

Dynamic Correlation Functions of 1D Quantum Liquids

Alex Kamenev



in collaboration with

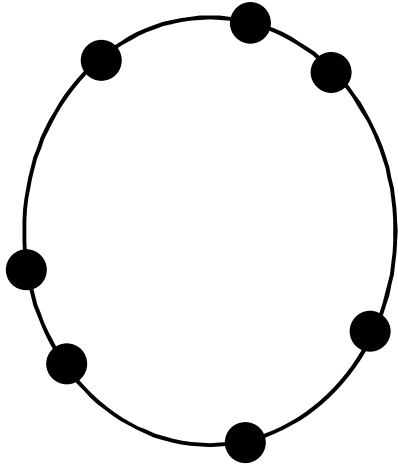
Leonid Glazman, Yale
Maxim Khodas, BNL
Michael Pustilnik, Georgia Tech

PRL **96**, 196405 (2006);
PRL **99**, 110405 (2007);
PRB (2007), PRA (2008).



Moscow, SC4 , May 2009

Models



N –interacting quantum particles on a ring

$$\hat{H} = -\frac{1}{2m} \sum_{i=1}^N \frac{\partial^2}{\partial x_i^2} + \sum_{i<j} V(x_i - x_j)$$

- ✓ Thermodynamic limit
- ✓ Bosons or spinless fermions
- ✓ Translationally invariant
- ✓ Bosons with $V(x_i - x_j) = c\delta(x_i - x_j)$ Lieb-Liniger, integrable
- ✓ (1+1)D complex field theory, “critical”

Observables

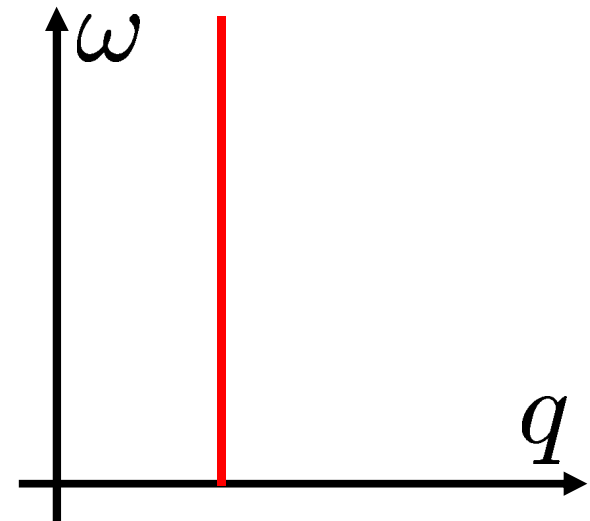
Density

$$\hat{\rho}(x) = \sum_{i=1}^N \delta(x - x_i)$$

$$\hat{\rho}(x, t) = e^{-i\hat{H}t} \rho(x, 0) e^{i\hat{H}t}$$

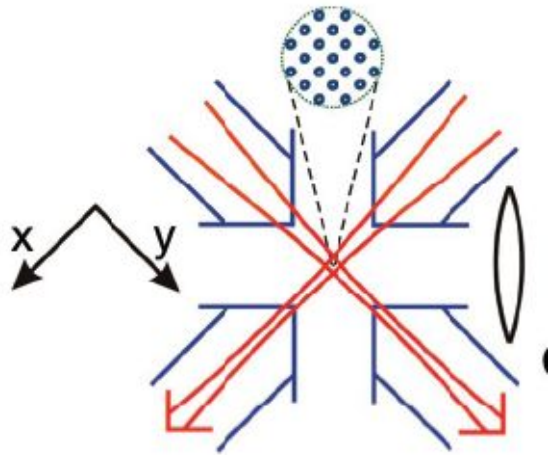
Dynamic Structure Factor (DSF) at T=0

$$S(x, t) = \langle 0 | \rho(x, t) \rho(0, 0) | 0 \rangle$$

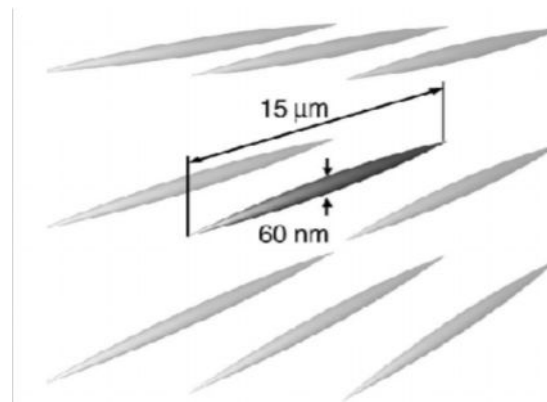


$$S(q, \omega) = \int dx dt e^{i(qx - \omega t)} S(x, t) = \sum_{n_q} |\langle n_q | \hat{\rho} | 0 \rangle|^2 \delta(\omega - \epsilon_{n_q})$$

Cold Atoms in Optical Lattices

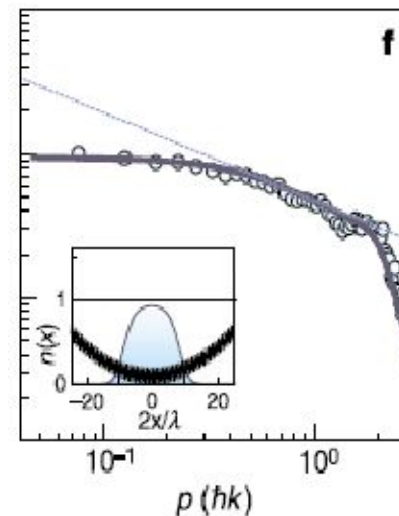
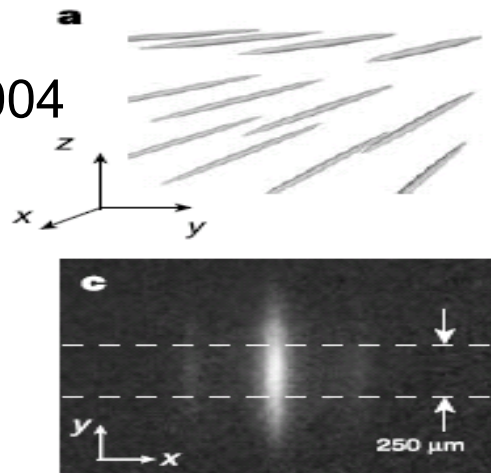


T. Kinoshita, et al 2004

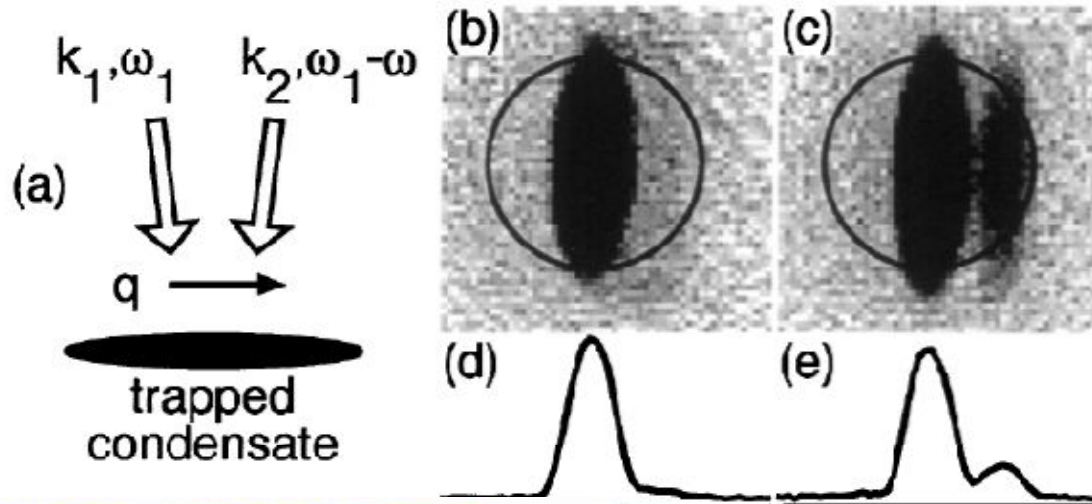


H. Moritz, et al 2003

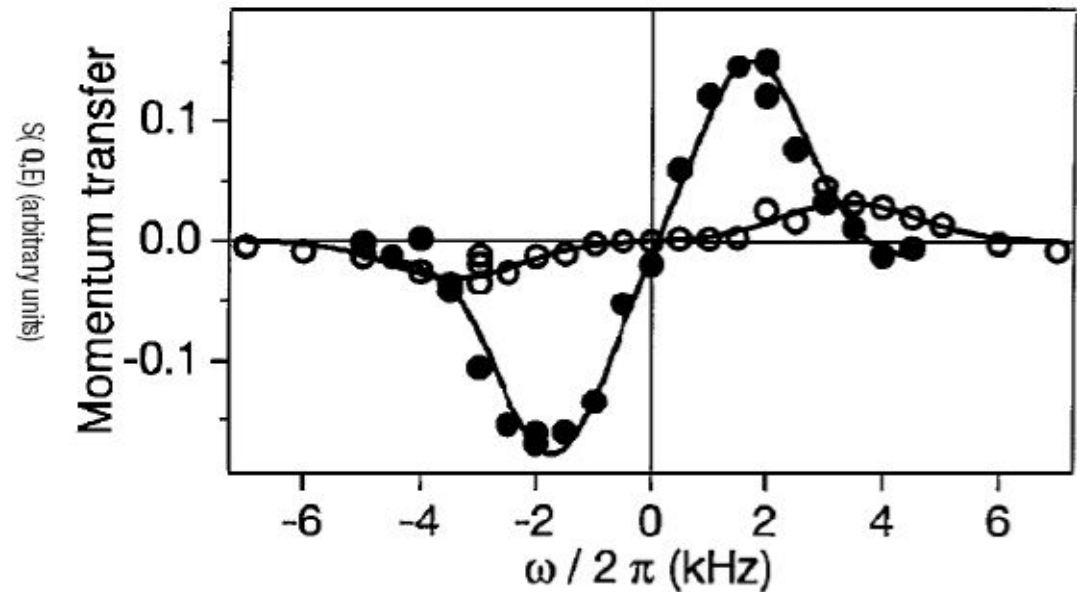
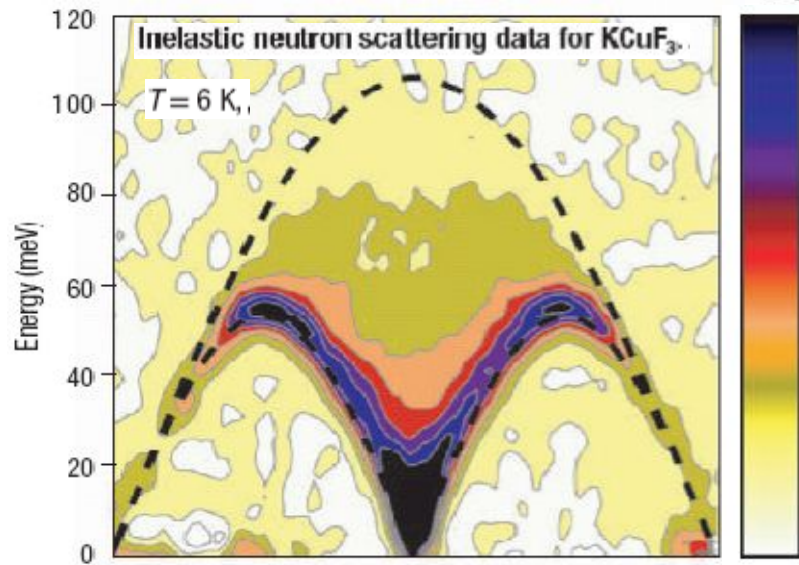
I. Bloch, et al 2004



