

Miracle of Conformal Bootstrap in 4-Dimensions

Slava Rychkov

(Scuola Normale Superiore, Pisa)



- with **R.Rattazzi, E.Tonni, A.Vichi** 0807.0004 (JHEP 0812)
- with **A.Vichi** 0905.2211

Status of CFT in 4D

- **Abstract theory** – scaling operators, OPEs, conformal partial waves
- **Few (non-SUSY) concrete examples**
- E.g. **conformal windows** of $SU(N_c)$ gauge theories with N_f flavors. *Evidence from*
 - large N **Belavin, Migdal'74**
 - SUSY **Seiberg'94**
 - lattice simulations
 $N_c=3, N_f=12$: **Appelquist et al, Deuzeman et al'09**
- **No theoretical control over these fixed points**

Why want to know more?

- curiosity
- what if plays a role in Nature?
 - Unparticles Georgi'07
 - Conformal Technicolor Luty, Okui'04
 - Conformal SUSY-breaking sectors
Roy, Schmaltz'08

CTC - Ideal theory of EWSB

Standard Model 😊 $y_{ij} H \bar{q}_i q_j$ unwanted flavor effects decouple $\frac{1}{\Lambda_{UV}^2} \bar{q}_i q_j \bar{q}_k q_l$

$[H] = 1$

☹️ very relevant operator $\Lambda_{UV}^2 |H|^2$
 $\Lambda_{UV} \rightarrow \infty$ problematic } **Hierarchy problem**

Usual Technicolor 😊 no relevant singlet scalar

$H_{TC} \sim \bar{\psi} \psi$

$[H_{TC}] = 3$ ☹️ Yukawas $y_{ij} \frac{H_{TC}}{\Lambda_{UV}^2} \bar{q}_i q_j$ as relevant as $\frac{1}{\Lambda_{UV}^2} \bar{q}_i q_j \bar{q}_k q_l$

«Conformal Technicolor»

Luty, Okui 2004

(based on earlier «walking TC» idea Holdom'81)

😊 Flavor: $[H] \approx 1$

😊 Hierarchy: $[H^+ H] \geq 4$

- We do not know of any such theories
- Is this at all possible?

\implies **A concrete question**

In *arbitrary* unitary CFT, take a Hermitean scalar operator, with OPE

$$\varphi \times \varphi = 1 + O + \dots$$

leading scalar $O \sim \varphi^2$

- Q:**
- Is there any *upper bound* on $[O]$, the dimension of O , in terms of $[\varphi]$?
 - In particular, is it true that $[O] \rightarrow 2$ as $[\varphi] \rightarrow 1$?

Recall unitarity bound:

$[\varphi] \geq 1$ for any scalar operator in a unitary CFT

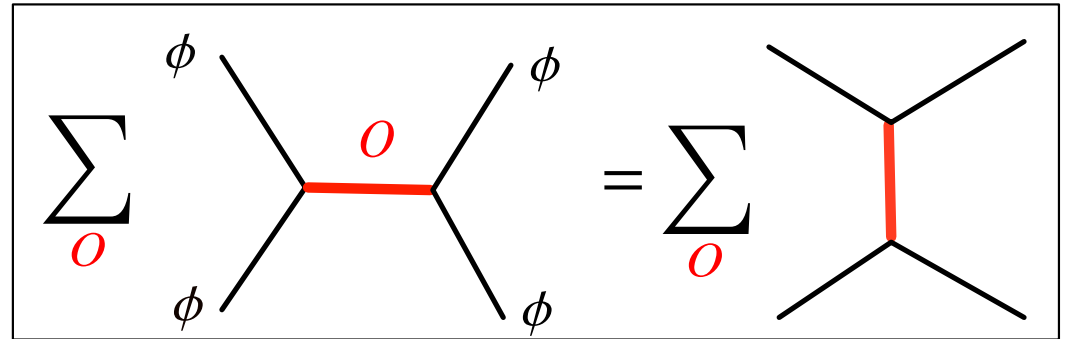
If $[\varphi] = 1$, the scalar is free, and in particular $O = :\varphi^2:$, $[O] = 2$

Conformal Bootstrap – our only weapon

OPE + crossing symmetry

Polyakov'74

Belavin, Polyakov, Zamolodchikov'84



A germ of an idea:

- The Unit Operator exchange is NOT crossing-symmetric

$$\frac{1}{x_{12}^{2d} x_{34}^{2d}} \neq \frac{1}{x_{14}^{2d} x_{23}^{2d}}$$

- Exchanges of scalars and higher spins should restore crossing
- Can it be that higher spins cannot do it by themselves, if all scalars are decoupled?

