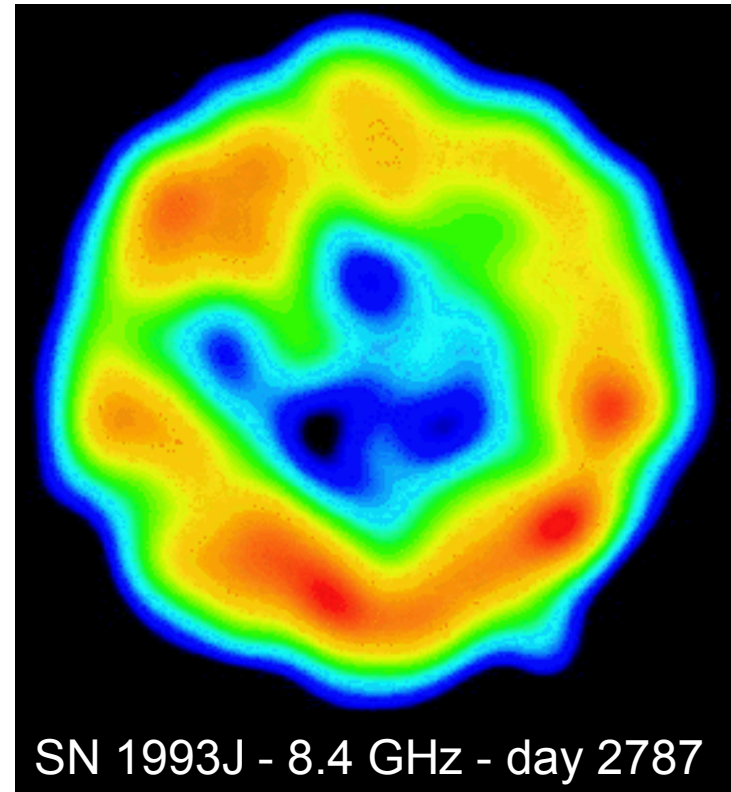
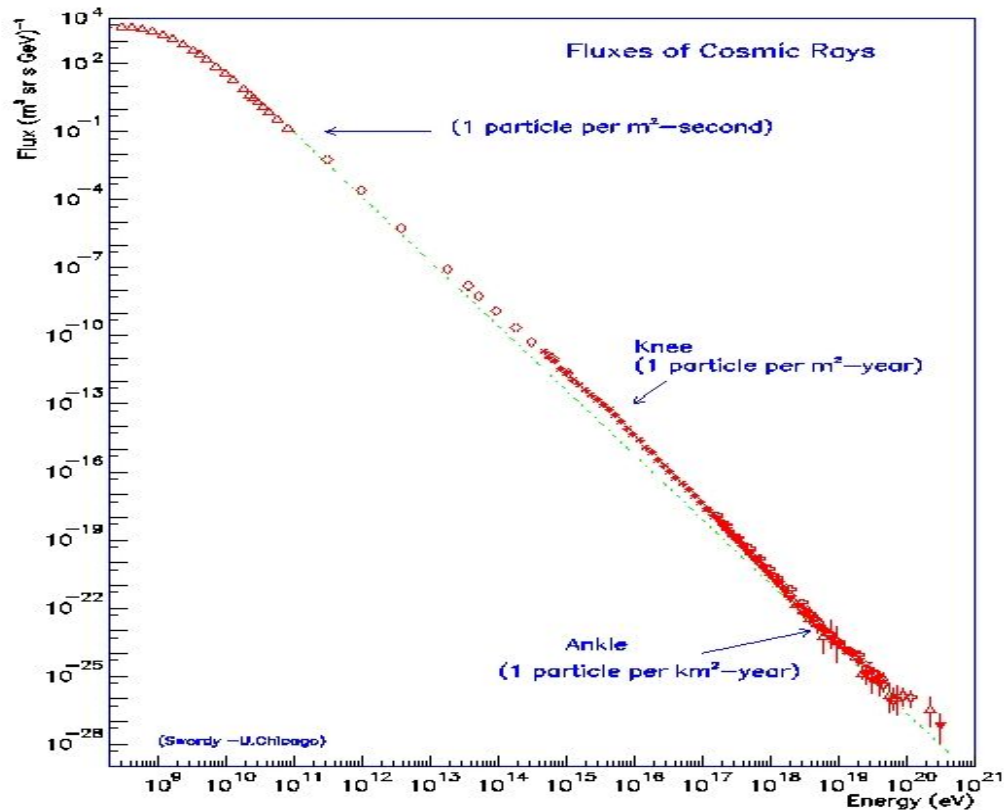


Cosmic-ray acceleration in radio supernovae

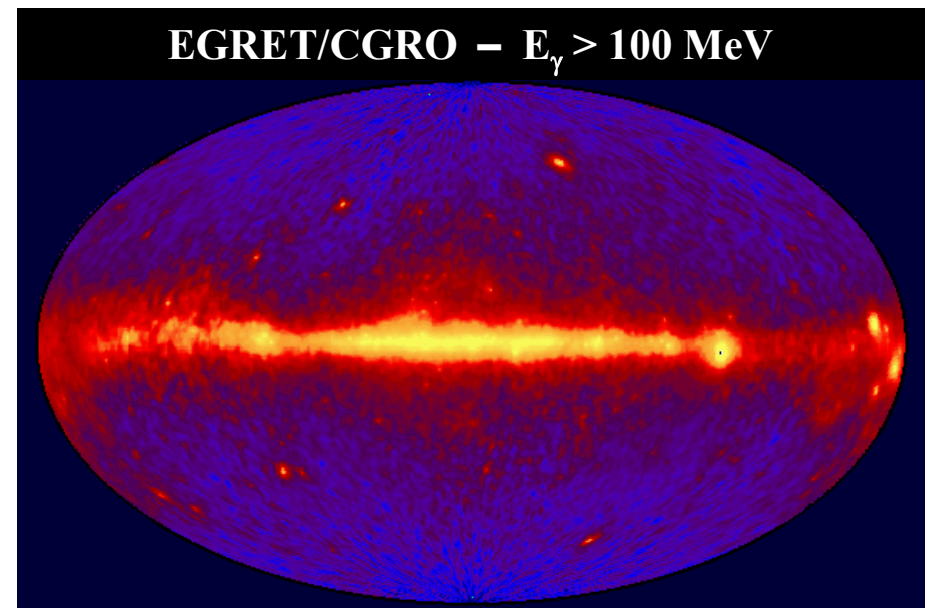
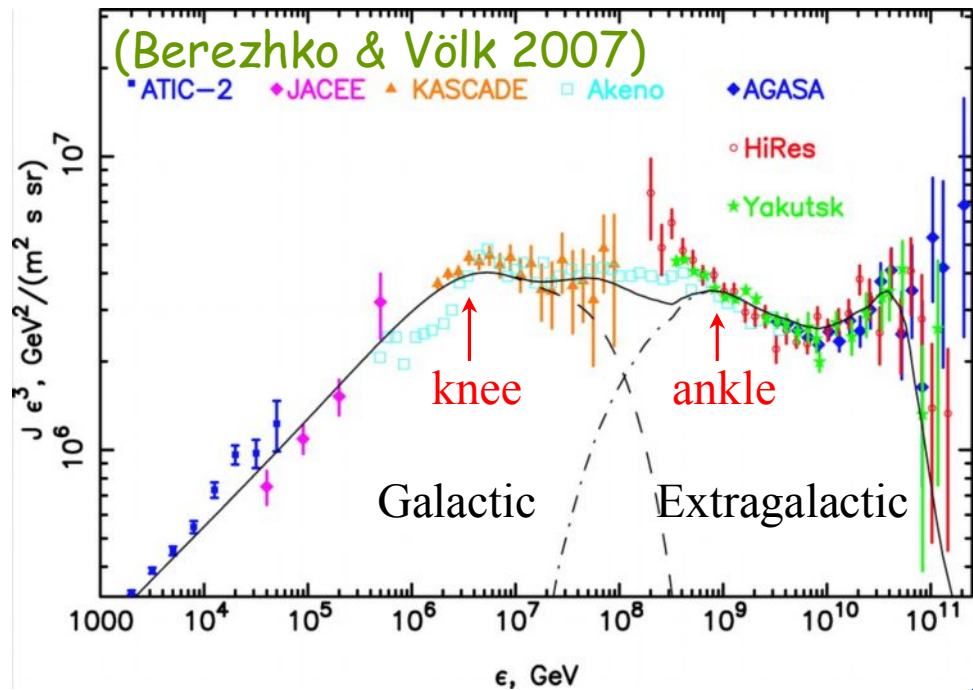
V. Tatischeff

CSNSM, Orsay, France



4th Sakharov Conference, Lebedev Institute, Moscow, May 18-23, 2009

Galactic cosmic-ray and supernova energetics



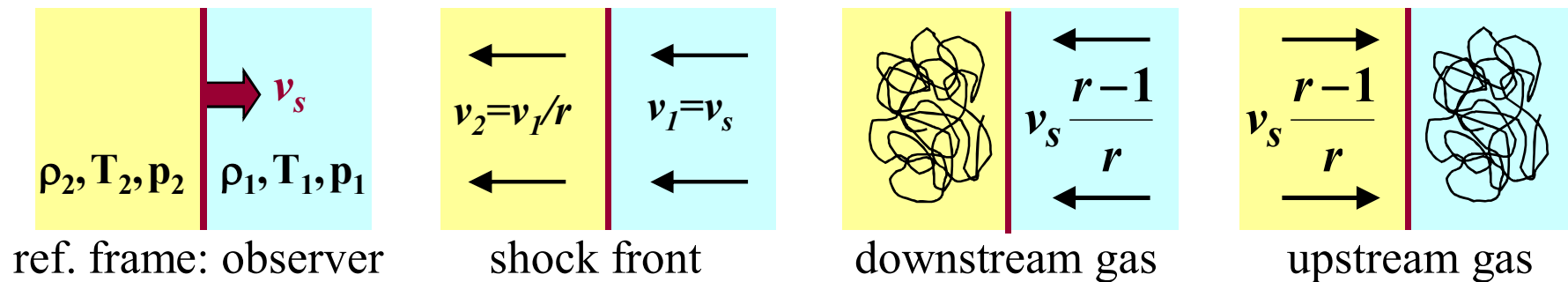
- Galactic origin below $\sim 10^{16} - 10^{18}$ eV
- Extrag. from AGNs (Auger 2007)

- Galactic diffuse γ -ray emission:
CR + ISM \rightarrow pion $\rightarrow \gamma$
- CR source luminosity: 5×10^{40} erg/s

- Total power supplied by SNe: $L_{SN} \sim 1.5 \times 10^{51}$ erg \times 50 yr $^{-1} \sim 10^{42}$ erg/s
 \Rightarrow SN acceleration efficiency $\sim 5\%$

Diffusive shock acceleration

- First-order Fermi (1949) acceleration process in SN shock waves (Krymskii 1977; Bell 1978; Axford et al. 1978; Blandford & Ostriker 1978)
- Particle diffusion on **magnetic turbulences** on both sides of the shock