Bragg Scattering



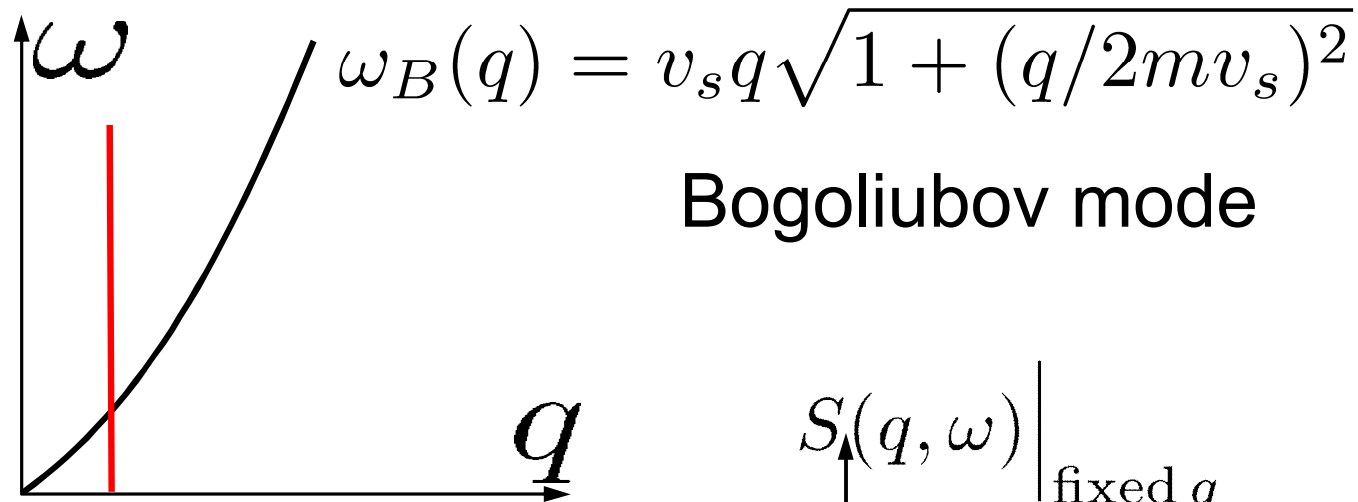
W. Ketterle, et al
2000-...



BELLA LAKE^{1,2*}, D. ALAN TENNANT^{2,3†},
CHRIS D. FROST³ AND STEPHEN E. NAGLER¹

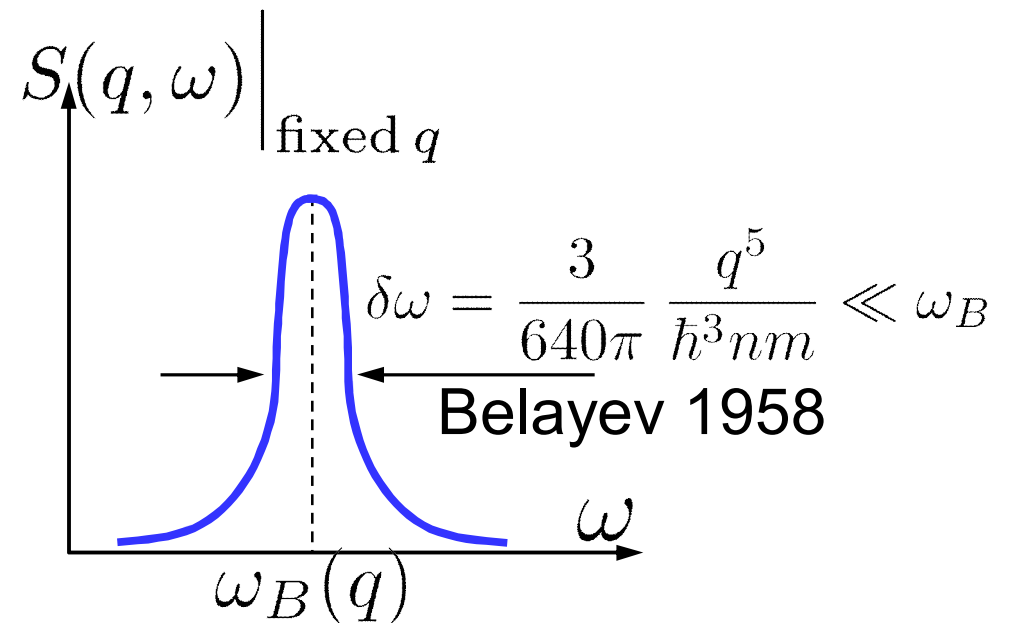
3D Condensates

$$T < T_c$$



Bogoliubov mode

✓ What about 1D?



Exact Analysis of an Interacting Bose Gas. I. The General Solution and the Ground State

ELLIOTT H. LIEB AND WERNER LINIGER

Thomas J. Watson Research Center, International Business Machines Corporation, Yorktown Heights, New York

(Received 7 January 1963)

- ✓ N bosons with delta-functional interactions on a 1D ring

$$H = -\frac{1}{2m} \sum_{i=1}^N \frac{\partial^2}{\partial x_i^2} + c \sum_{i<j} \delta(x_i - x_j)$$

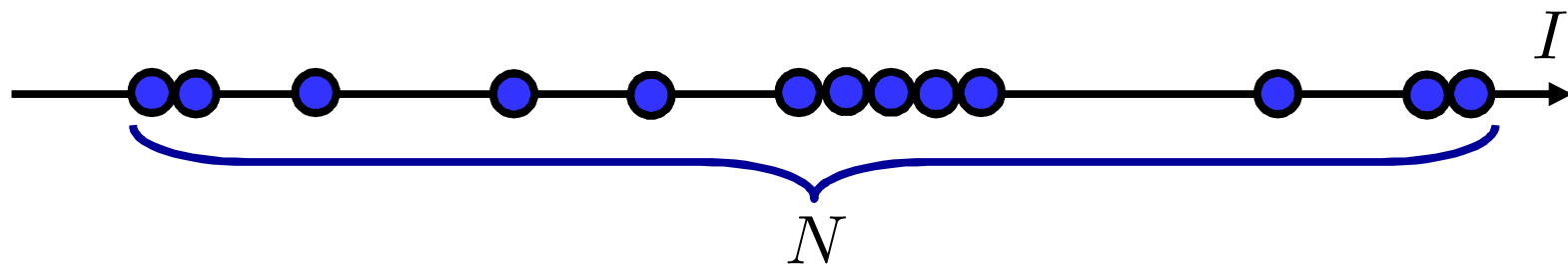
- ✓ Two characteristic momenta: mc and n
 $=N/L$

- ✓ Dimensionless coupling constant: $\gamma = \frac{mc}{n}$

Bethe Ansatz

$$\psi(x_1, x_2, \dots, x_N) = \sum_P a(P) e^{i \sum_{j=1}^N x_j \lambda_j}$$

$$\lambda_j + \frac{1}{L} \sum_k 2 \arctan \frac{\lambda_j - \lambda_k}{m c} = \frac{2\pi}{L} I_j \quad \leftarrow \text{integers}$$

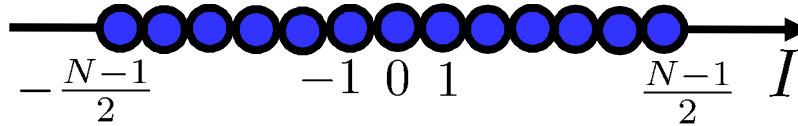


$$E = \frac{1}{2m} \sum_j \lambda_j^2$$

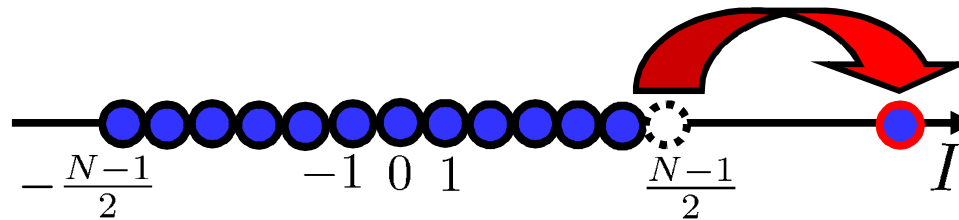
$$q = \sum_j \lambda_j = \frac{2\pi}{L} \sum_j I_j$$

