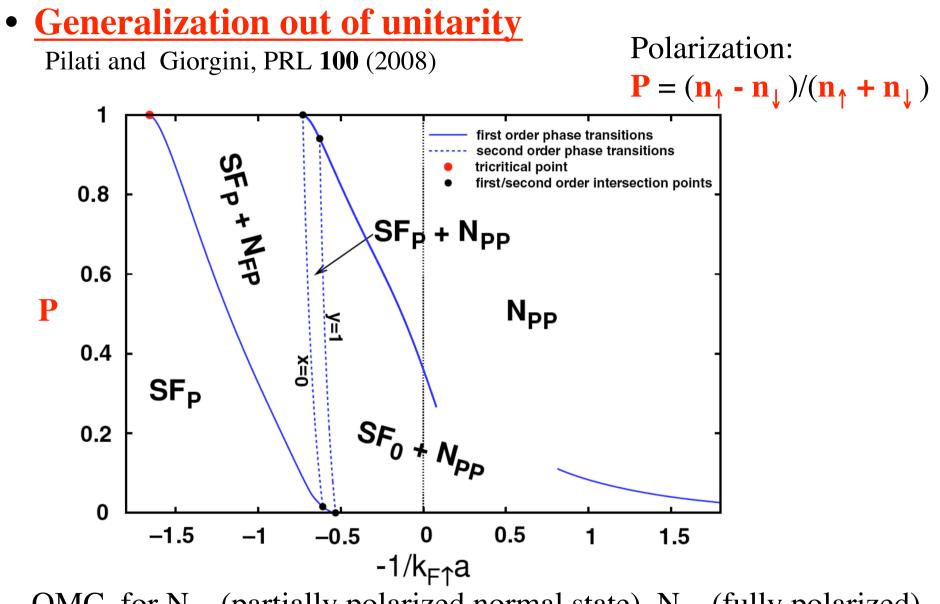
FIG. 3: Radii of the three phases in the trap in units of the radius $R^0_{\uparrow} = a_{\rm ho} (48N_{\uparrow})^{1/6}$ of a noninteracting fully polarized gas, where $a_{\rm ho}$ is the harmonic oscillator length.

• Critical polarization $P_c = 0.77$ in good agreement with experiments



• Very good agreement for density of trapped gas





- QMC for N_{PP} (partially polarized normal state), N_{FP} (fully polarized) for SF_0 (unpolarized superfluid) for SF_P (polarized superfluid)

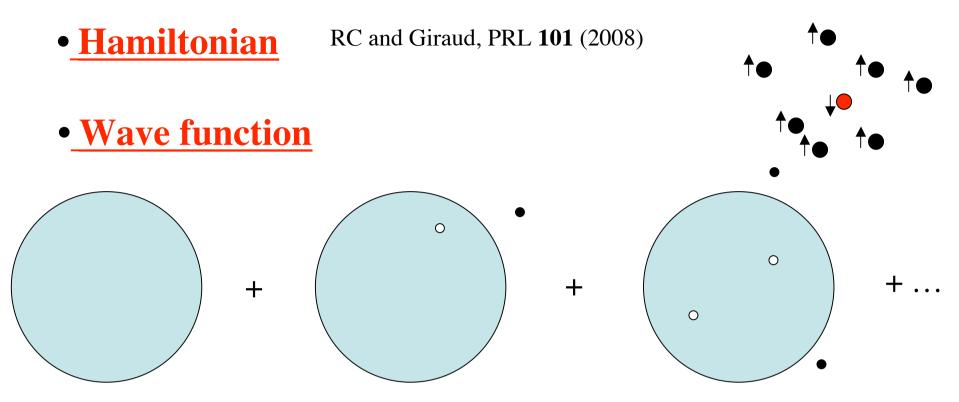
• <u>1 spin ↓ + N spins</u> ↑

- For full control of "Clogston-Chandrasekhar", solution needed for this very interesting "many body" (normal state) problem
- "Simple" since spins ↑ non interacting
- First step: T- matrix approach (quite often used)
 - Unitarity

$$\frac{\mu_{\downarrow}}{E_F} = -0.6066$$
 $\frac{m*}{m} = 1.17$

- Very near Monte Carlo !
- Very surprising → **coincidental** ?





- $\alpha_{\{\mathbf{k}_i\}\{\mathbf{q}_j\}}$ antisymetric in particle and hole variables

$$-H|\psi\rangle = E|\psi\rangle$$

-If weak dependence of kinetic energies on hole variables q_i neglected

Exact decoupling of higher order terms by destructive interference

- In practice very fast convergence \rightarrow n = 2 quite enough

• Check
$$m_{\downarrow} \to \infty$$

- Fermi sea + impurity : exactly solvable (one-body)

- Unitarity $ho = E_b/E_F = 0.5$
 - 1^{rst} order $\rho = 0.465$ 2nd order $\rho = 0.498$
- Convergence essentially complete !
- $\underline{1D}$ Excellent agreement with exact results for energy + effective mass

• Results for
$$m_{\uparrow} = m_{\downarrow}$$

1^{rst} order $\rho = 0.6066$ 2nd order $\rho = 0.6156$

- Extremely fast convergence
- Monte Carlo

Pilati-GiorginiFN-DMC $\rho = (3/5) (0.99 \pm 0.01) = 0.59 - 0.60$ Prokof'ev-SvistunovDiag.MC $\rho = 0.618$ $\rho = 0.615$

- Excellent agreement not coincidental !



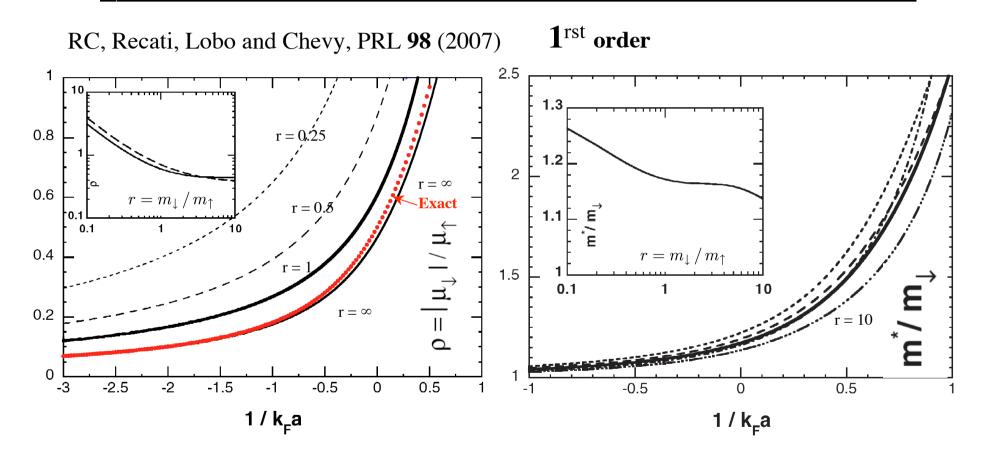
• Effective mass

1^{rst} order

 $m^* / m = 1.17$

2nd order

 $m^* / m = 1.20$



- Nothing special happens at unitarity : unitarity physically not different from BCS side





Very interesting field !

Thank you for your attention !