Conformal block decomposition – divide and conquer

$$\langle \phi(x_1) \phi(x_2) \phi(x_3) \phi(x_4) \rangle = \frac{g(u, v)}{|x_{12}|^{2d} |x_{34}|^{2d}}$$

$$d \equiv [\varphi]$$

$$\langle \phi\phi\phi\phi \rangle = \sum_0 \underbrace{\begin{array}{c} 1 \quad \quad \quad 3 \\ \diagdown \quad \diagup \\ \text{O} + \text{descendants} \\ \diagup \quad \diagdown \\ 2 \quad \quad \quad 4 \end{array}}_{\text{conformal block } \{\Delta, l\}}$$

$$g(u, v) = 1 + \sum_{\Delta, l} (\lambda_{\Delta, l})^2 g_{\Delta, l}(u, v)$$

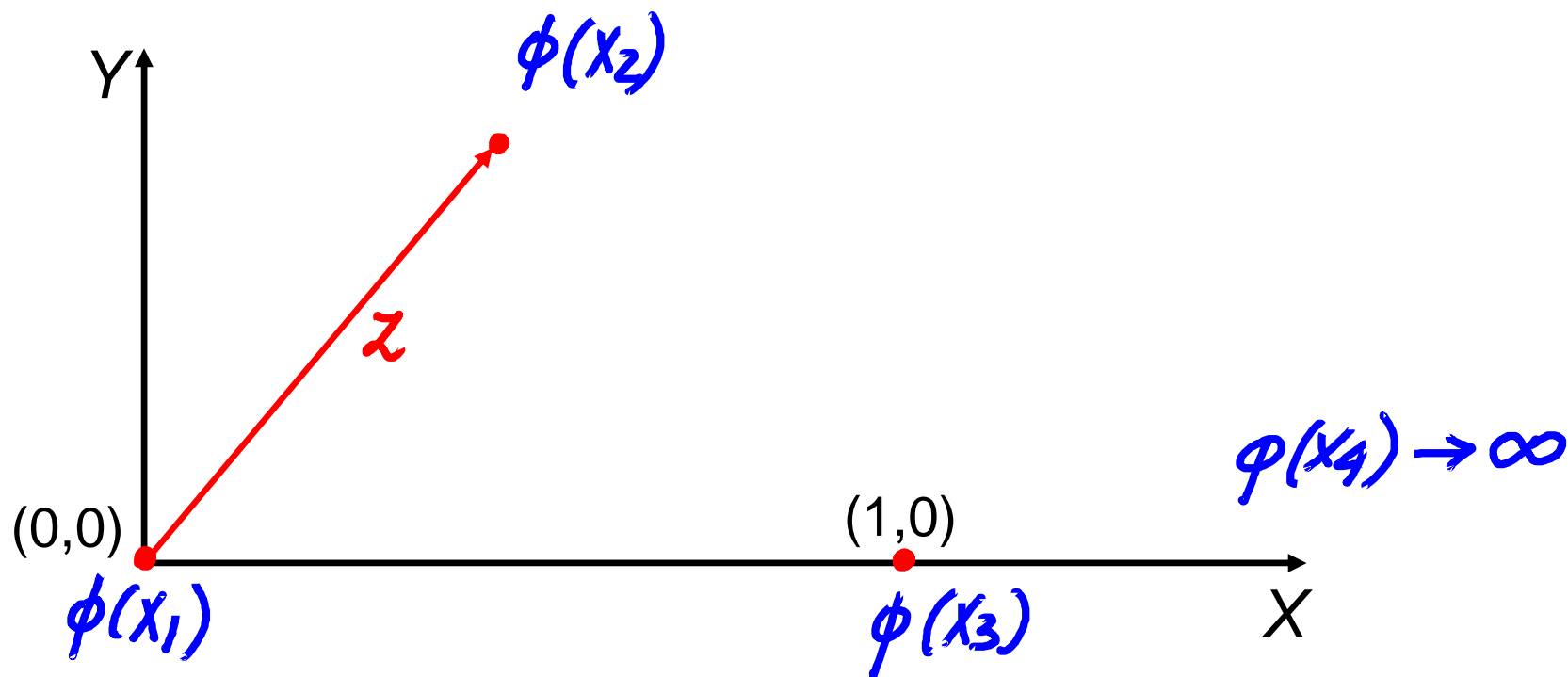
4D Conformal Blocks in closed form [Dolan, Osborn, 2001]