Lieb's Modes

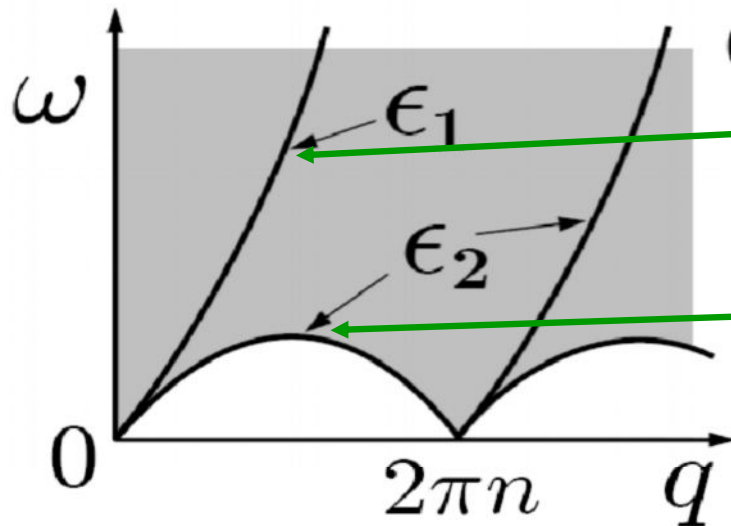
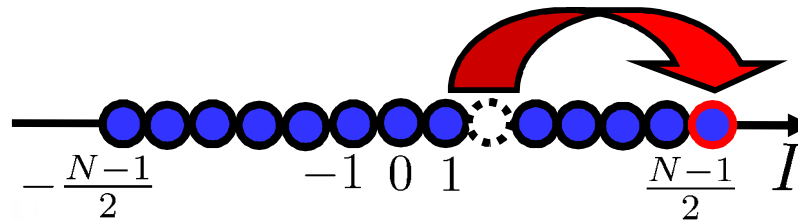
Ground state:



Lieb's I mode
"particles":



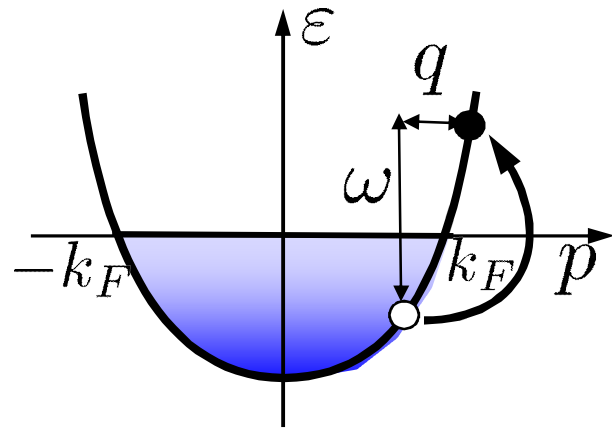
Lieb's II mode
"holes":



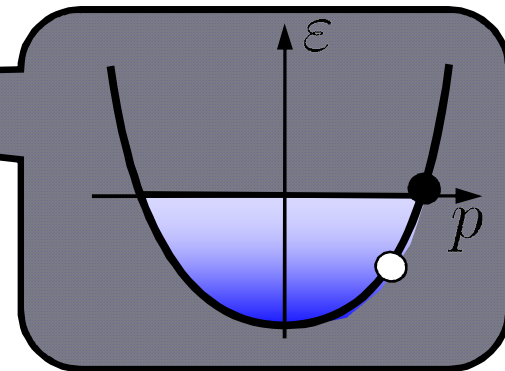
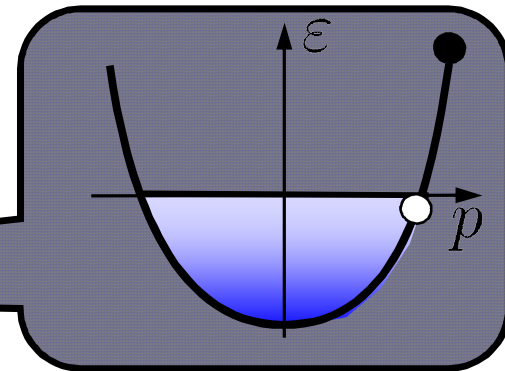
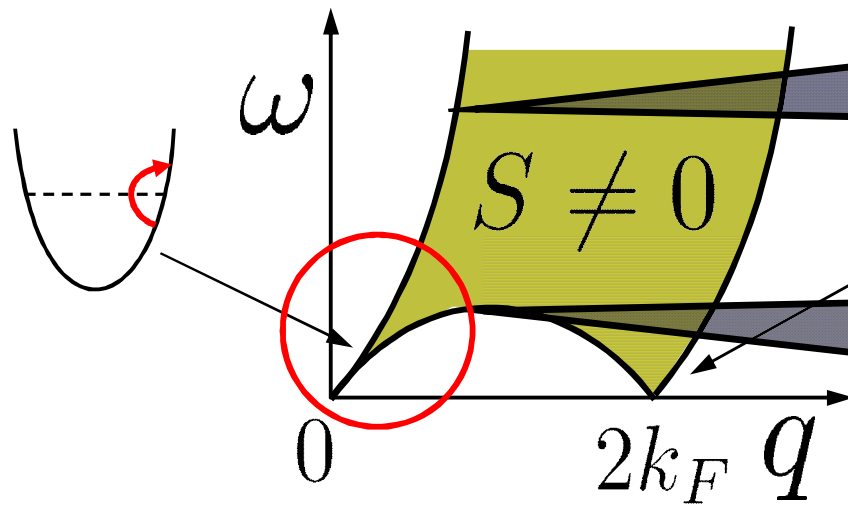
Bogoliubov for $\gamma \rightarrow 0$

Lower bound of the
spectral continuum

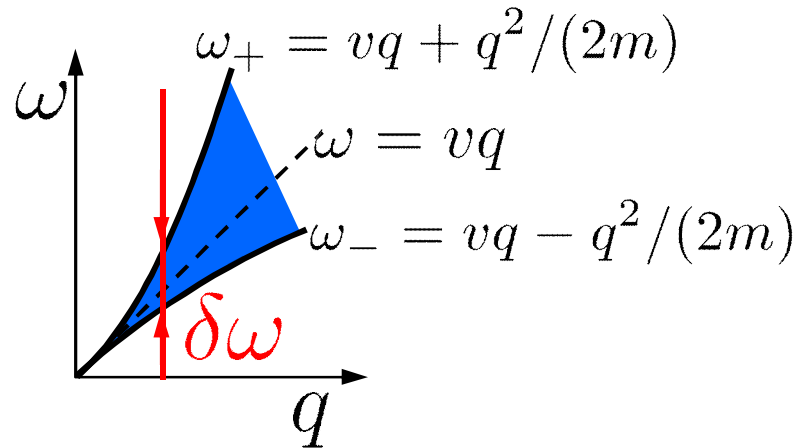
Strongly Interacting Bosons = Free Fermions



$$S(q, \omega) = \langle \rho(q, \omega) \rho(-q, -\omega) \rangle$$



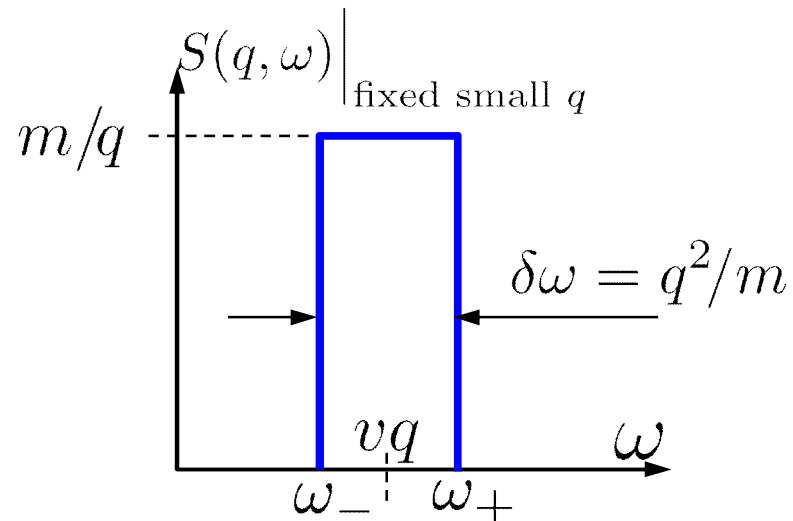
Structure Factor (free fermions)



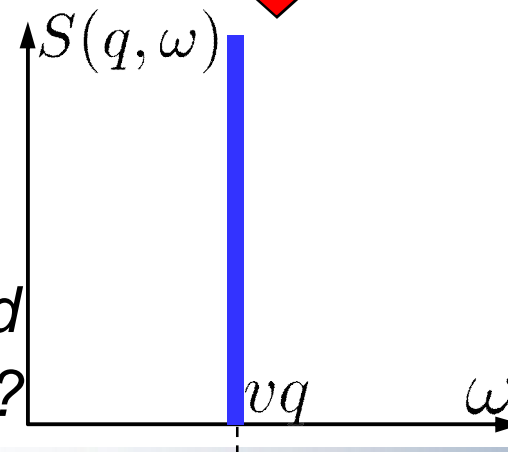
$$S(q, \omega) = q\delta(\omega - vq)$$

✓ **Exact result within the Luttinger approximation.**

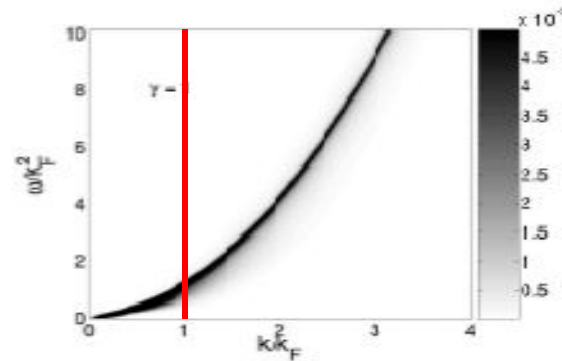
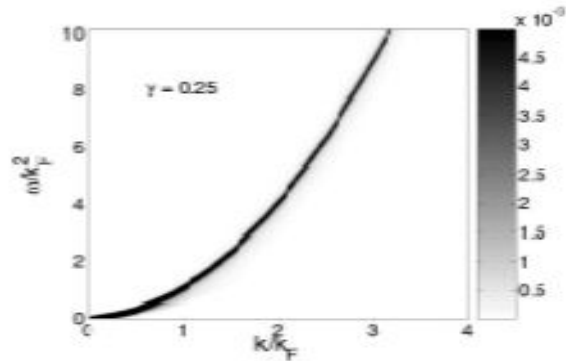
How does the dispersion curvature and interactions affect the structure factor?