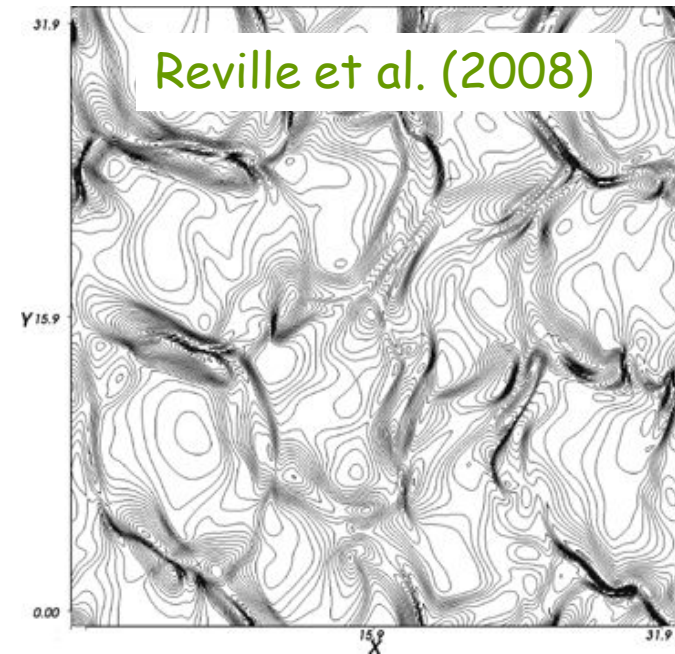
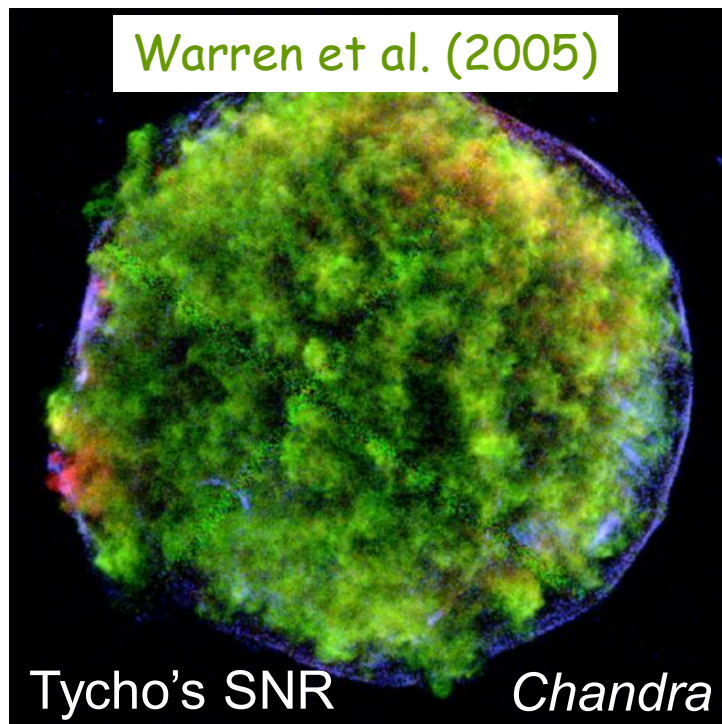


- Fractional momentum gain per cycle: $\frac{\Delta p}{p} = \frac{4}{3} \frac{r-1}{r} \frac{v_s}{c}$ (relativistic)
- Particle energy spectrum: $N(E) \propto E^{-q}$ with $q = 1 + 3 / (r-1)$ (for a test-particle strong shock $r = 4 \Rightarrow q = 2$)
- Relatively slow process $\Rightarrow E_{\max} = 23 \text{ TeV} \frac{Z B_{\mu\text{G}} E_{51}}{n_0^{1/3} M_{\text{ej}}^{1/6}}$ (Lagage & Cesarky 1983)

$\Rightarrow E_{p,\max} \ll 3 \times 10^{15} \text{ eV}$ for $B \sim \text{few } \mu\text{G} !$

Magnetic field amplification

- CRs can excite magnetic fluctuations in the upstream plasma by both resonant (Bell and Lucek 2001) and nonresonant (Bell 2004) streaming instabilities: $\delta B/B \gg 1$
- MHD simulations: **yes** (e.g. Zirakashvili et al. 2008) or **no** (Niemi et al. 2008; $\delta B/B \sim 1$)

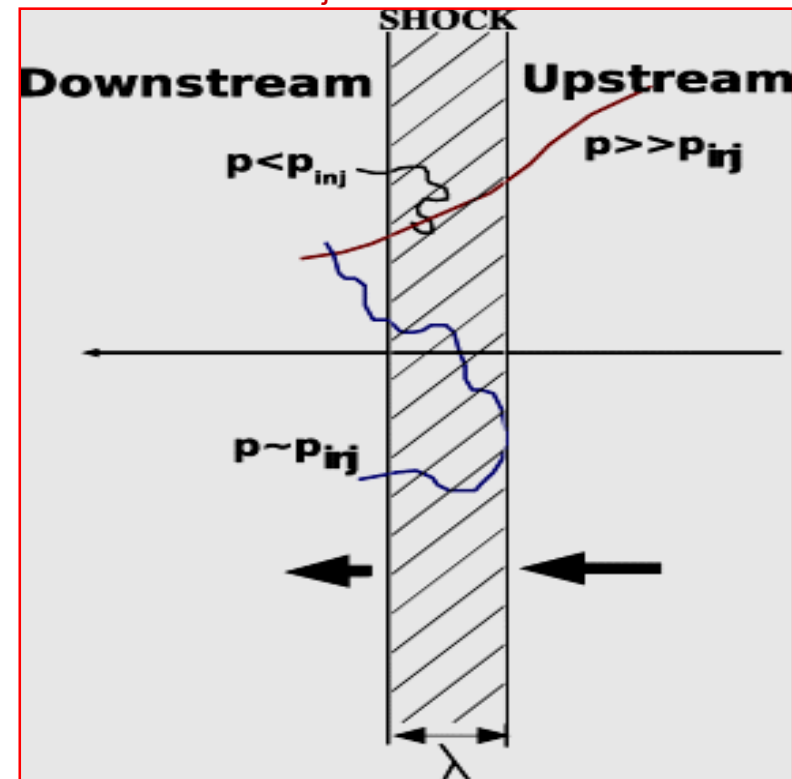
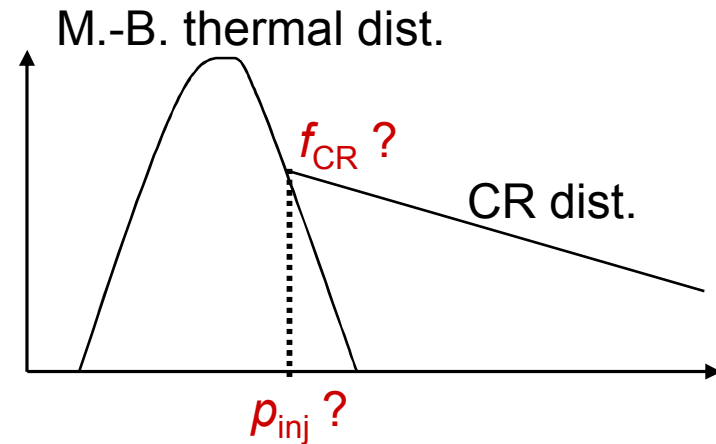


- Evidence for B-field amplification in SNRs ($\delta B/B > 20$) from synchrotron **X-ray filaments** (e.g. Parizot et al. 2006) (but δB damping? Pohl et al. 2005)

\Rightarrow Acceleration to $\approx 10^{15}$ eV. And beyond?

The injection problem

- Particles injected into the DSA process: **postshock thermal particles with $p > p_{inj}$** (Ellison et al.; Monte-Carlo simulations)
 - Condition for injection (Blasi et al. 2005): $r_L > \lambda = \alpha r_L^{th}$, with the shock thickness parameter $\alpha = 1 - 2$
 - But depending on the shock strength and α , the fraction of thermal particles converted into CRs $\eta_{inj} \in [10^{-5}, 10^{-2}]$!
- ⇒ **Acceleration efficiency ?**



Cosmic-ray modified shock

See Berezhko & Ellison (1999)

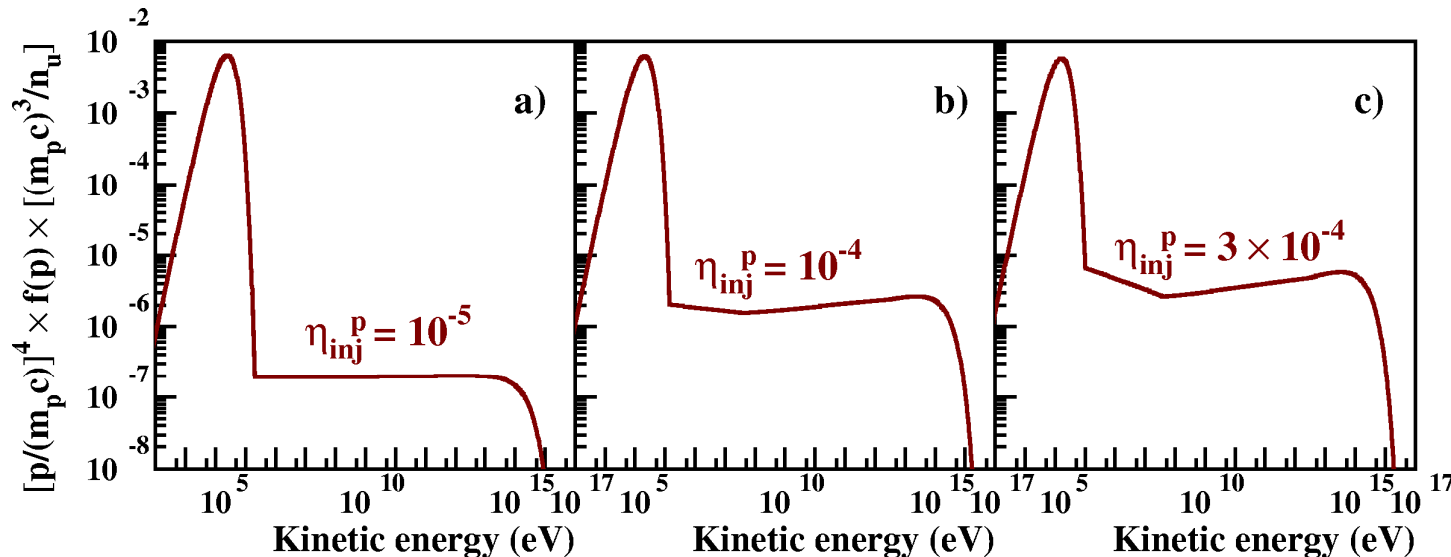
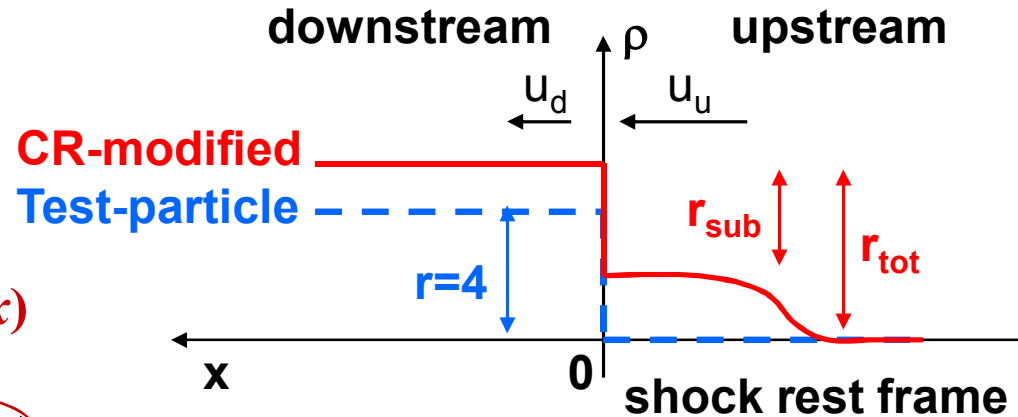
- Momentum flux conservation:

$$\rho_0 u_0^2 + P_{g,0} = \rho(x) u(x)^2 + P_g(x) + P_{CR}(x)$$

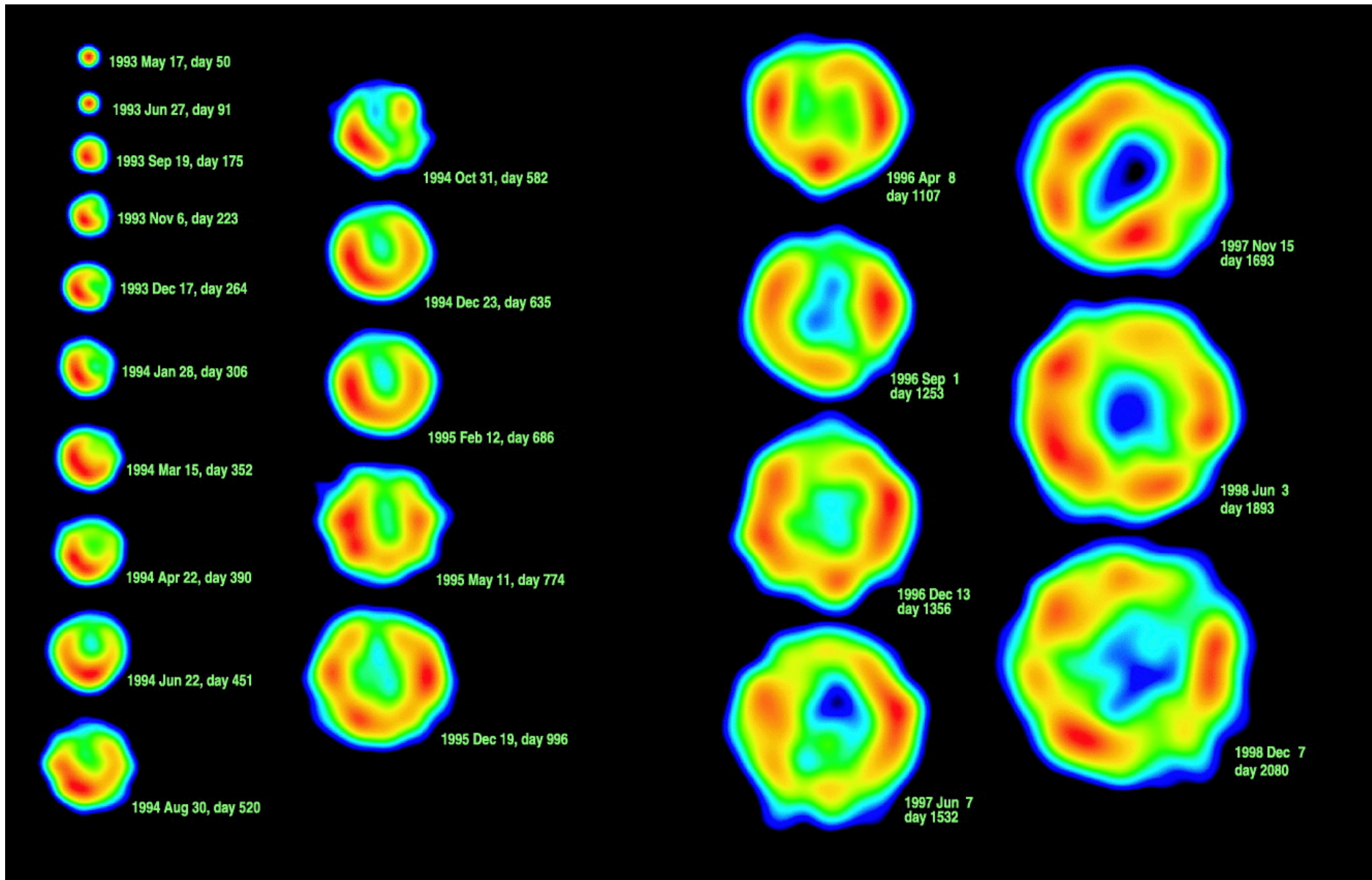
with $P_{CR}(x) = \frac{1}{3} \int_{p_{inj}}^{p_{max}} dp 4\pi p^3 v(p) f(x, p)$

particle dist. funct. (\Leftarrow diffusive transport eq.)

- Higher energy particles feel a higher compression ratio \Rightarrow concave spect.

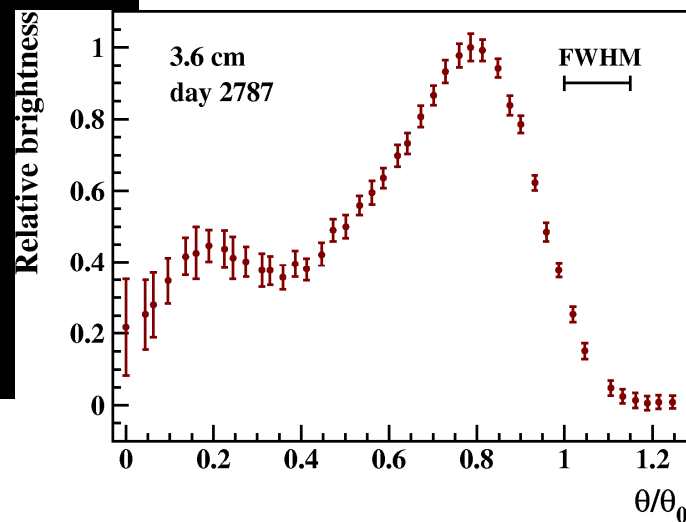
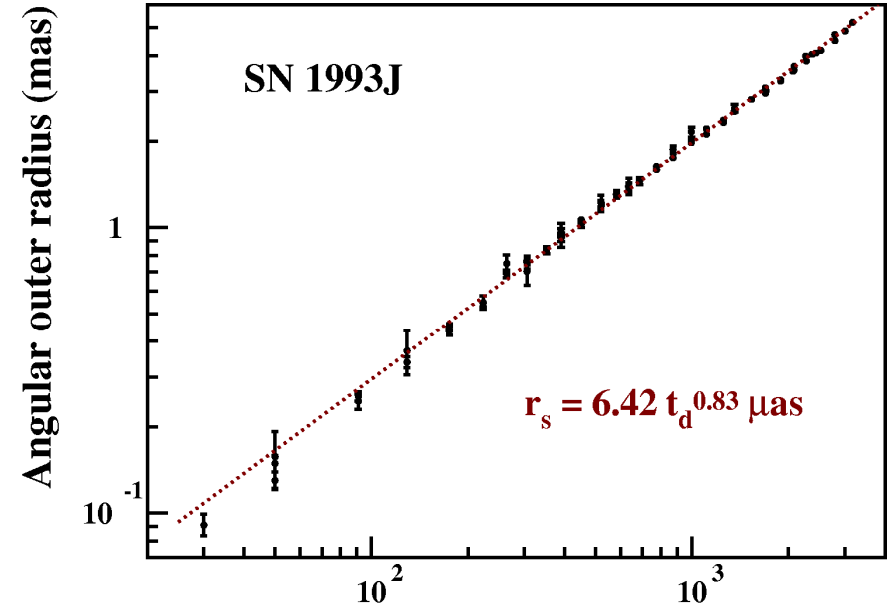
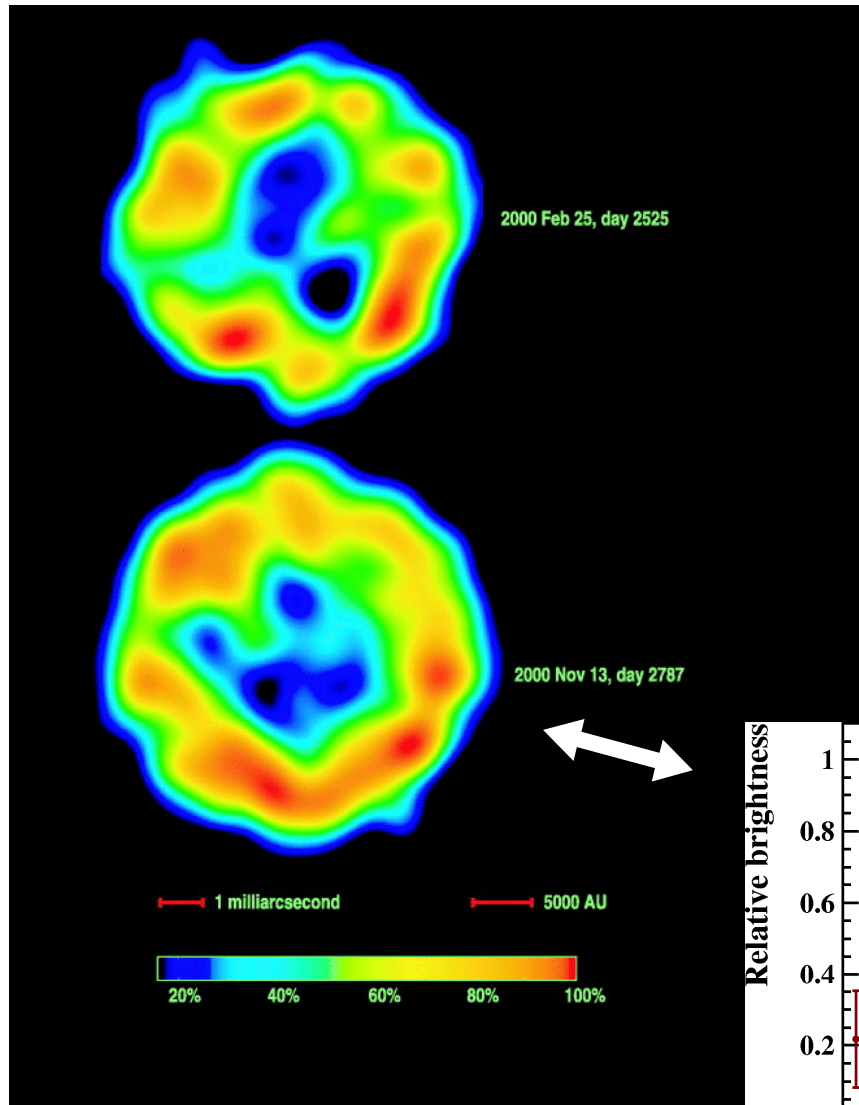