It makes you feel powerful!

$$g_{\Delta,l}(u,v) = \frac{z\bar{z}}{z-\bar{z}} [f_{\Delta+l}(z)f_{\Delta-l-2}(\bar{z}) - (z \leftrightarrow \bar{z})]$$

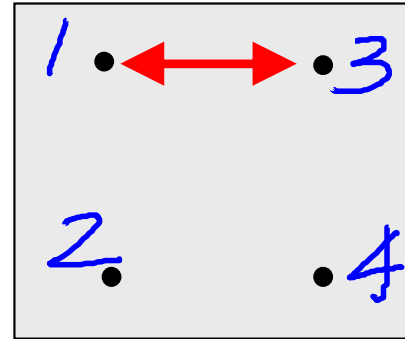
$$f_{\beta}(z) = z^{\beta/2} {}_2F_1\left(\frac{\beta}{2}, \frac{\beta}{2}, \beta; z\right)$$

$$u = z\bar{z}, \quad v = (1-z)(1-\bar{z})$$



Crossing Symmetry – can we balance the budget deficit?

$$v^d g(u, v) = u^d g(v, u)$$



crossing deficit
from unit operator

$$u^d - v^d = \sum_{\Delta, l} (\lambda_{\Delta, l})^2 \left[v^d g_{\Delta, l}(u, v) - u^d g_{\Delta, l}(v, u) \right]$$

Sum Rule:

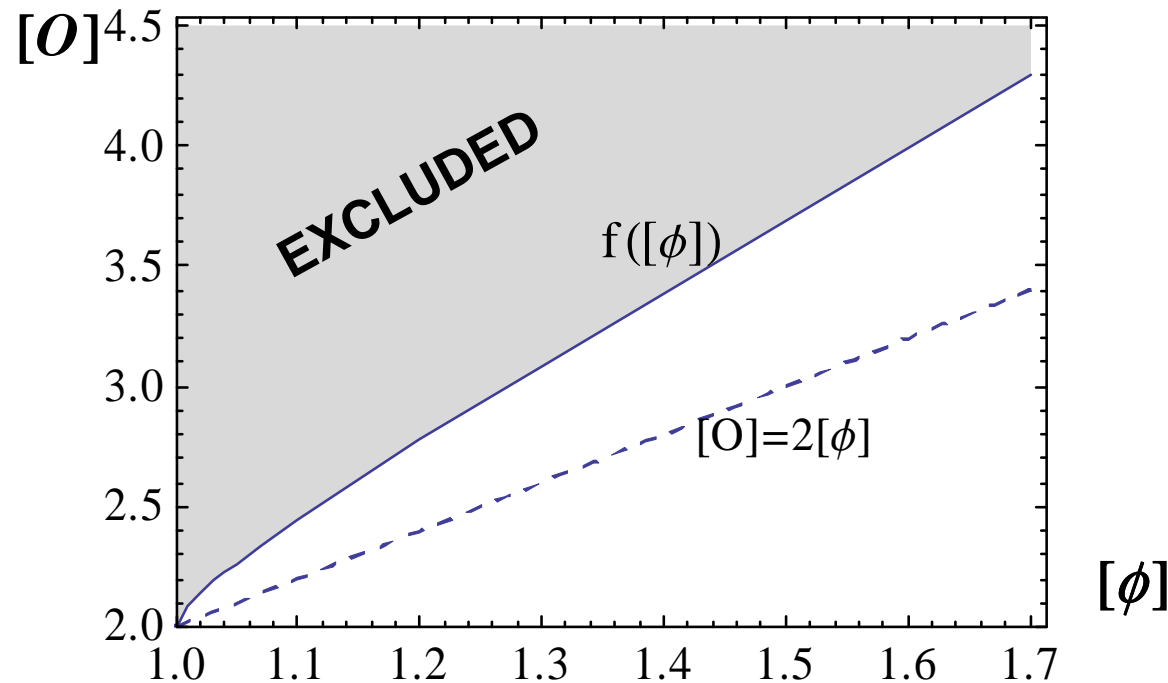
$$1 = \sum_{\Delta, l} \lambda_{\Delta, l}^2 F_{d, \Delta, l}(u, v)$$