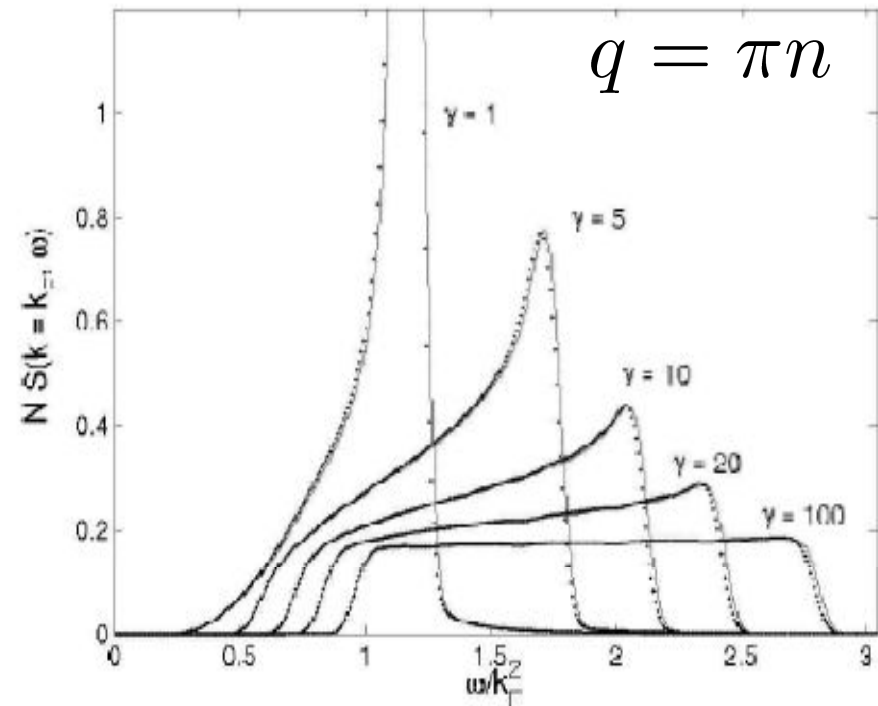
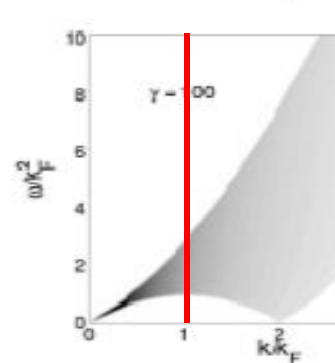
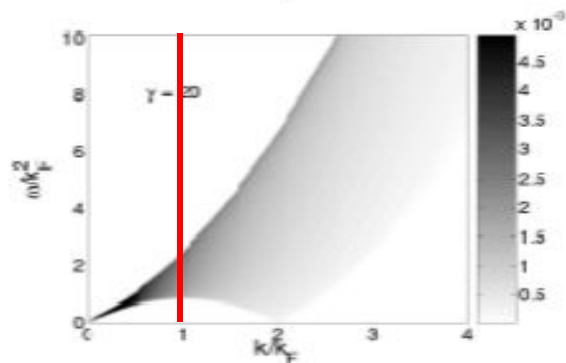
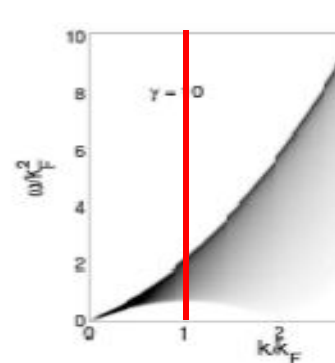
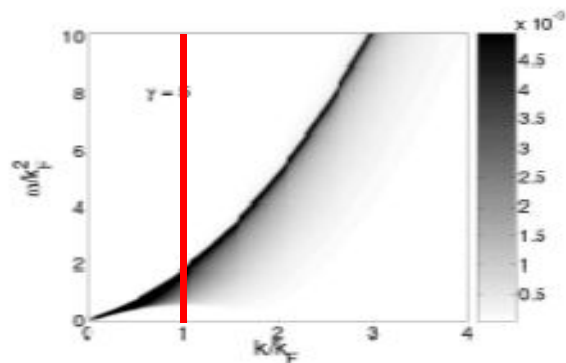
Linear dispersion



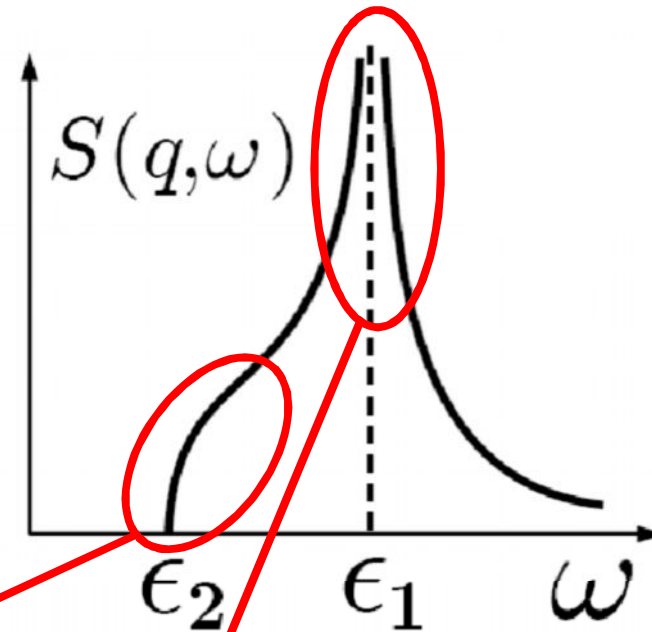
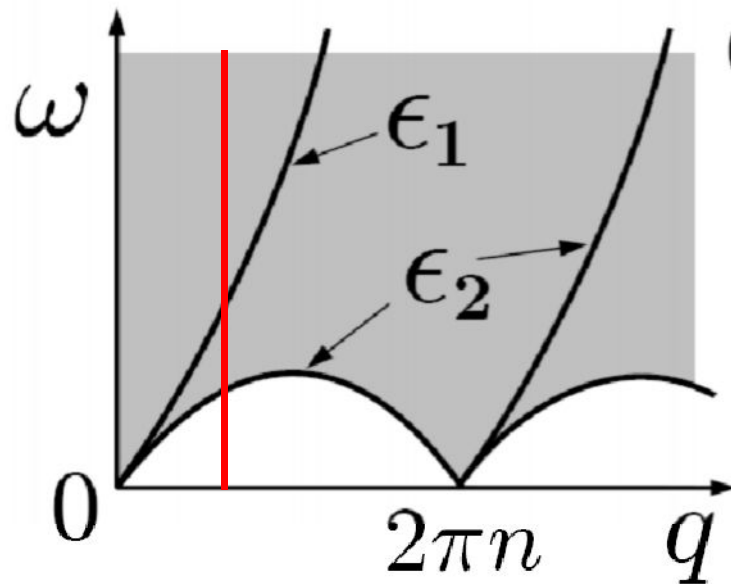
Algebraic BA exact numerics



J-S. Caux, P. Calabrese,
2006
N. Slavnov, 1989



DSF singularities at Lieb's modes



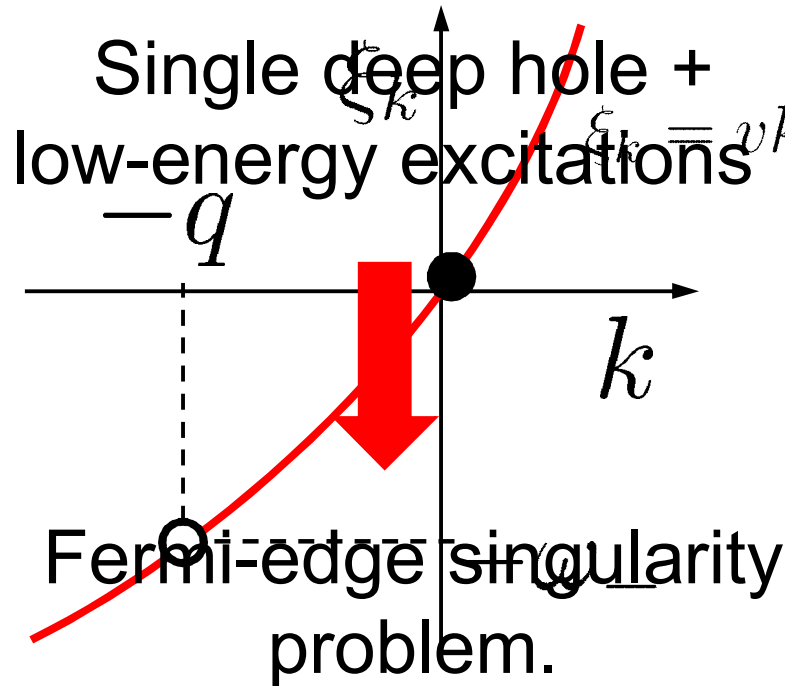
$$S(q, \omega) \sim \frac{m}{q} \left[\frac{\omega - \epsilon_2}{\delta\epsilon} \right]^{\mu_2(q)} \theta(\omega - \epsilon_2)$$