SN 1993J — One of the best observed radio SN



SN 1993J VLBI observations

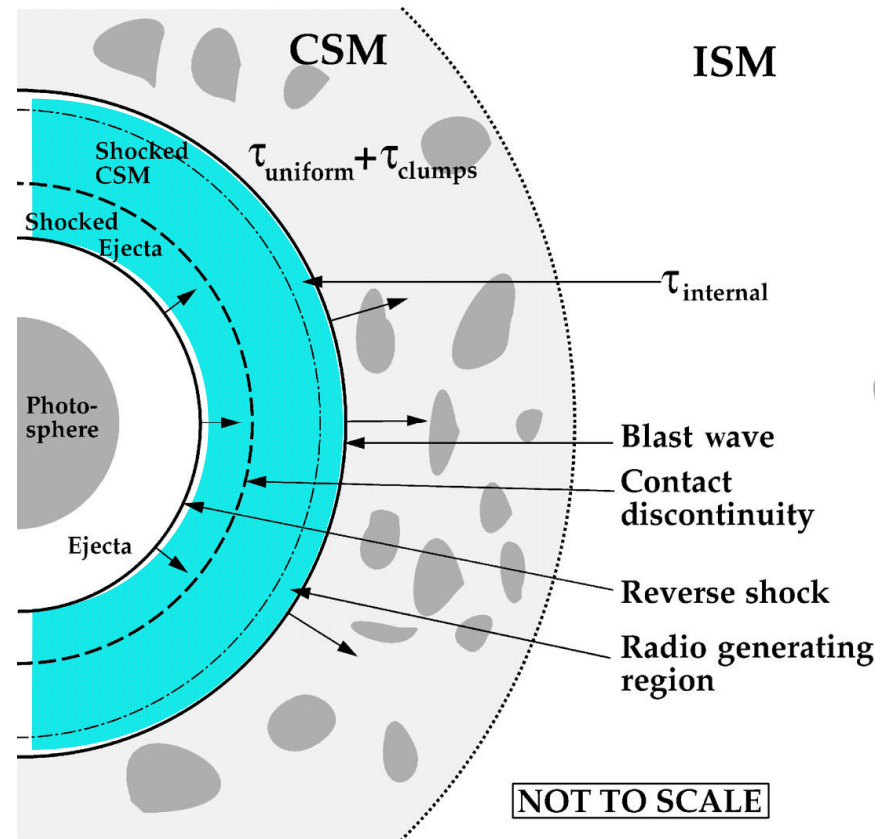
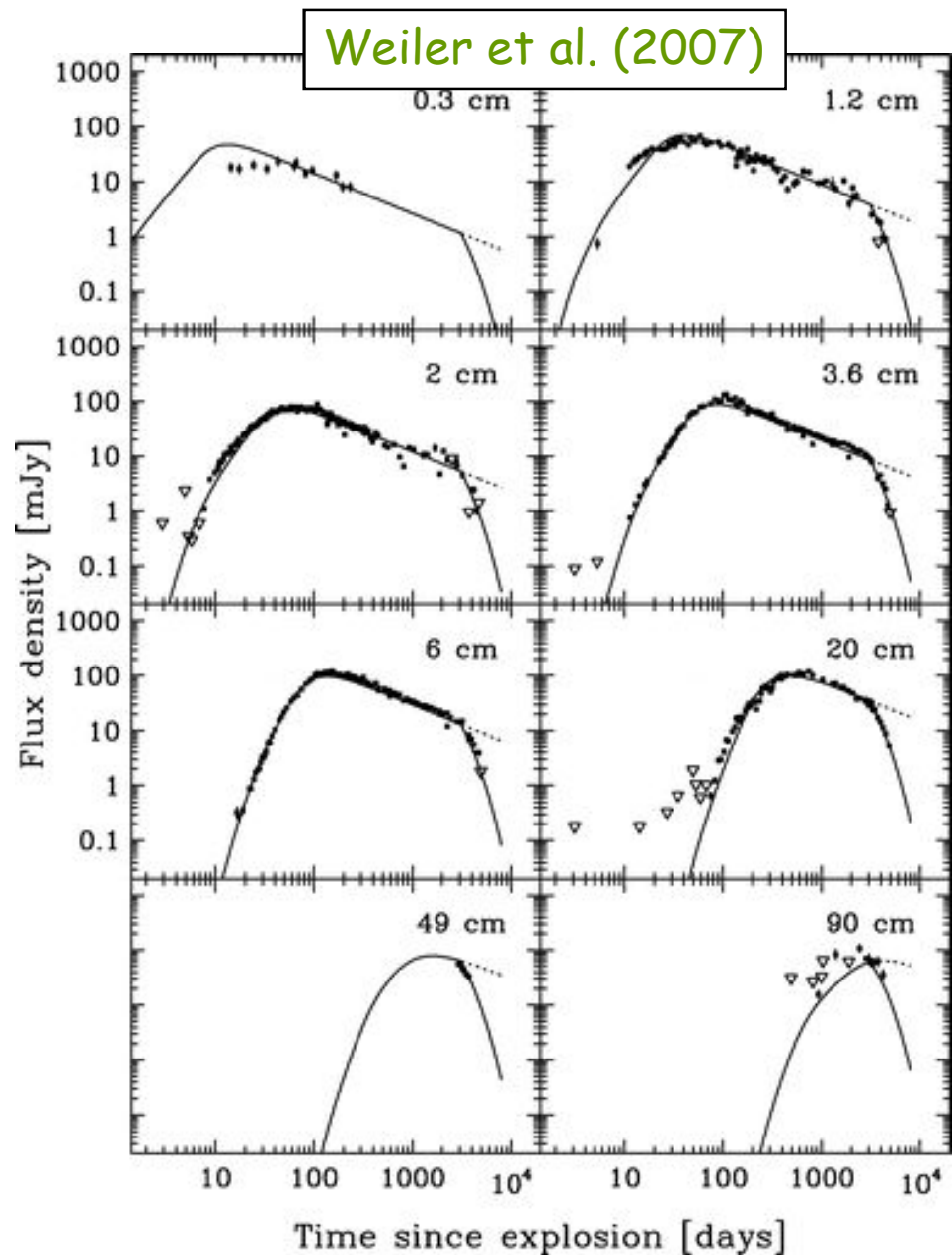
- The expansion is almost self similar



Days after outburst

See, e.g.,
Marcaide et al.
(1995, 1997)
Bartel et al.
(2002, 2007)

Radio light curves



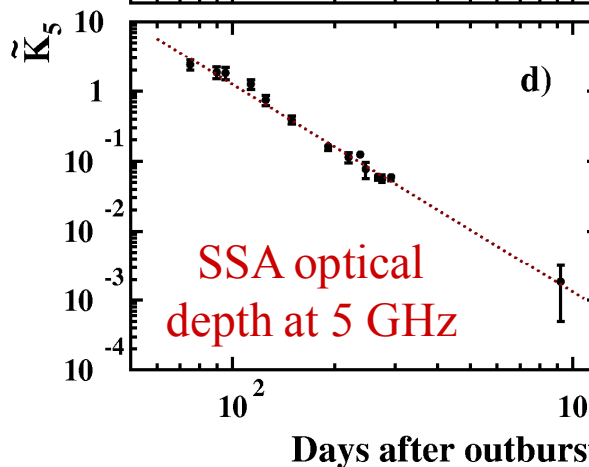
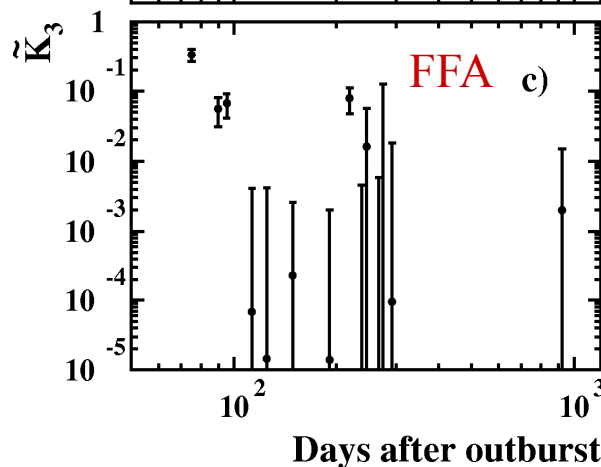
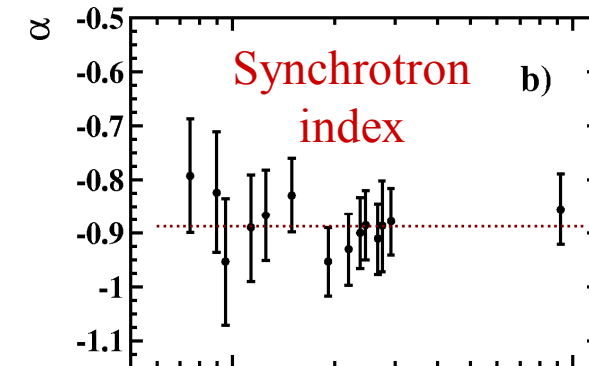
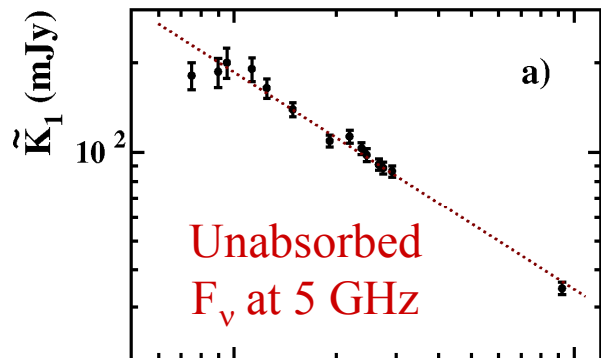
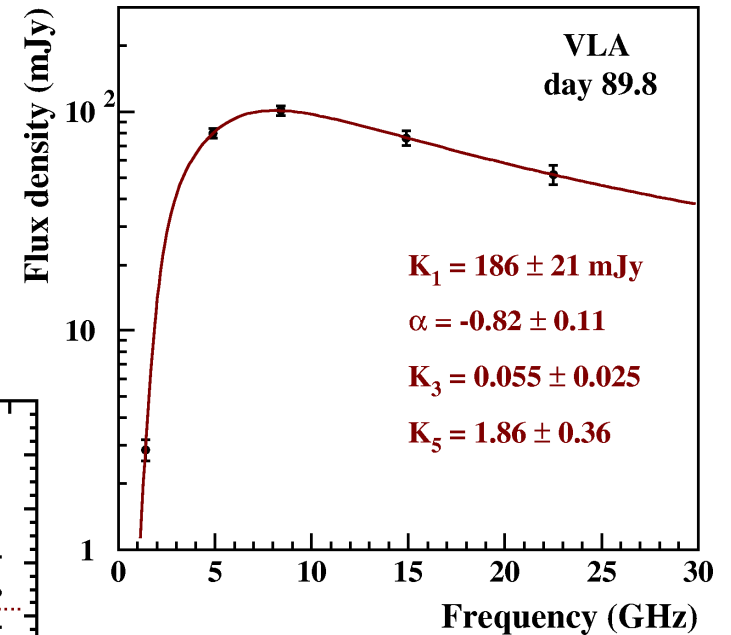
- **Abrupt decline after day ~ 3100** as the shock has reached the outer limit of the dense CSM
- **Weiler et al.:** overall fit to all the data with 9 free parameters

Synchrotron self-absorption (SSA)

- Assuming a **power-law** distribution of electrons radiating in a **uniform shell**:

$$\langle B \rangle \propto_e \left(\frac{K_5}{K_1} \right)^2$$

- Fits to individual VLA spectra



$$\langle B \rangle = \left(2.4 \pm 1.0 \right) \left(\frac{t}{100 \text{ days}} \right)^b \text{ G}$$

with $b = -1.16 \pm 0.20$

\Rightarrow Strong amplification of the stellar B-field

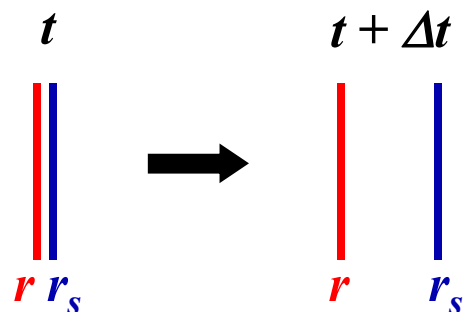
Radio SN model (1)