$$F_{d, \Delta, l}(u, v) := \frac{v^d g_{\Delta, l}(u, v) - u^d g_{\Delta, l}(v, u)}{u^d - v^d}$$

Sum rule \implies Results

Theorem (with Rattazzi, Tonni, Vichi):

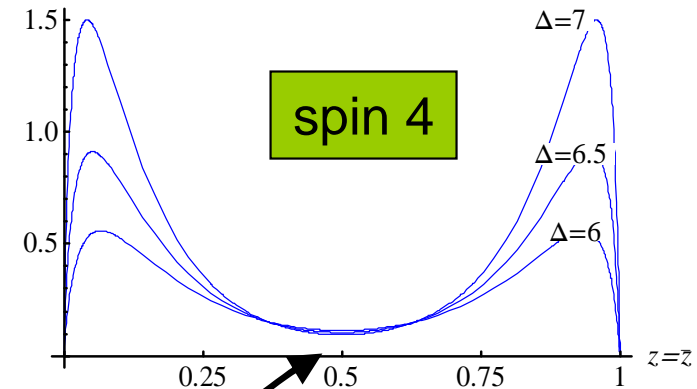
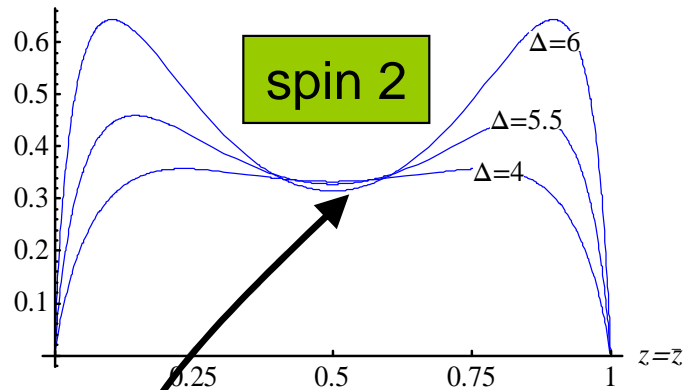
OPE $\varphi \times \varphi$ must contain at least one scalar O of dimension $[O] < f([\varphi])$



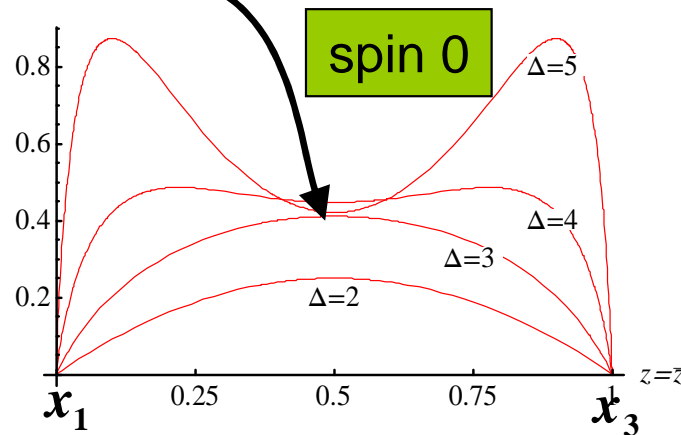
- Numerical fit: $f(d) \approx 2 + 0.7\sqrt{d-1} + 2.1(d-1) + 0.43(d-1)^{3/2}$
- $f(d) \rightarrow 2$ as $d \rightarrow 1$

How can this be at all possible?