$$\delta\epsilon = \epsilon_1 - \epsilon_2$$

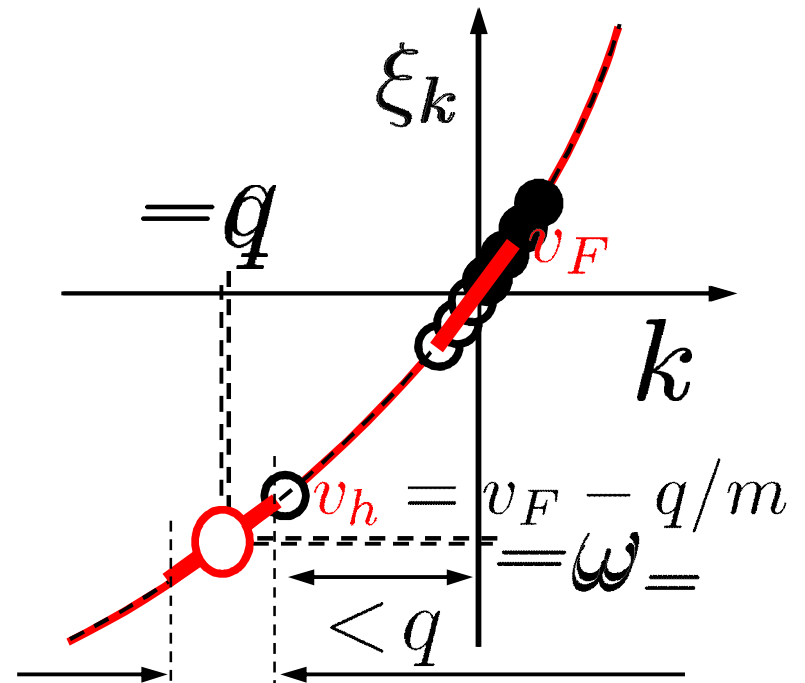
$$S(q, \omega) \sim \frac{m}{q} \left| \frac{\delta\epsilon}{\omega - \epsilon_1} \right|^{\mu_1(q)} [\theta(\epsilon_1 - \omega) + \nu_1 \theta(\omega - \epsilon_1)] \quad \mu_1 < 1$$

Effective model $\omega_- \lesssim \omega$

multipair states with momentum q



single pair
Power-law edge
singularities.



states, contributing to the **leading** logarithm corrections in each order of the perturbation theory.
the idea: project all other states out, linearize remaining spectrum.

Why Power-Law ?

Band of low energy excitations:

$$H_0 = \frac{v}{2\pi} \int dx \left[\frac{1}{K} (\partial_x \phi)^2 + K (\partial_x \theta)^2 \right]$$

Popov, 1973

Efetov, Larkin, 1975

Haldane, 1981

Deep hole creation operator
(instantaneous shift of density and current):

$$\hat{D}(x, t) = e^{i[\delta_\theta \phi(x, t) + \delta_\phi \theta(x, t)]}$$

Dynamic structure factor

$$S(q, \omega) = \int dx dt e^{i(qx - \omega t)} \left\langle \hat{D}(x, t) \hat{D}^\dagger(0, 0) \right\rangle_{H_0}$$



power-law of $x - v_d t$

Exactly solvable models

$\delta_{\phi,\theta}(q)$ can be fixed by:

(i) comparing finite size spectrum of the effective model and Bethe Ansatz spectrum with fixed total momentum q

Pereira, White, Affleck, 2008, 2009

Cheianov, Pustilnik, 2008

Imambekov, Glazman, 2008

Khodas, 2009

(ii) Galilean invariance + dispersion relation of the mode

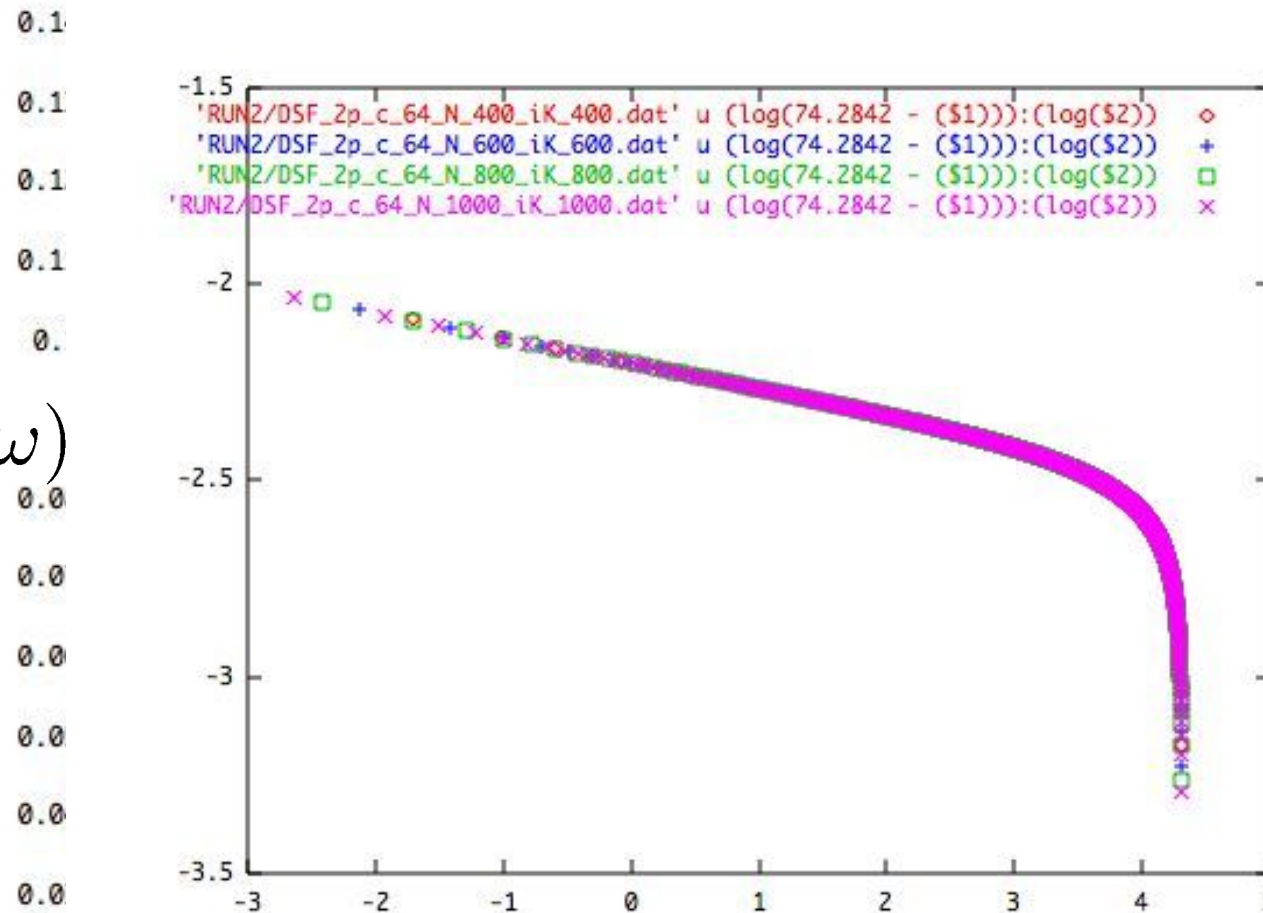
Kamenev, Glazman, 2008

Lamacraft, 2009

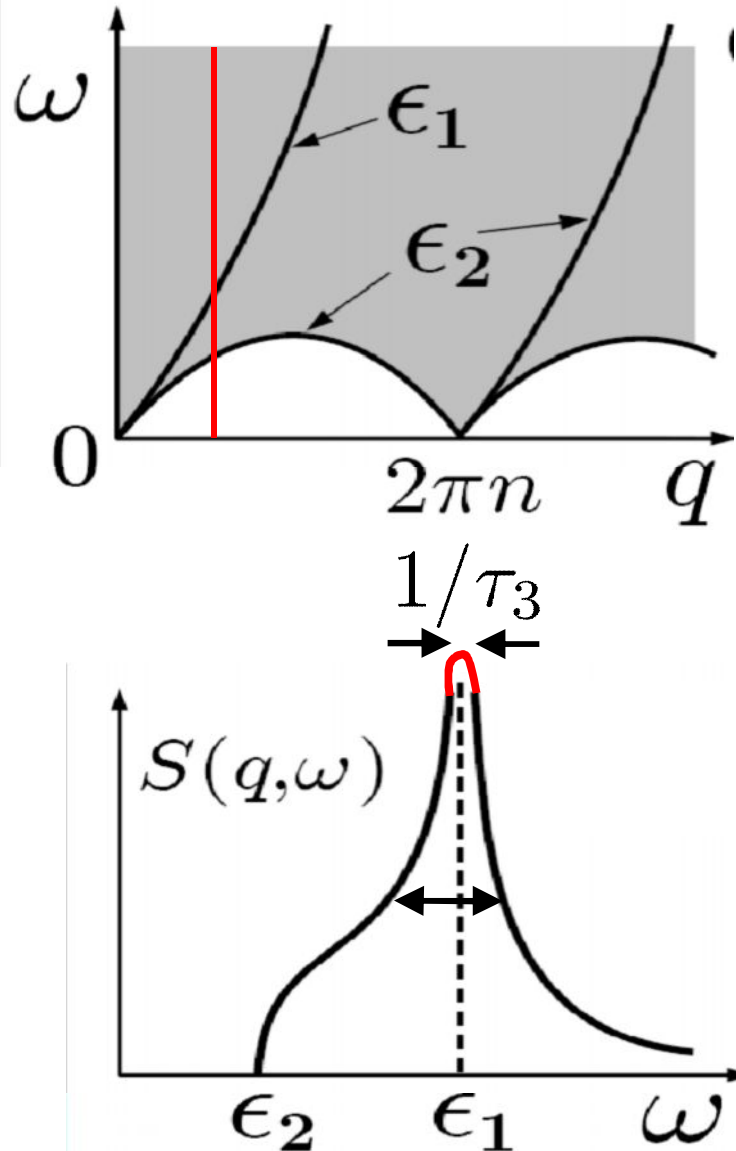
Numerics (preliminary)