- Inspired by *Cassam-Chenaï et al. (2005)* for Galactic SNRs

- Nonlinear diffusive shock acceleration model

(*Berezhko & Ellison 1999*) $\Rightarrow f_p(p, r_s)$ and $f_e(p, r_s)$

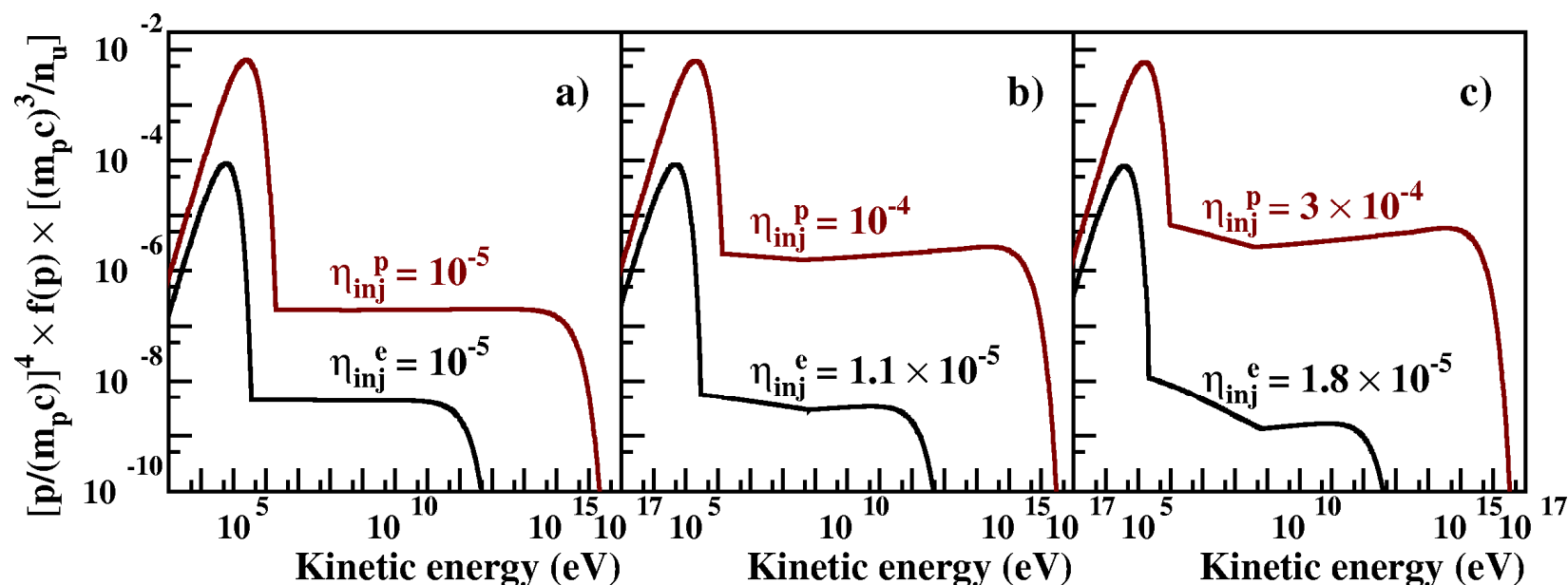
f_e depends on the proton injection parameter η_{inj}^p



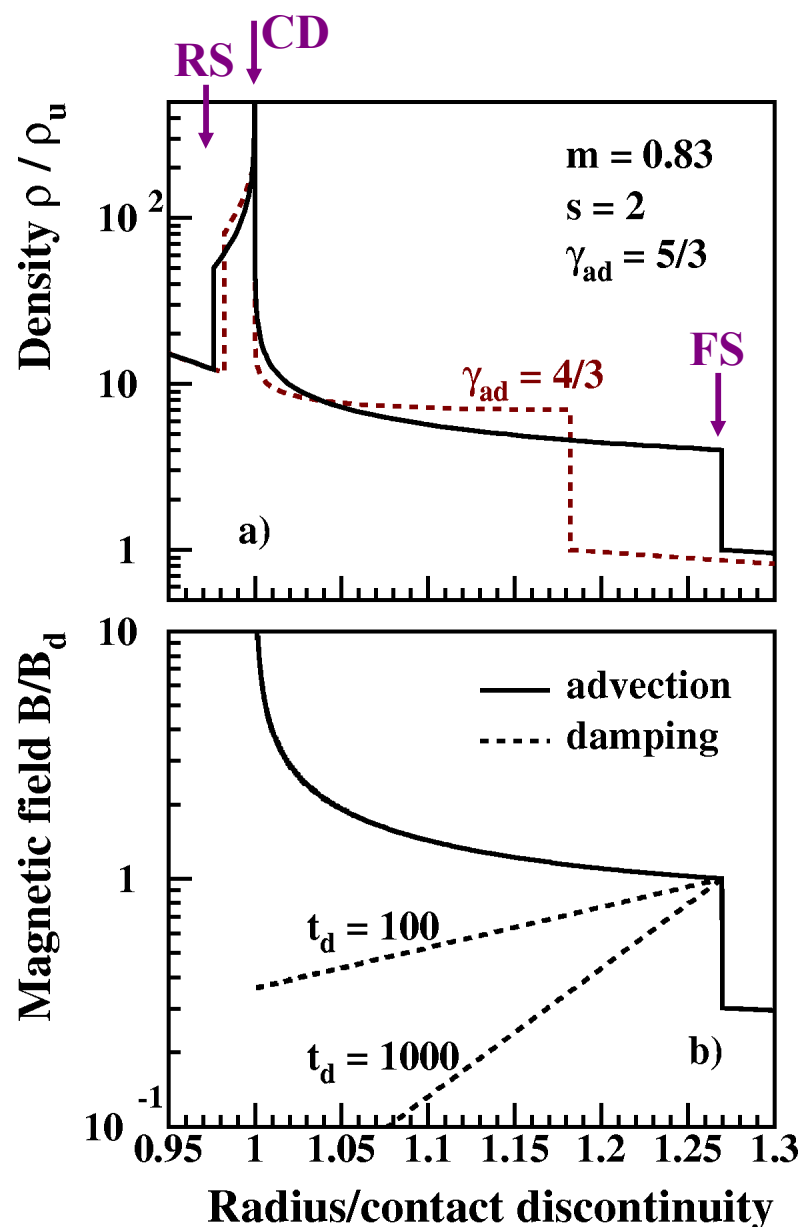
- Electron energy losses during the expansion:

adiabatic, synchrotron and inverse Compton cooling (*Reynolds 1998*)

(radiation density dominated by the ejecta; *Fransson & Björnsson 1998*)

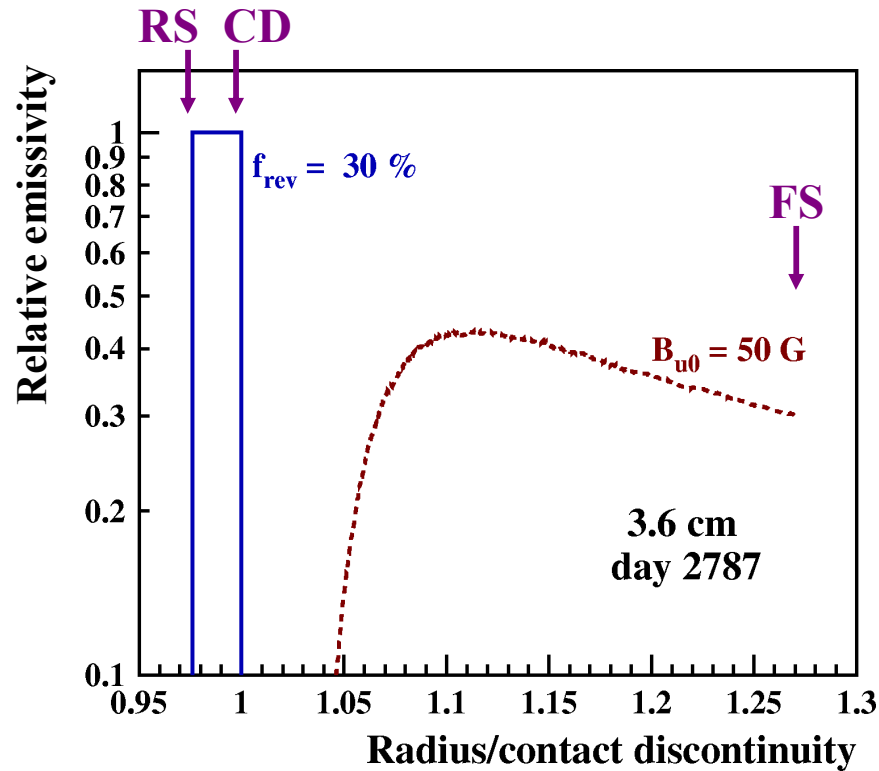


Radio SN model (2)

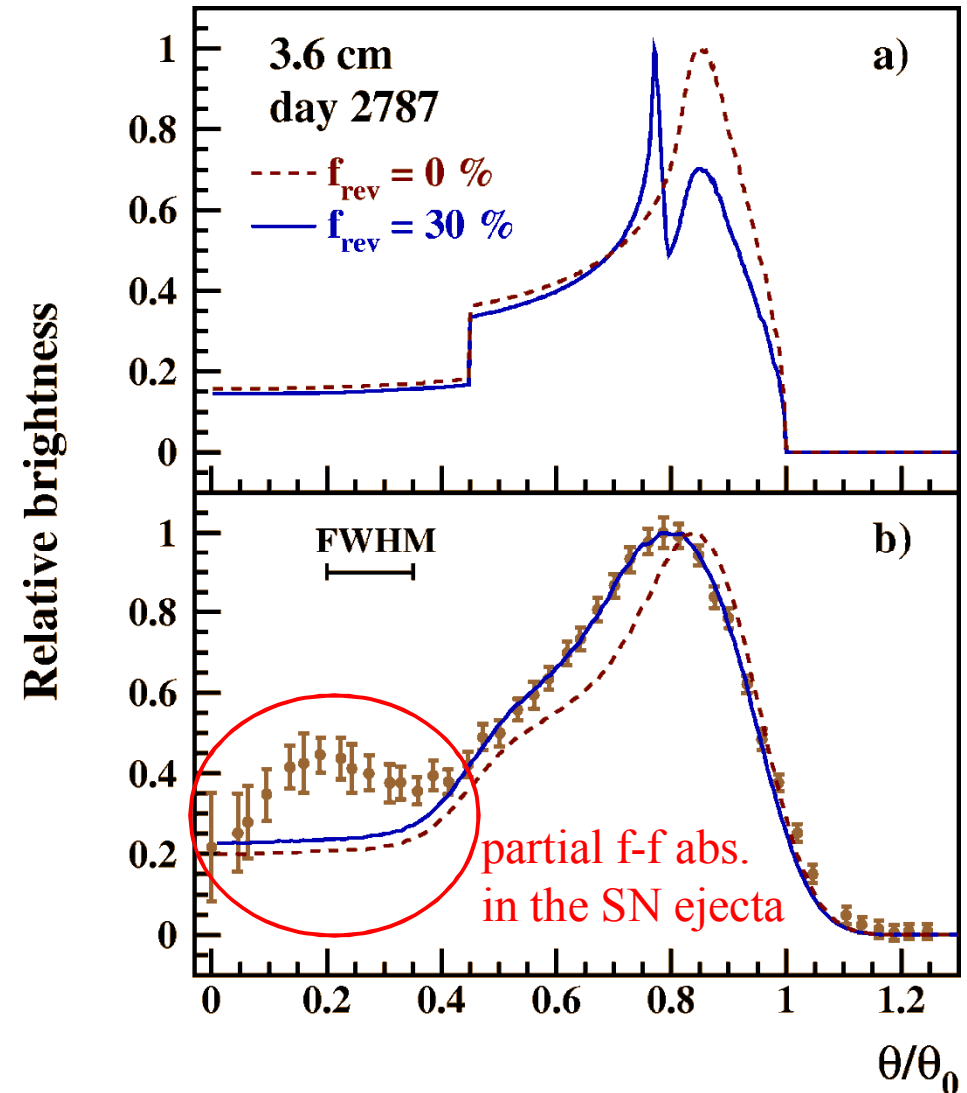


- Hydrodynamic of the postshock plasma: **self-similar solutions** (Chevalier 1983)
 - Postshock magnetic field evolution:
 - **advected** in the plasma flow or
 - **damped** by cascading of MHD wave energy (Pohl et al. 2005)
 - Synchrotron emission: radiative transfer calculations including **synchrotron self-absorption**
 - Free-free absorption in the **clumpy** wind lost from the progenitor star
- ⇒ 4 free parameters:
 $(dM/dt)_{\text{RSG}}$, B_{u0} , η_{inj}^p , and η_{inj}^e