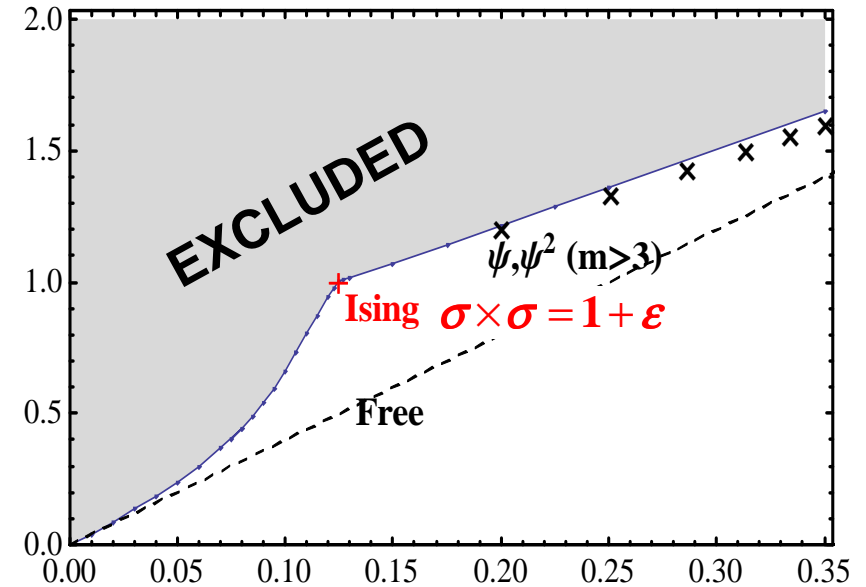
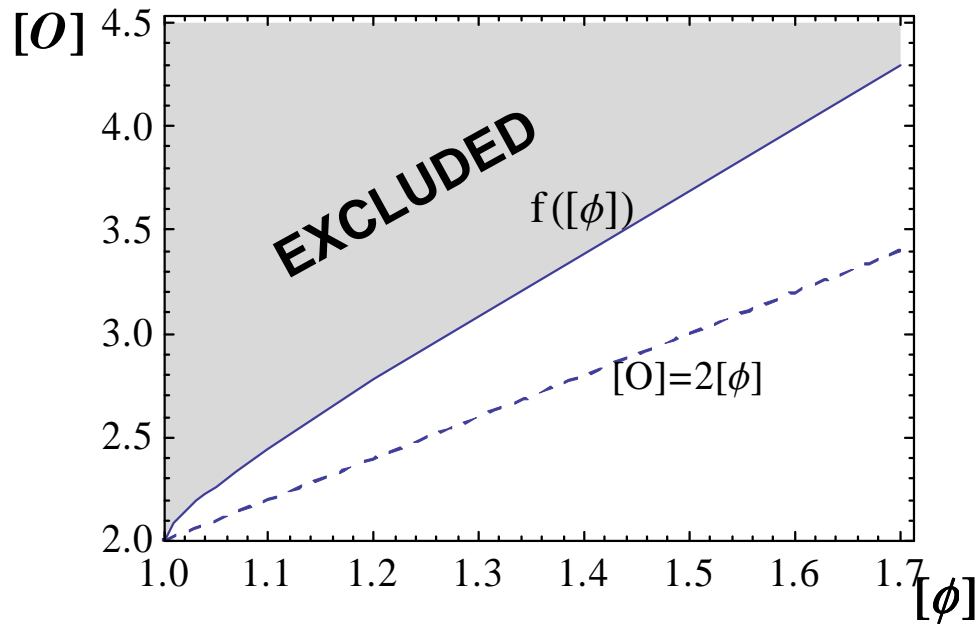
Example:
 $d=1.01$



- **Must sum up to $\equiv 1$ with positive coefficients**
- **$F''(1/2) > 0$ for all higher spins and for all scalars with small Δ**
- **\Rightarrow a scalar with small Δ must be present**