Courtesy of J-S. Caux

$$S(2\pi n, \omega)$$

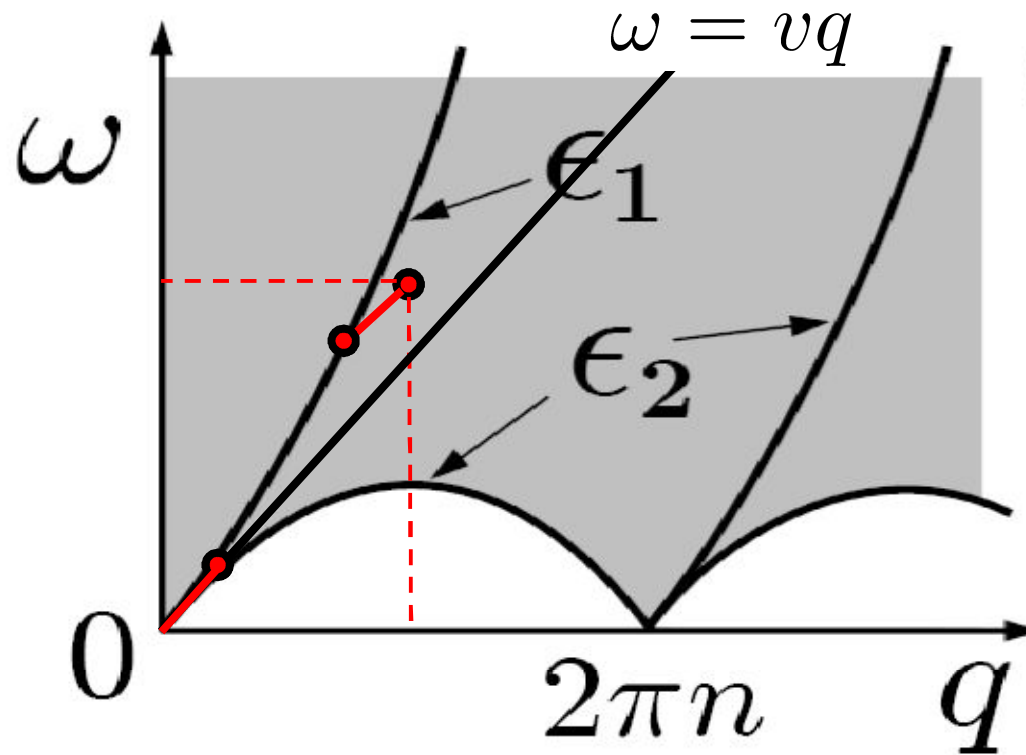


Singularities at Lieb's modes



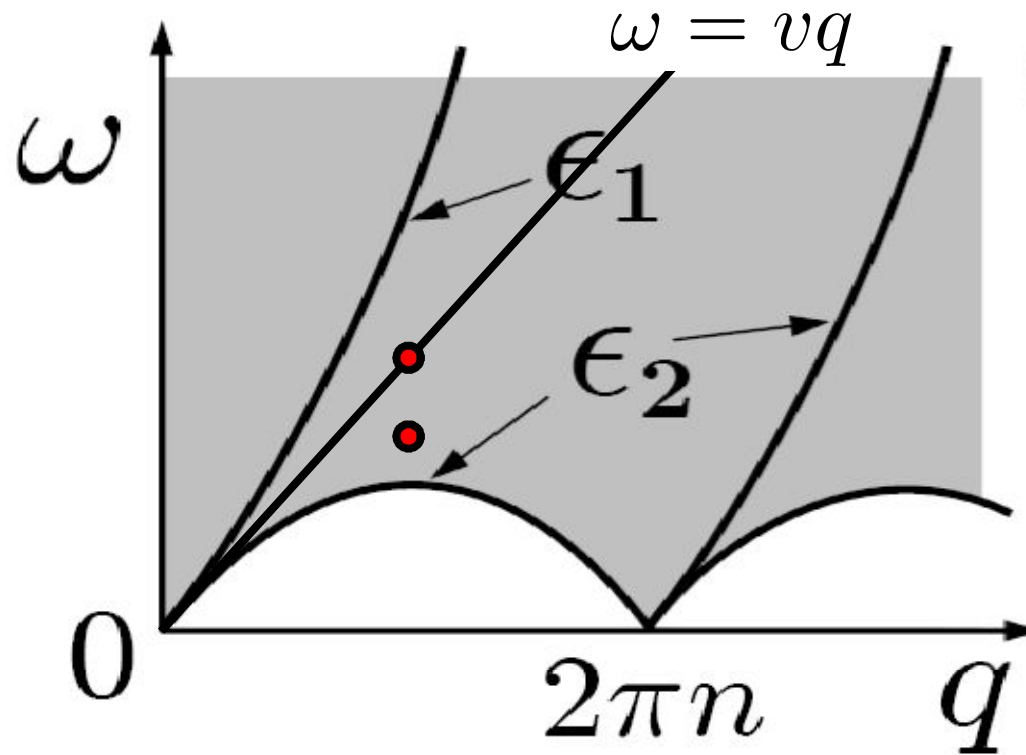
- ✓ For integrable models: positions and exponents are known from TBA
- ✓ Scaling functions ???
- ✓ Singularities are NOT smeared by temperature
- ✓ Non-integrable models: three-body scattering smears singularity in the bulk, but NOT at the edge

Weakly Interacting Bosons



Phonon modes shake-up \rightarrow power-law singularity

Weakly Interacting Bosons



For $\omega = vq$ infinite number of quasiparticles is excited.

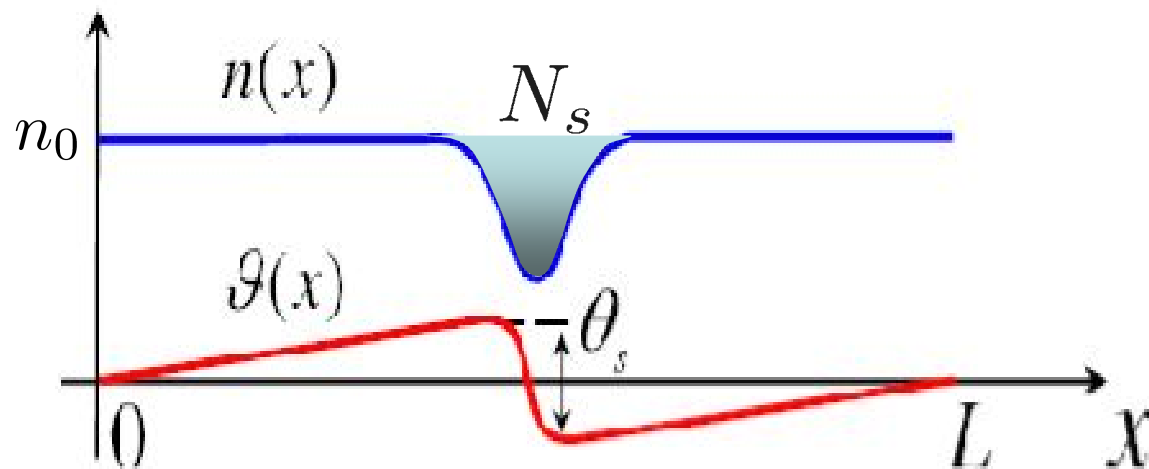
What are the excitations below $\omega = vq$ line ?

Dark Solitons

$$i\partial_t\psi = -\frac{1}{2m}\nabla^2\psi + c|\psi|^2\psi - \mu\psi$$

Gross-Pitaevskii
equation

$$\psi(x - Vt) = \sqrt{n(x - Vt)} e^{i\vartheta(x - Vt)}$$

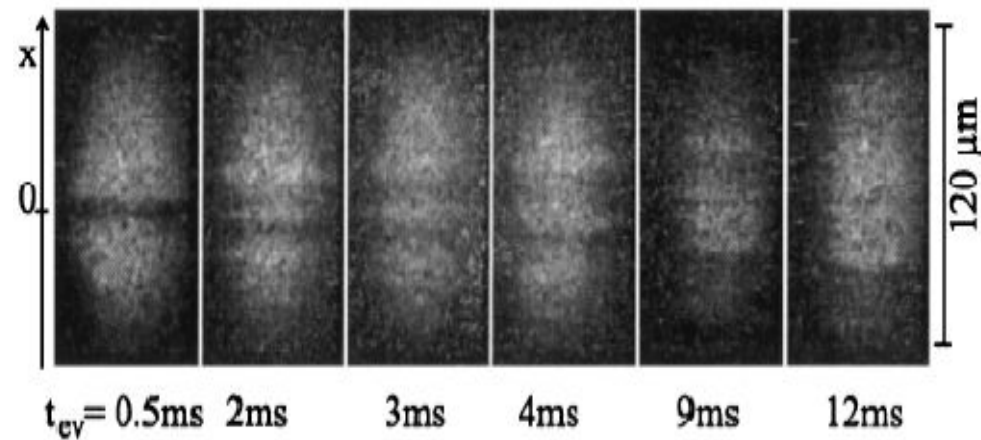


$$N_s = \frac{2}{\sqrt{\gamma}} \sqrt{1 - \frac{V}{v_B}} \gg 1$$

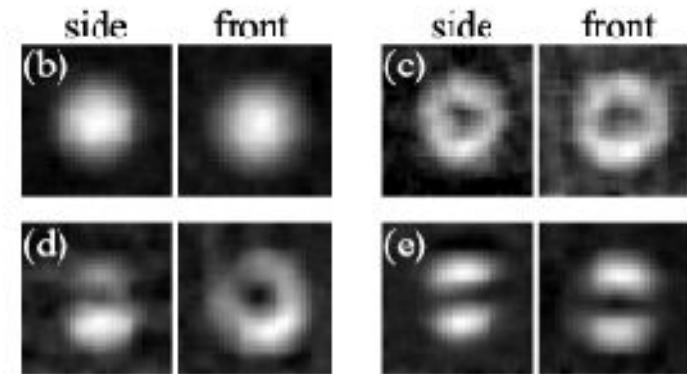
$$\theta_s = 2 \arccos(V/v_B)$$

$$p_s = n_0 (\theta_s - \sin \theta_s) ; \quad \varepsilon_s = \frac{4}{3} n_0 v_B \sin^3(\theta_s/2)$$