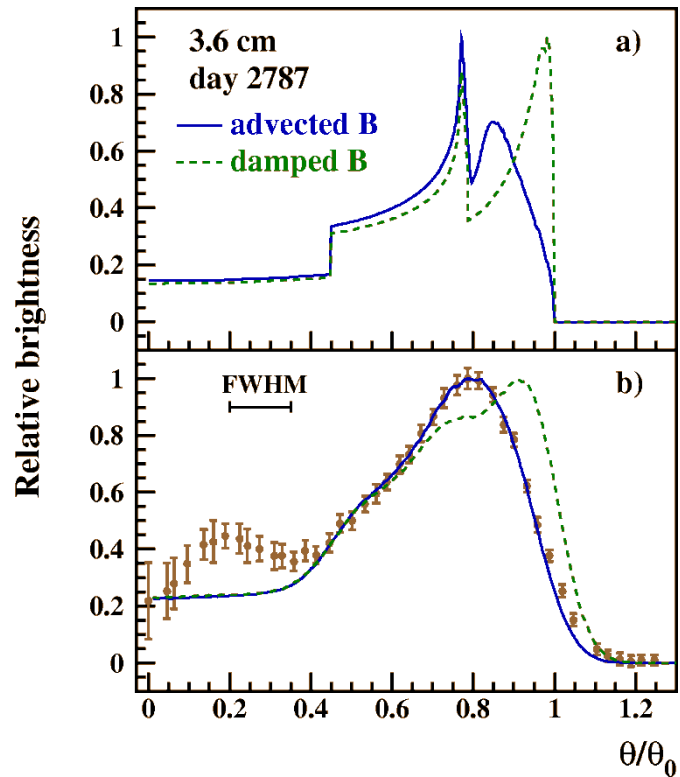
Acceleration at the reverse shock



- Electrons **accelerated at the reverse shock** account for $\leq 17\%$ of the total radio flux



No damping of the postshock magnetic field

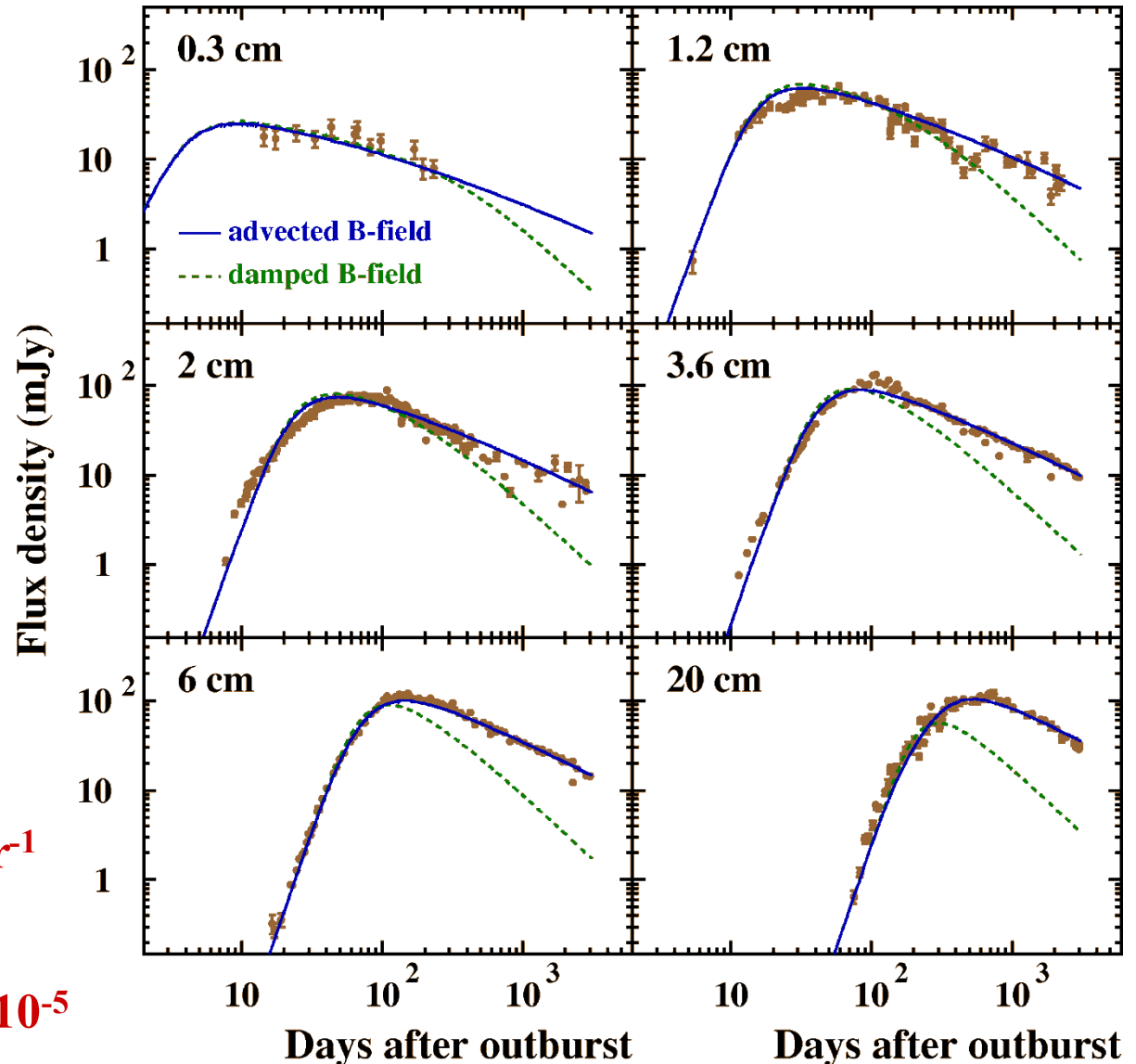


- Best parameter values:

- $(dM/dt)_{\text{RSG}} \approx 3.8 \times 10^{-5} M_{\odot} \text{yr}^{-1}$

- $B_{u0} = 50 \pm 20 \text{ G}$

- $\eta_{\text{inj}}^{\text{p}} \approx 10^{-4}$ and $\eta_{\text{inj}}^{\text{e}} \approx 1.1 \times 10^{-5}$



Magnetic field amplification

- Saturated δB from the **Bell's nonresonant** streaming instability (Pelletier et al. 2006):

$$\frac{\delta B_{\text{nr}}^2}{8\pi} \approx 0.1 \left(\frac{P_{\text{CR}}}{\rho_u v_s^2} \right) \frac{\rho_u v_s^3}{c}$$

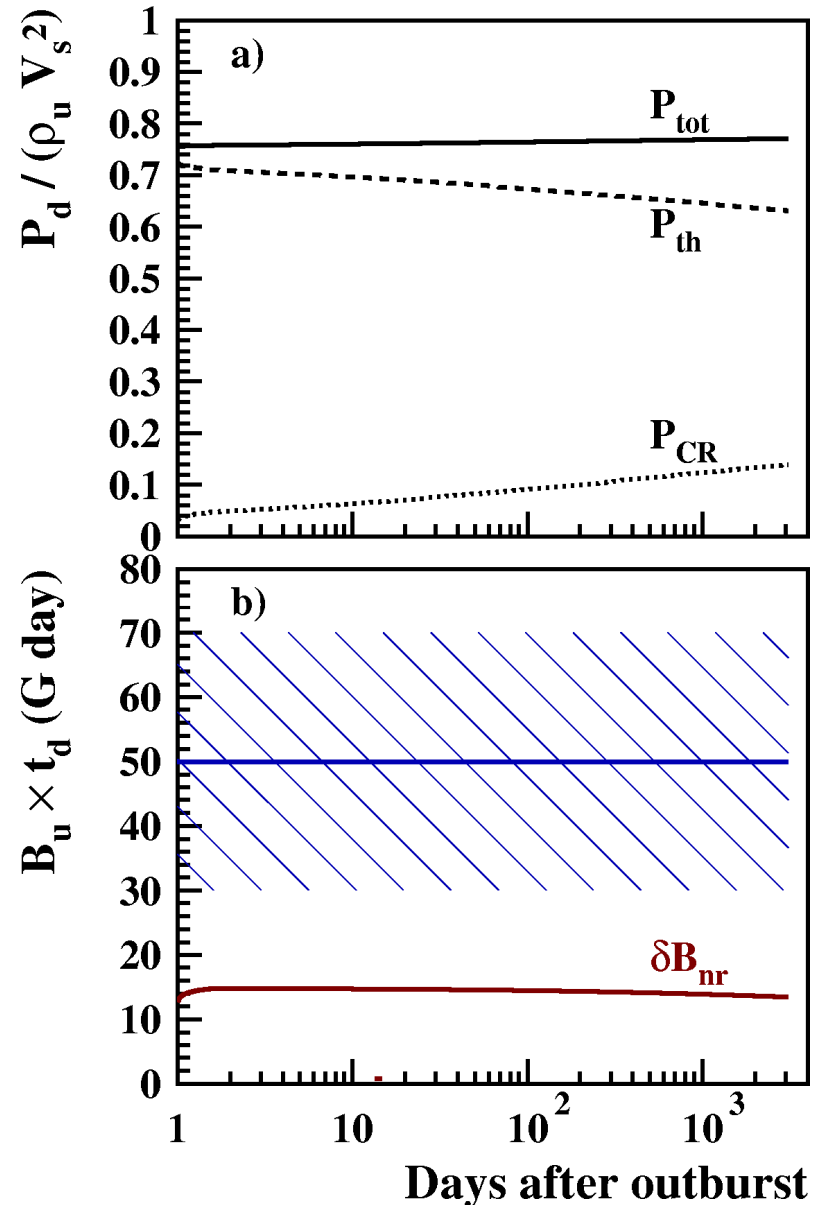
2-5 lower than the "measured" B-field

- Further amplification by the resonant instability (Pelletier et al. 2006)?