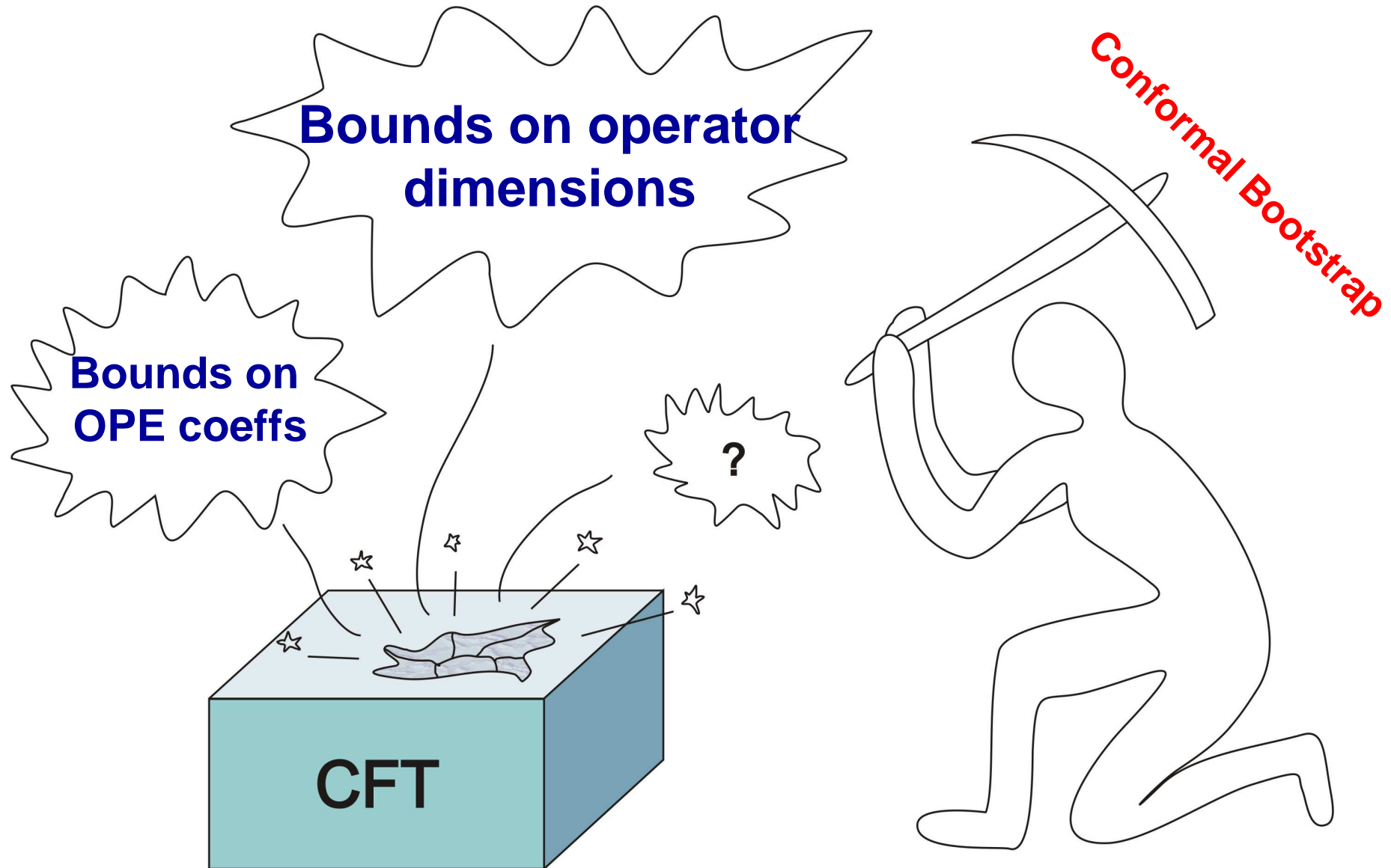


4-D versus 2-D



- In 2D, the bound is saturated by Minimal Models and Free Scalar, interpolating in between
- We hope that also in 4D CFTs near-saturating the bound must exist