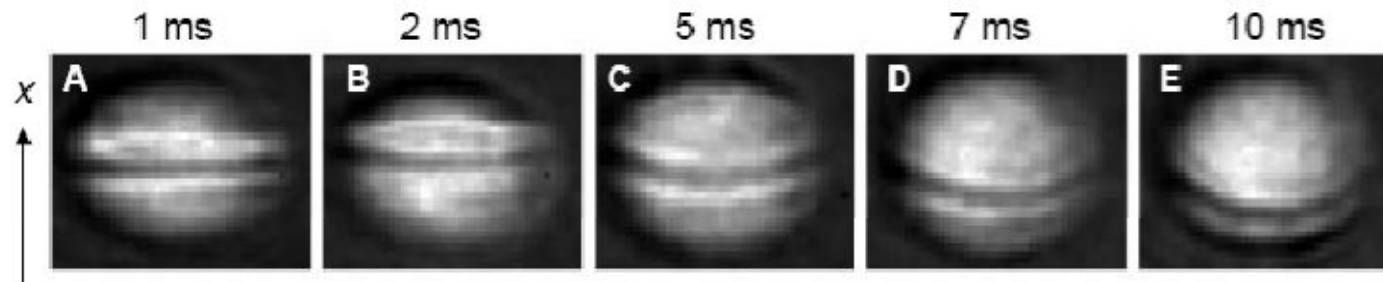
Observations of Dark Solitons



Sengstock, et al. 1999



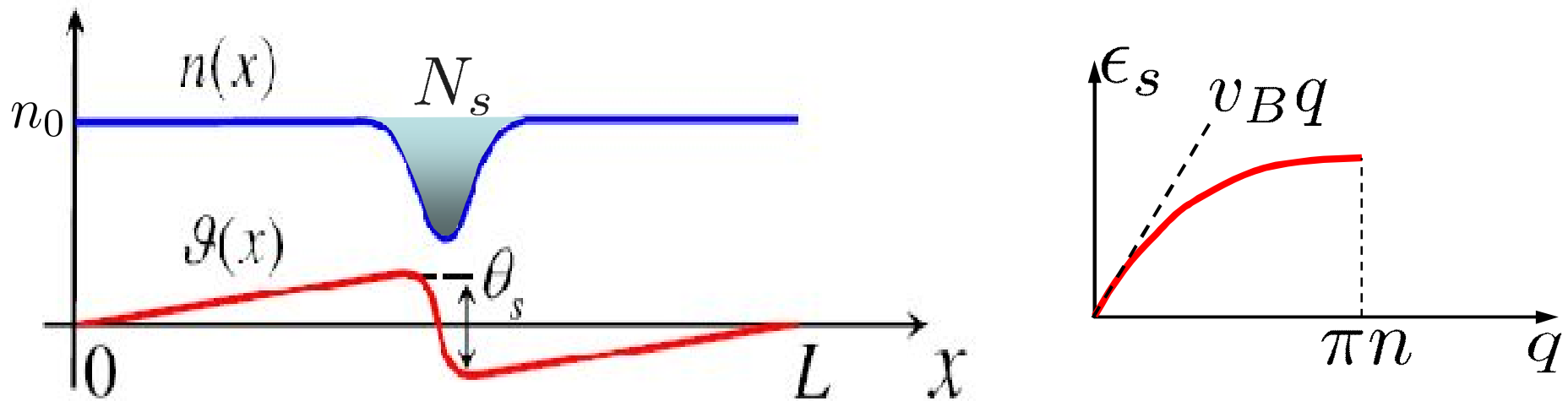
Cornell, et al. 2001



Phillips, et al. 2000

Dark Solitons as Lieb II excitations

$$p_s = n_0 (\theta_s - \sin \theta_s) ; \quad \epsilon_s = \frac{4}{3} n_0 v_B \sin^3(\theta_s/2)$$



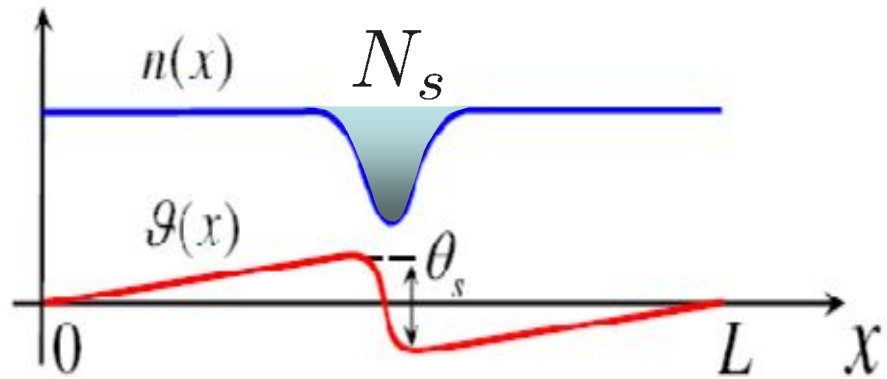
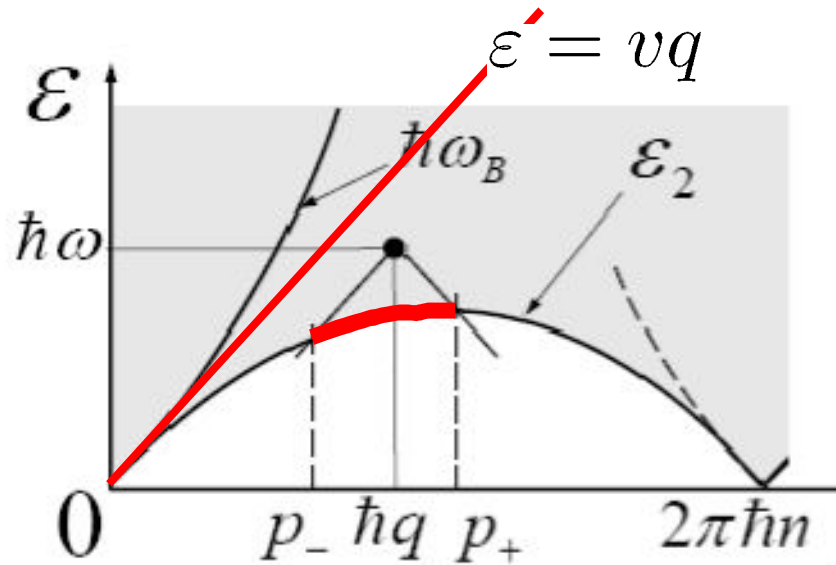
$$\epsilon_s(q) \xrightarrow{\gamma \ll 1} \epsilon_2(q)$$

P.P. Kulish, S.V. Manakov
and L.D. Faddeev 1976

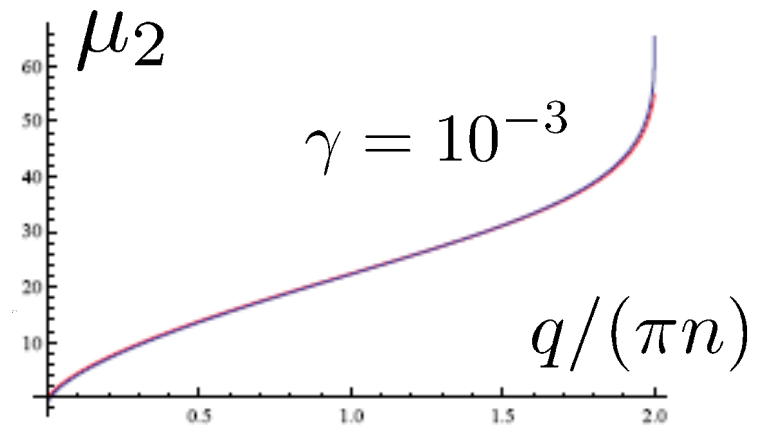
Dark soliton

Lieb II mode

Photo-Solitonic Effect



$$S \sim (\delta\varepsilon)^{\mu_2}$$



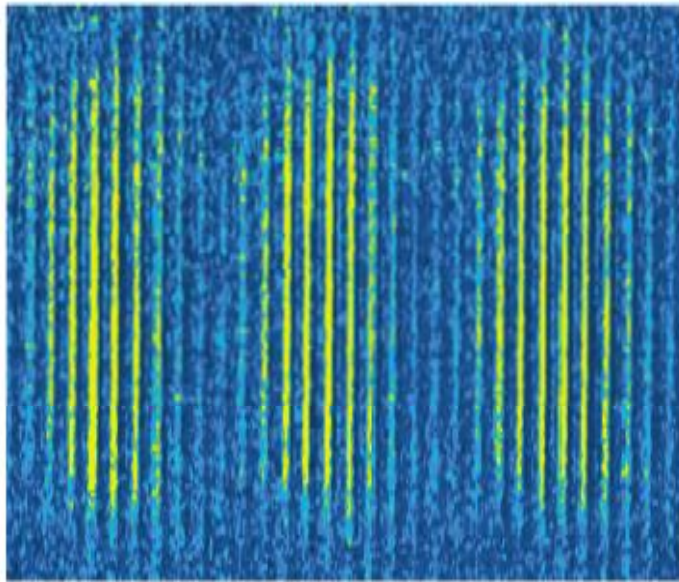
Probability to excite a soliton is suppressed by orthogonality

Wake up!