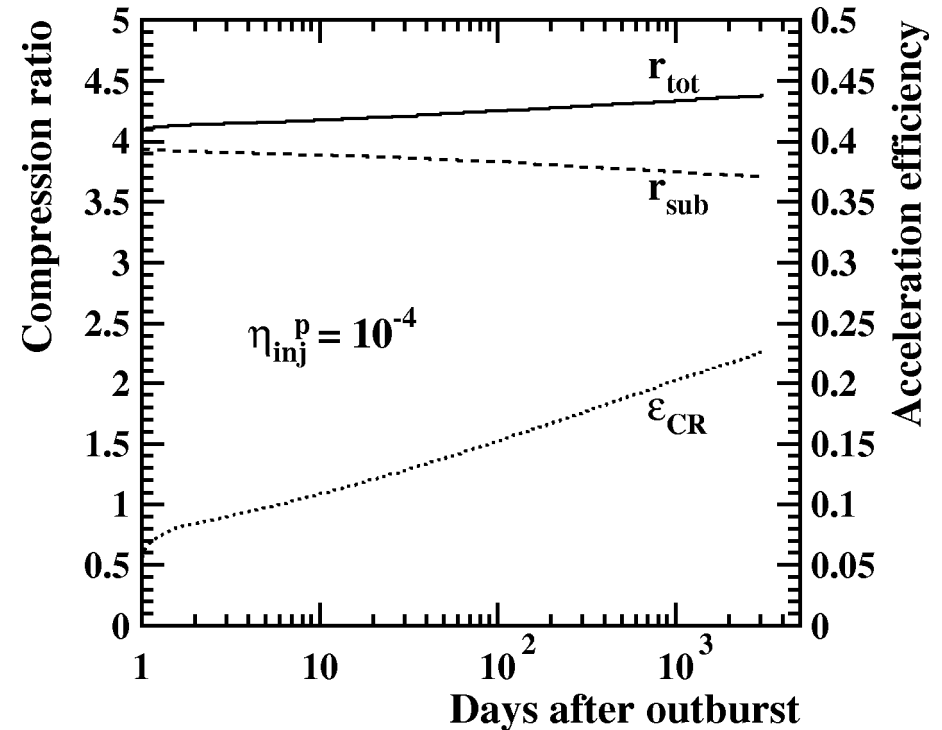
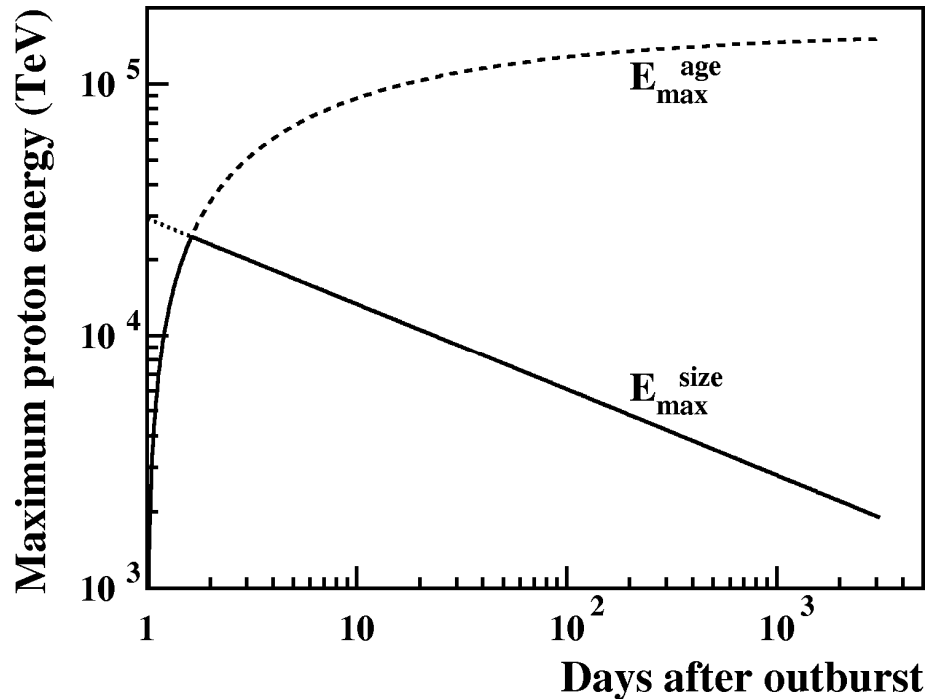
$$\frac{\delta B_{\text{res}}^2}{\delta B_{\text{nr}}^2} \sim \sqrt{\frac{P_{\text{CR}}}{\rho_u v_s^2} \frac{c}{v_s}} \Rightarrow \text{No}$$

- Empirical formula for both **Galactic SNRs** (Berezhko 2008) and **SN 1993J**:

$$\frac{\delta^2}{87} \approx 10^{-1} P_{\text{CR}} \left(\frac{v_s}{3 > 10^4 \text{ km s}^{-1}} \right)$$



SN 1993J and the origin of cosmic rays

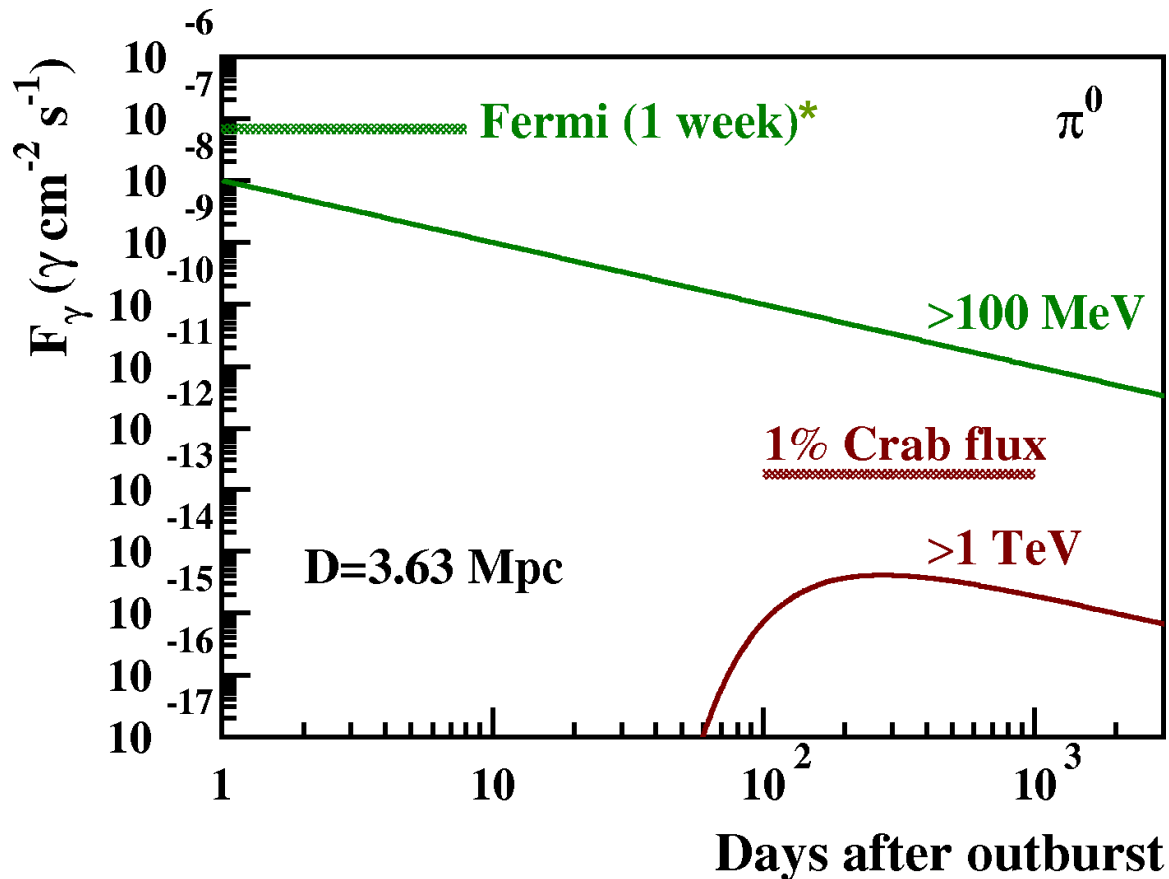


- Rapid acceleration **above the "knee" energy** of 3×10^{15} eV

- Total CR energy: $E_{\text{CR}} \cong \int_{\text{day 1}}^{\text{day 3100}} \varepsilon_{\text{CR}}(t) \times 0.5 \rho_{\text{CSM}} v_s^3 \times 4\pi r_s^2 dt = 7.4 \times 10^{49} \text{ erg}$

- **Escape** of high-energy CRs after day ~ 3100 as $\rho_{\text{CSM}} \searrow \Rightarrow B_u \searrow \Rightarrow l_{\text{diff}} \nearrow$

Gamma-ray emission from π^0 production



* *Fermi* LAT sensitivity for a 5σ detection in all-sky survey operation

The early TeV emission was strongly attenuated by $\gamma + \gamma \rightarrow e^+ + e^-$ in the dense radiation field from the SN ejecta

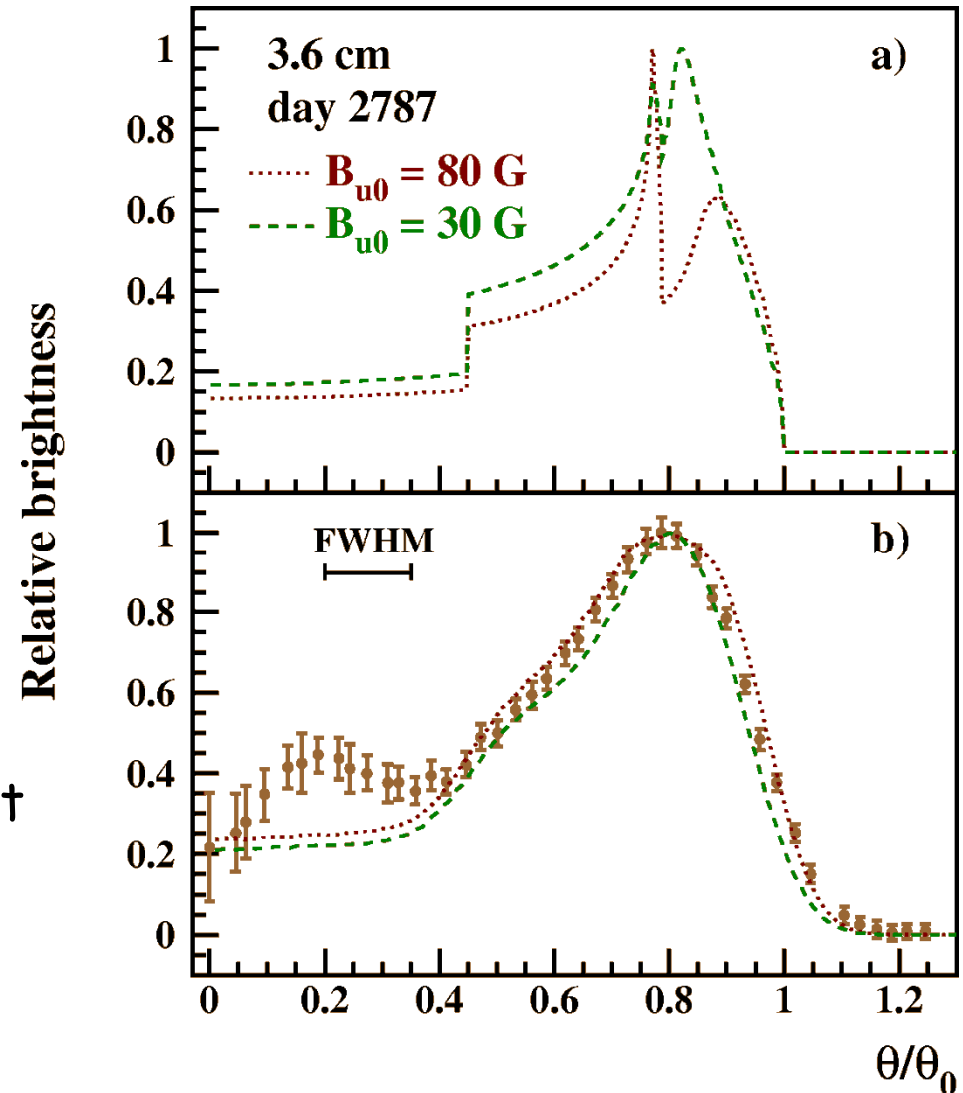
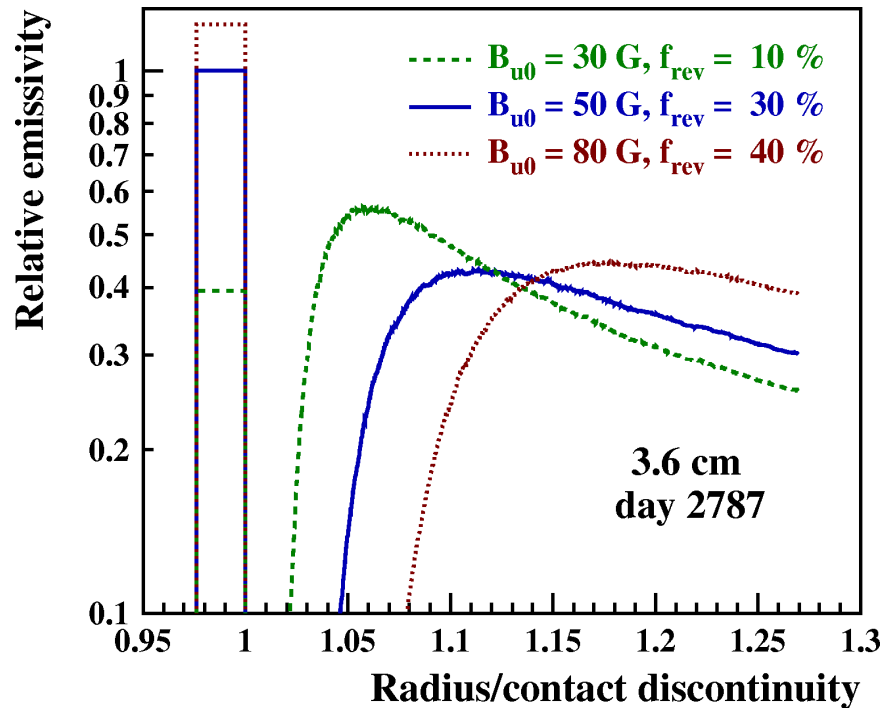
\Rightarrow Type II SNe could be detected in π^0 -decay γ -rays out to a maximum distance of $\approx 1 \text{ Mpc}$