Conclusions/Подведем итоги



OPE

Primaries

Descendants
(fixed by symmetry)

$$\phi(x)\phi(0) \sim \frac{1}{|x|^{2d}} \left\{ \mathbf{1} + \sum_{l=2n} c_{\Delta,l} \left[|x|^{-\Delta} K_l(x) \cdot O_{\Delta,l}(0) + \dots \right] \right\}$$

$d \equiv [\varphi]$

by Bose symmetry

$$K_l(x) = \frac{x^{\mu_1} \dots x^{\mu_l}}{|x|^l}$$

$\Delta \geq 1$ ($l=0$)
 $\Delta \geq l+2$ ($l=2,4,6\dots$)
Unitarity bounds
Mack'77

2D and 3D examples

show that $\gamma_{\phi^2} \gg \gamma_{\phi}$ is not impossible.

Ising model: $\sigma \times \sigma = 1 + \varepsilon$

2-dimensions (Onsager)	$[\sigma] = 1/8, \quad [\varepsilon] = 1$
3-dimensions (ε - and high-T expansions, Monte-Carlo)	$\gamma_{\sigma} \approx 0.02, \quad \gamma_{\varepsilon} \approx 0.4$

OPE

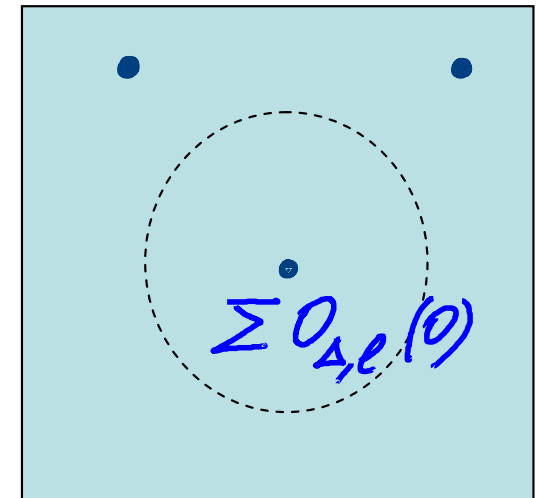
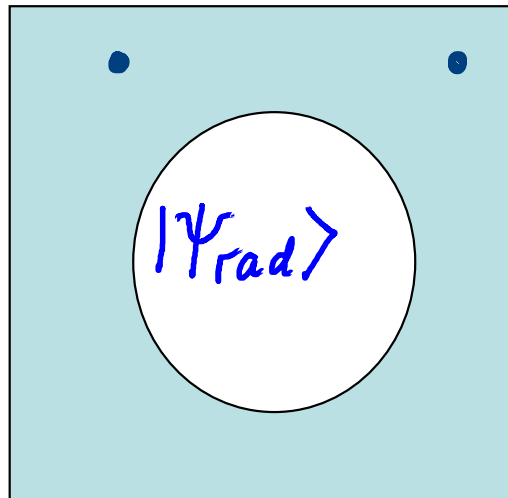
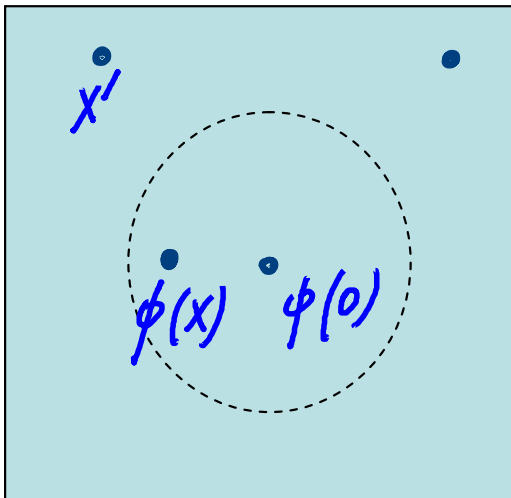
$$\phi(x)\phi(0) = \frac{1}{x^{2d}} + \sum_{l=0,2,4,\dots} \sum_{\Delta} \lambda_{\Delta,l} [C(x)O(0) + \dots]$$

descendants

NB. Only even spins appear

Convergence at finite separation:

should converge if no other operators with $|x'| < |x|$

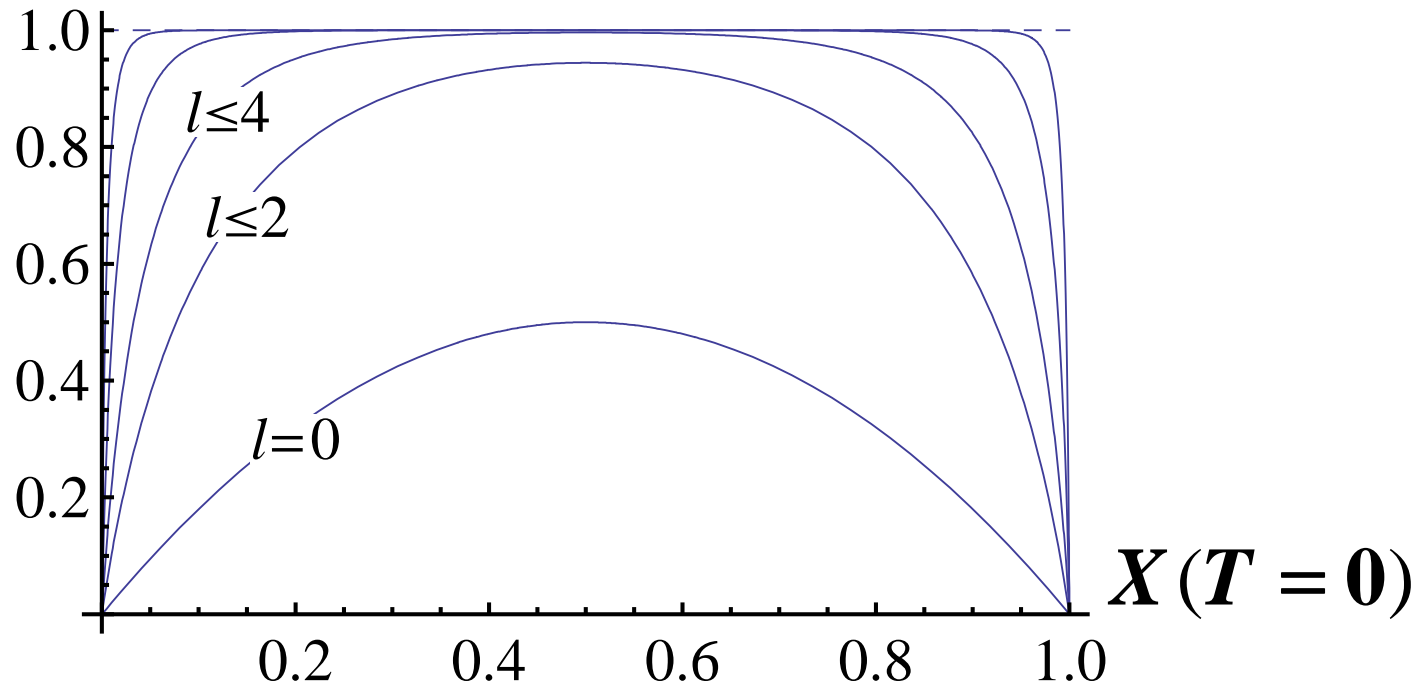


Sum rule convergence in free scalar theory

$$\phi \times \phi = \sum_{l=2n} \phi \overset{\leftrightarrow}{\partial}{}^{2n} \phi$$

twist 2 fields only

$$\lambda_l^2 = 2^{l+1} \frac{(l!)^2}{(2l!)^2}$$



Monotonic convergence

-Paradox in 4-epsilon dimensions

Naive extrapolation of our 4d bound:

$$\gamma_{\phi^2} \leq 0.7 \sqrt{\gamma_{\phi}} \quad (\gamma_{\phi} \ll 1)$$

to 4-epsilon is in contradiction with Wilson-Fischer fixed point anomalous dimensions for $N=1,2$:

$$\gamma_{\phi} = \frac{N+2}{4(N+8)^2} \varepsilon^2$$
$$\gamma_{\phi^2} \equiv \gamma_T = \frac{2}{N+8} \varepsilon$$