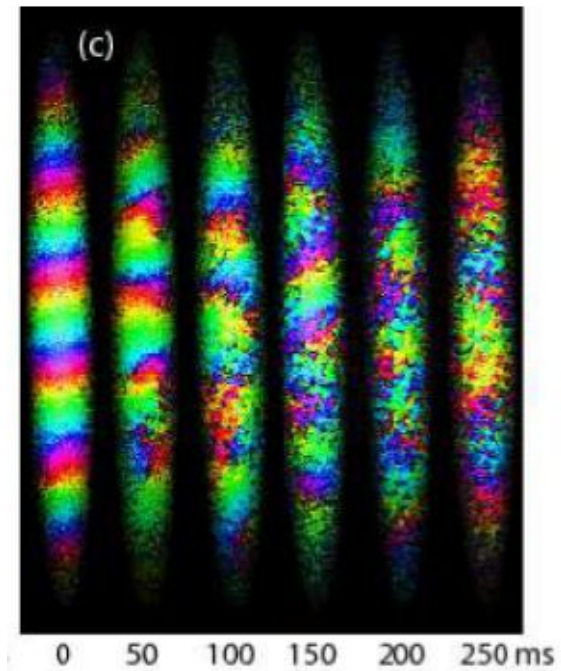
- ✓ Power-law singularities at Lieb modes, where exponents are functions of q
- ✓ Single energetic particle + low-energy excitations
- ✓ Single particle = quasiparticle or soliton
- ✓ Photo-Solitonic Effect

1D Spinor Condensates

Hyperfine states, e.g. $F=1$. Ferromagnetic ground state

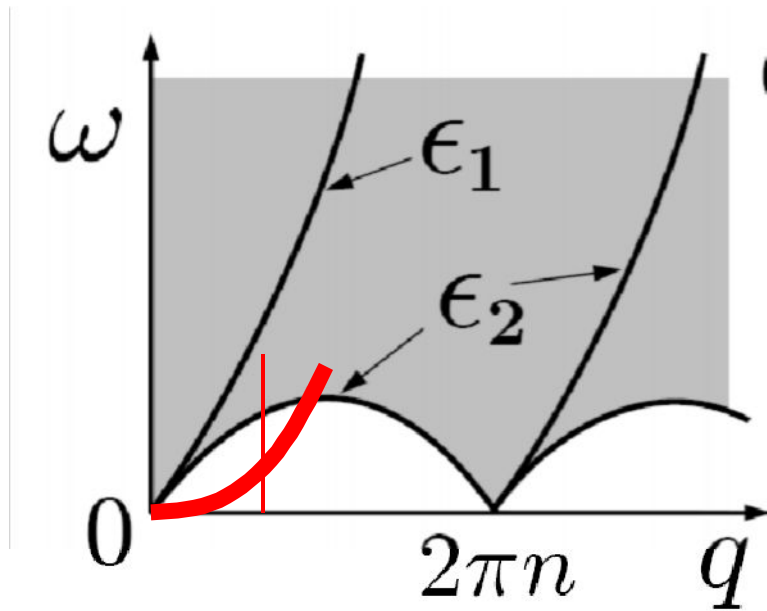


J.M. Higbie, et al 2008



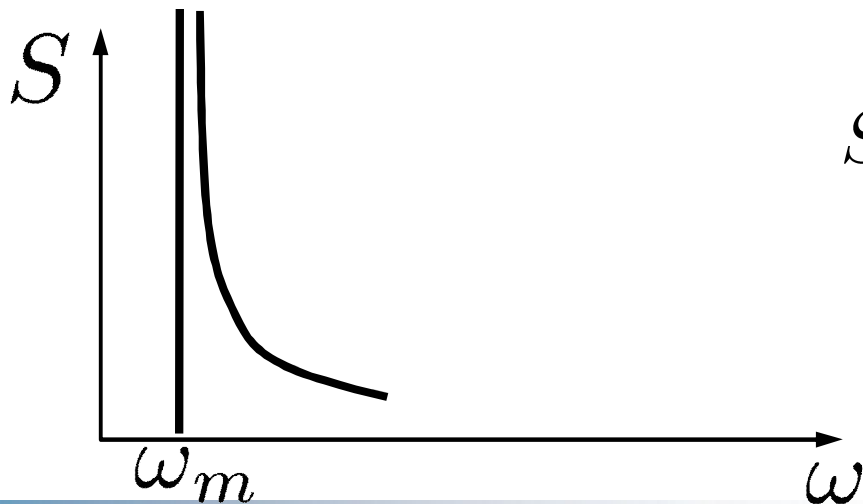
M. Vingalatore, et al 2008

1D Spinor Condensates



Ferromagnetic magnon

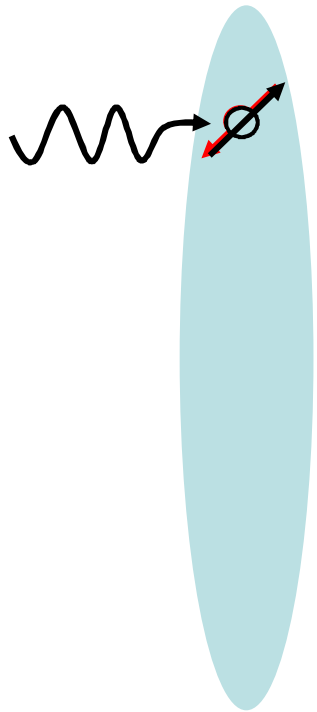
$$\omega_m = \frac{q^2}{2m^*}$$



$$S(q, \omega) \sim \frac{1}{(\omega - \omega_m(q))^{\mu_m(q)}}$$

Gravitational Fall in Spinor Condensate

M. Kohl




$$\dot{v} = g - \Gamma v$$

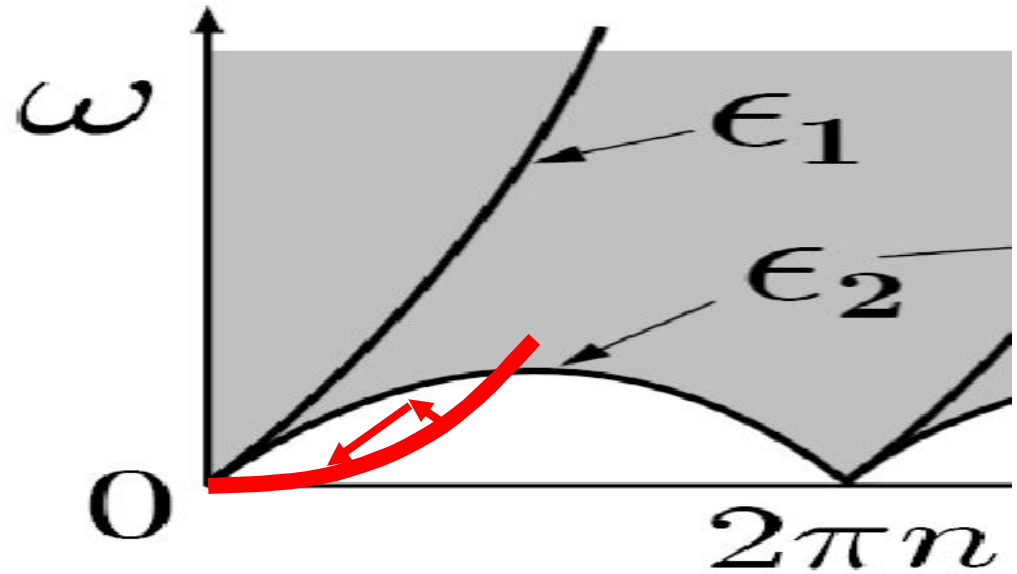
$$v = g/\Gamma \sim T^{-4} \left(\frac{v_g}{v_B} \right)^3$$

$$\frac{\hbar}{mv_g} = \frac{v_g^2}{g};$$

$$v_g = \left(\frac{4.1}{A} \right)^{1/3} \frac{cm}{s}$$

Atomic number


detector



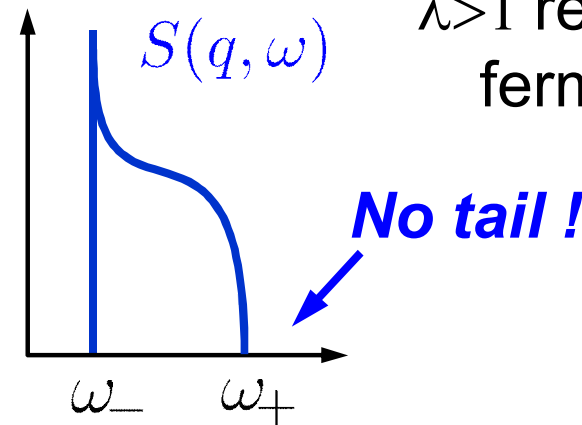
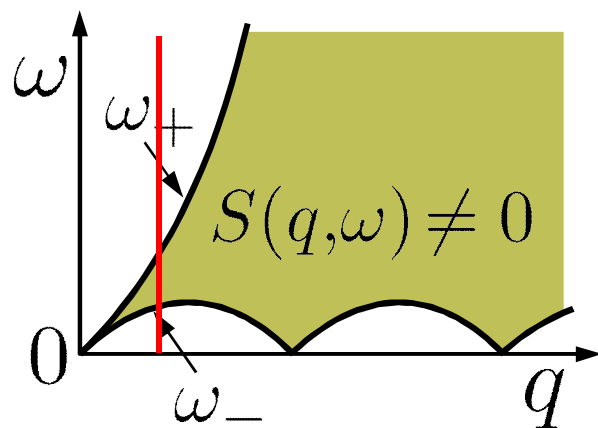
Calogero-Sutherland model

$$H = - \sum_i \frac{1}{2m} \frac{\partial^2}{\partial x_i^2} + \sum_{i < j} V(x_i - x_j) \quad V(x) = \frac{\lambda(\lambda - 1)\pi^2/m}{\sin^2(\pi x)}$$

Haldane 1994, Ha 1995

$$\lambda = \frac{k}{n}$$

S = $k + n$ dim
integral
 $\lambda > 1$ repulsive
fermions



$S(q, \omega)$: power-law singularities at $\omega \rightarrow \omega_{\pm}$

$$S \propto (\omega - \omega_-)^{1/\lambda - 1}$$

$$S \propto (\omega_+ - \omega)^{\lambda - 1}$$

Pustilnik, 2006