Conclusions

- Evidence from the morphology of the radio emission from SN 1993J that electrons are accelerated at the reverse shock
- The blast wave is a weakly cosmic-ray-modified shock, $\eta_{inj}^p \approx 10^{-4}$
- B-field amplification, possibly by the Bell's nonresonant streaming instability in the precursor region. Consistent with SNR results
- The magnetic turbulence is not damped behind the shock
- Massive stars exploding into their former stellar wind could be a major source of GCRs above $\sim 10^{15}$ eV (Völk & Biermann 1988)
- A new model for radio SNe (e. g. SN 2008D...)
- Type II SNe could be detected at γ -ray energies out to only ≈ 1 Mpc

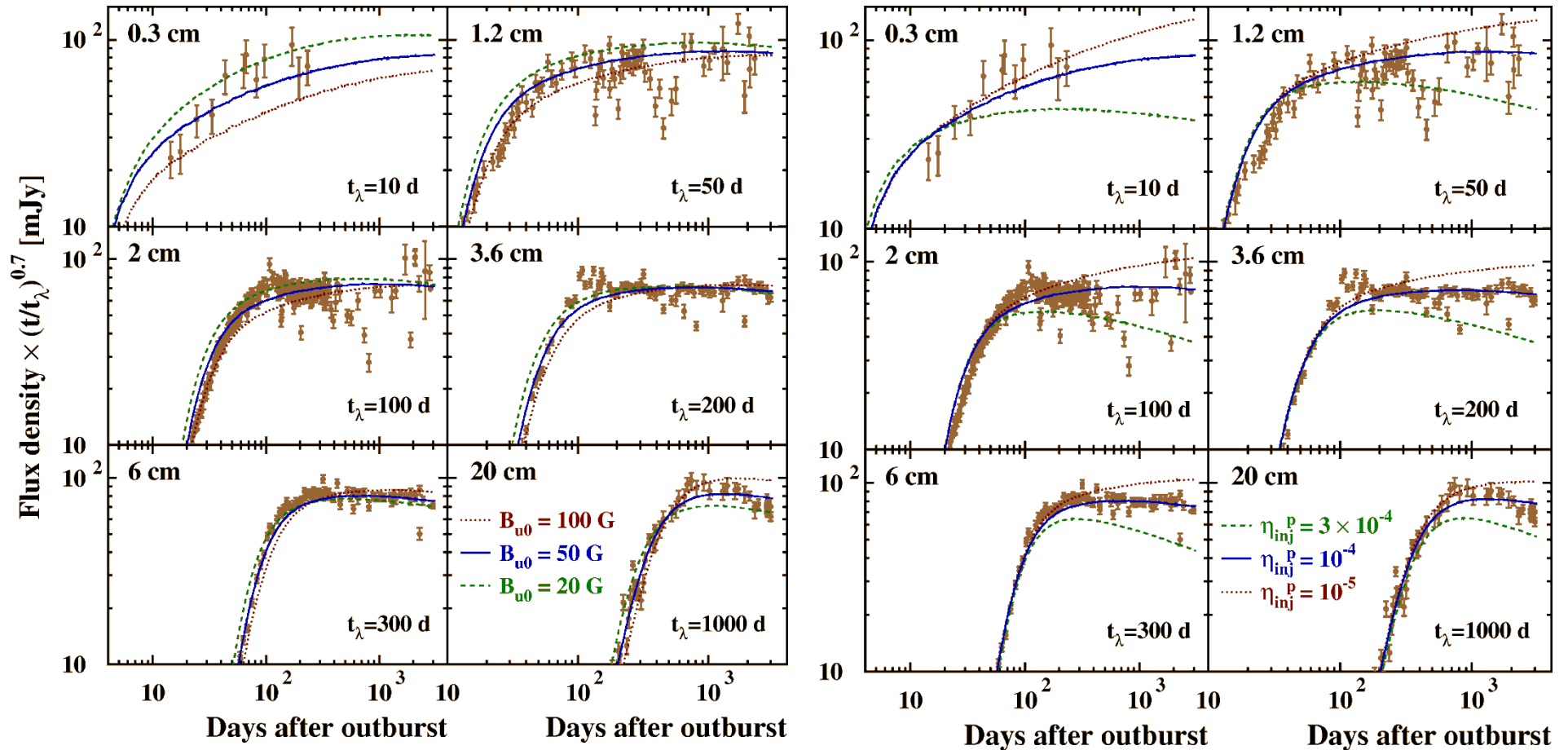
Refs: Tatischeff (2008) PoS [arXiv:0804.1004]; A&A (2009) in press [arXiv:0903.2944]

Radio profile — B-field estimate



- Synchrotron losses are important to the radial emissivity
- From the width of the observed radial profile: $30 < B_{u0} < 80 \text{ G}$

Parameter uncertainties



- The **degeneracy** between B and N_e is **lifted** by the synchrotron losses

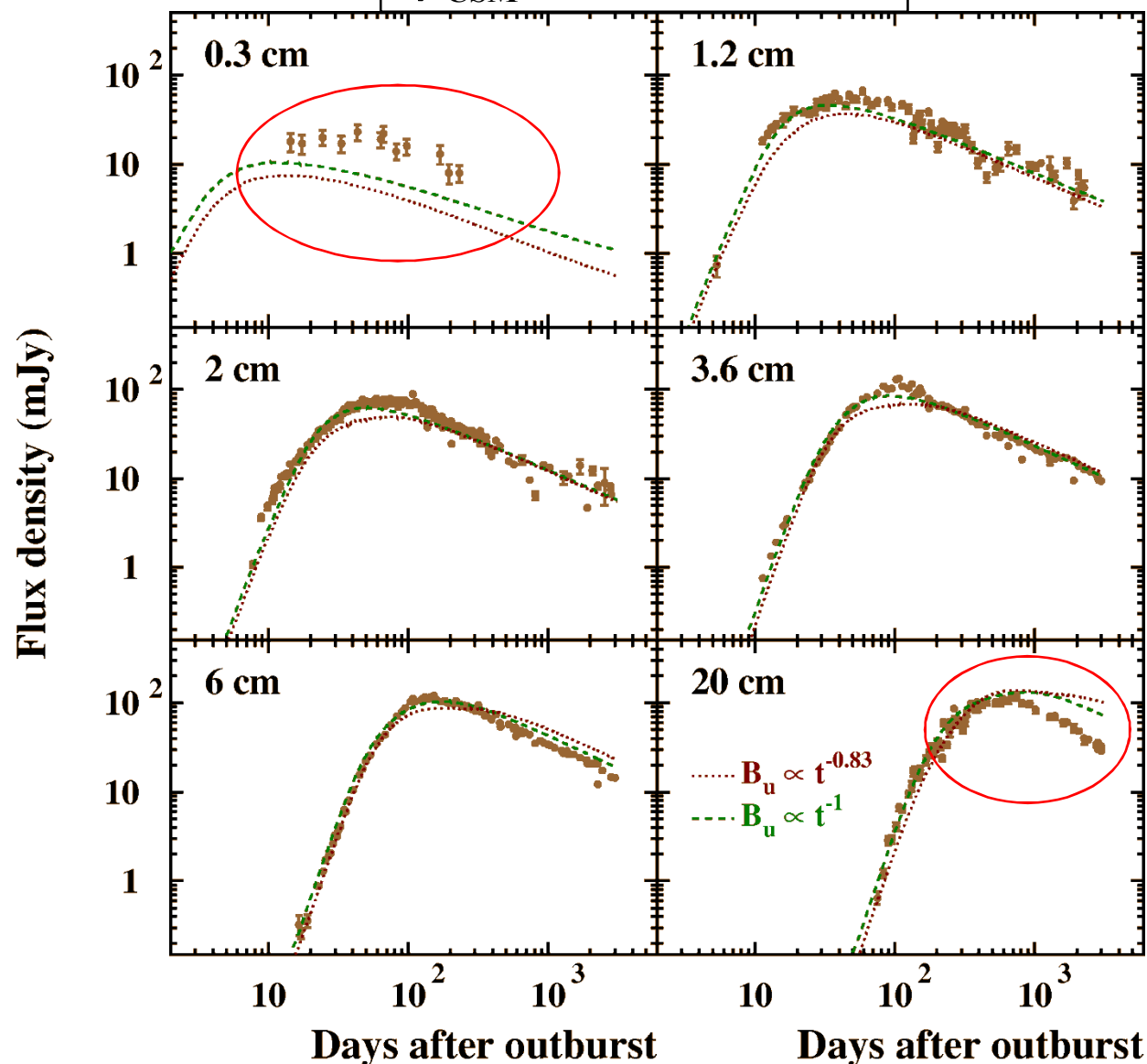
$$B_{u0} = 50 \pm 20 \text{ G}$$

- The shock is weakly modified:

$$5 \times 10^{-5} < \eta_{inj}^p < 2 \times 10^{-4}$$

Density profile of the CSM

$$\rho_{\text{CSM}} \propto r^{-s} \text{ with } s = 1.6$$



- Van Dyk et al. (1994), Fransson et al. (1996), Immler et al. (2001), Weiler et al. (2007): the CSM density profile is flatter, $s=1.5-1.7$, than the standard $s=2$ case

- Fransson & Björnsson (1998, 2005): $s=2$

- With the present model, the **optically thin** emission cannot be reproduced with $s=1.6 \Rightarrow s=2$